

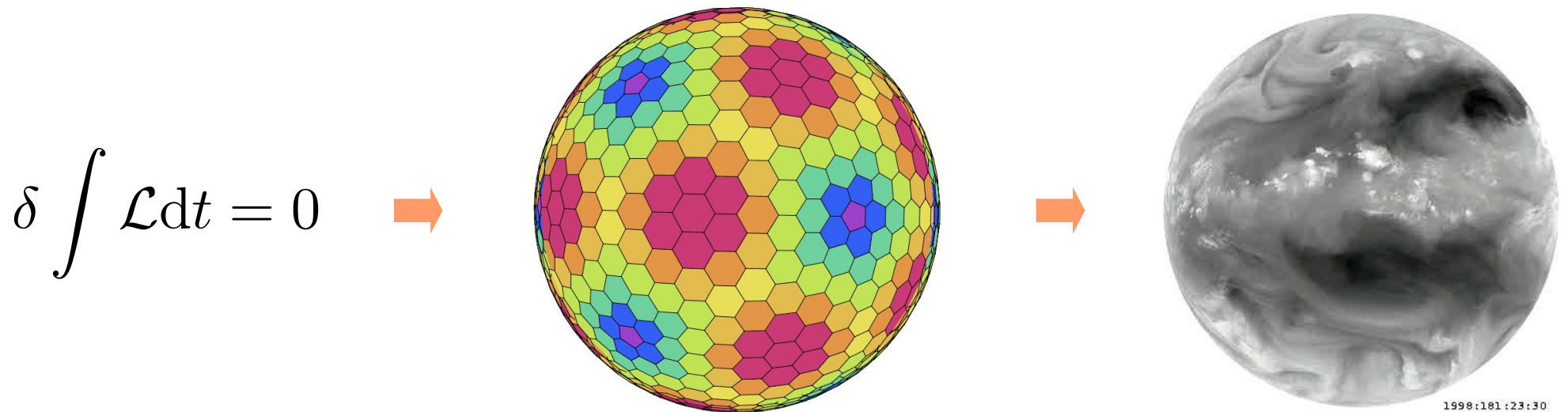
The DYNAMICO dynamical core : theory, implementation and outlook

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DYNAMICO

Equations of motion	<i>shallow-water</i> <i>shallow-atmosphere, hydrostatic</i> <i>ongoing : deep-atmosphere, fully compressible</i>
Conservation properties	<i>Mass (air and species)</i> <i>Energy</i>
Formulation	<i>Mass : flux-form</i> <i>Momentum : Hamiltonian vector-invariant form</i>
Vertical coordinate	<i>Terrain-following mass-based</i>
Numerics	<i>Mass : finite volume</i> <i>Momentum : low-order mimetic finite difference</i> <i>Mesh : icosahedral-hexagonal C-grid</i> <i>Time : (additive) Runge-Kutta (HEVI)</i>
Computing	<i>MPI / OpenMP</i> <i>Scales at least to $O(10^4)$, including I/O</i>

- Hamiltonian formulations
- DYNAMICO status

Spherical geoid
Shallow-atmosphere

Traditional

(Quasi-)Hydrostatic

Semi-hydrostatic

Boussinesq / Anelastic /

Pseudo-incompressible

Boussinesq

Anelastic
(Ogura & Phillips)

Pseudo-incompressible
(Durran ; Klein & Pauluis)

Acoustic
Lamb
Inertia-gravity
Rossby

Compressible
Euler

Spherical-geoid
Euler

Non-traditional shallow-
atmosphere
(Tort & Dubos, 2014a)

Traditional shallow-
atmosphere
(Phillips, 1966)

Acoustic
Lamb
Inertia-gravity
Rossby

Quasi-hydrostatic
(White & Wood, 2012)
(Tort & Dubos, 2014b)

Spherical-geoid
Quasi-hydrostatic
(White & Wood, 1995)

NT shallow-atmosphere
quasi-hydrostatic
(Tort & Dubos, 2014a)

Primitive equations
(Richardson, 1922)

Semi-hydrostatic
(Dubos & Voitus, 2014)

Acoustic
Lamb
Inertia-gravity
Rossby

Least action principle

$$\delta \int L(\mathbf{x}, \overset{\text{density}}{\rho}, \overset{\text{entropy}}{s}, \overset{\text{velocity}}{\mathbf{u}}) dm dt = 0$$

$$(x, y, z) \rightarrow (\xi^1, \xi^2, \xi^3)$$

$$\mathbf{u} \begin{cases} \rightarrow u^i = D\xi^i / Dt \\ \rightarrow v_i = \partial L / \partial u^i \end{cases}$$

Euler-Lagrange equations
Tort & Dubos,
JAS 2014

$$\frac{Dv_i}{Dt} - \frac{\partial L}{\partial \xi^i} = \frac{1}{\mu} \partial_i \left(\mu^2 \frac{\partial L}{\partial \mu} \right)$$

L=kinetic+Coriolis
-internal-potential

contravariant relative velocity
 $D/Dt = \partial_t + u^i \partial_i$

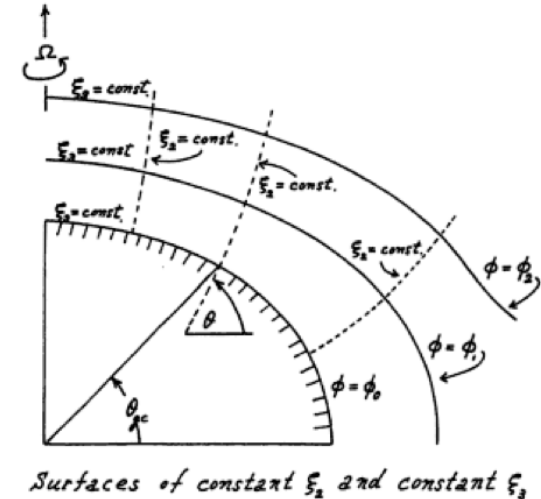
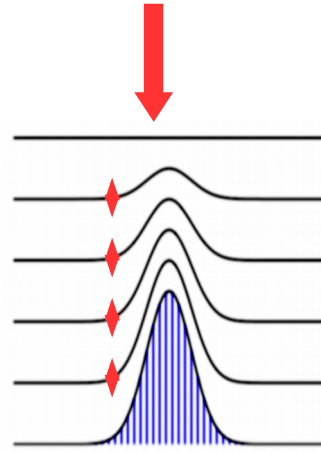
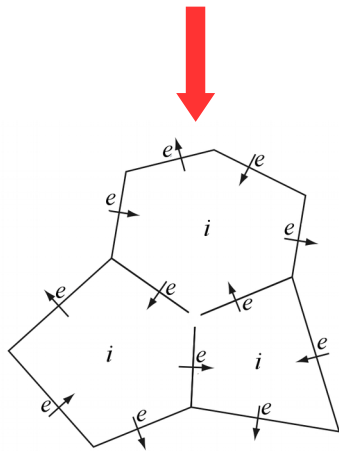
covariant absolute velocity
 $v_i = G_{ij} u^j + R_j$

Hamiltonian formulation
Dubos & Tort, MWR
2014

Curl form
 $\frac{\partial v_i}{\partial t} = \dots$

Flux form
 $\frac{\partial}{\partial t} (\mu v_i) = \dots$

Discretize !



Kinematics / dynamics separation

Mass budget in **shallow-atmosphere** vs **deep-atmosphere** geometry :

$$\frac{\partial \rho}{\partial t} + \frac{1}{a \cos \phi} \frac{\partial}{\partial \lambda} \rho u + \frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \rho \cos \phi v + \frac{\partial}{\partial r} \rho w = 0 \quad \begin{aligned} u &= a \cos \phi \dot{\lambda}, v = a \dot{\phi}, w = \dot{r} \\ \mu &= a^2 \cos \phi \rho \end{aligned}$$

$$\frac{\partial \rho}{\partial t} + \frac{1}{r \cos \phi} \frac{\partial}{\partial \lambda} \rho u + \frac{1}{r \cos \phi} \frac{\partial}{\partial \phi} \rho \cos \phi v + r^{-2} \frac{\partial}{\partial r} \rho r^2 w = 0 \quad \begin{aligned} u &= r \cos \phi \dot{\lambda}, v = r \dot{\phi}, w = \dot{r} \\ \mu &= r^2 \cos \phi \rho \end{aligned}$$

$$\frac{\partial \mu}{\partial t} + \frac{\partial}{\partial \lambda} (\mu \dot{\lambda}) + \frac{\partial}{\partial \phi} (\mu \dot{\phi}) + \frac{\partial}{\partial r} (\mu \dot{r}) = 0 \quad \begin{aligned} (\lambda, \phi) &\rightarrow \mathbf{x} \in S^2 \\ r &\rightarrow \eta \end{aligned}$$

Kinematics / dynamics separation

- Transport works in a metric-independent computational space $(\mathbf{x}, \eta) \in S^2 \times [0, 1]$
- Metric is needed only when converting between covariant and contravariant
- Diagnosing contravariant (flux) from covariant (momentum) is the business of the dynamics

=> separation of concerns

$$\begin{aligned}\frac{\partial \mu}{\partial t} + \frac{\partial}{\partial \mathbf{x}} (\mu \dot{\mathbf{x}}) + \frac{\partial}{\partial \eta} (\mu \dot{\eta}) &= 0 \\ \frac{\partial}{\partial t} \mu \theta + \frac{\partial}{\partial \mathbf{x}} (\theta \mu \dot{\mathbf{x}}) + \frac{\partial}{\partial \eta} (\theta \mu \dot{\eta}) &= 0\end{aligned}$$

Energy / vorticity conserving schemes and the curl (vector-invariant) form

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} + \mathbf{f} \times \mathbf{u} + \nabla \Phi + \frac{1}{\rho} \nabla p = 0$$



$$\mathbf{u} \cdot \nabla \mathbf{u} = (\nabla \times \mathbf{u}) \times \mathbf{u} + \nabla \frac{\mathbf{u} \cdot \mathbf{u}}{2}$$

$$\partial_t \mathbf{u} + (\mathbf{f} + \nabla \times \mathbf{u}) \times \mathbf{u} + \nabla \left(\frac{\mathbf{u} \cdot \mathbf{u}}{2} + \Phi \right) + \frac{1}{\rho} \nabla p = 0$$

- Curl-form : standard starting point to devise energy/vorticity-conserving schemes
- intimately connected to Hamiltonian formulation (Salmon, 1983 ; Gassmann, 2012)
- for **compressible flow**, Hamiltonian formulation well-established if coordinate system is **Eulerian**

A change of perspective

- Usual perspective : choose *prognostic fields* and write down *evolution equations*
- Now suppose *evolution equations are known for all prognostic variables*. Then how an arbitrary functional of those prognostic variables evolves is entirely known :

$$\frac{d}{dt}\mathcal{F}[\mathbf{v}, \mu, \Theta] = \int \left(\frac{\delta\mathcal{F}}{\delta v_i} \partial_t v_i + \frac{\delta\mathcal{F}}{\delta\mu} \partial_t \mu + \frac{\delta\mathcal{F}}{\delta\Theta} \partial_t \Theta \right) dx^1 dx^2 dx^3$$

- Conversely, *such an evolution equation contains all the equations of motion*

New perspective : write a single evolution equation for an arbitrary functional of prognostic fields

- Remark : by design $\frac{d}{dt}\mathcal{F}$ is linear in the functional derivatives of $\mathcal{F}$

Hamiltonian formulation

- A Hamiltonian formulation is such an evolution equation of the special form :

$$\frac{d}{dt}\mathcal{F} = \{\mathcal{F}, \mathcal{H}\}$$

where $\mathcal{H}$ is total energy and the Poisson bracket is bilinear, antisymmetric and satisfies the Jacobi identity

- nice, but where is my curl form ?

The curl (vector-invariant) form and Hamiltonian formulation

$$\partial_t \mathbf{u} + (\mathbf{f} + \nabla \times \mathbf{u}) \times \mathbf{u} + \nabla \left(\frac{\mathbf{u} \cdot \mathbf{u}}{2} + \Phi \right) + \frac{1}{\rho} \nabla p = 0$$

$$\partial_t v_i + (\partial_j v_i - \partial_i v_j) u^j + \partial_i (k + \Phi) + \theta \partial_i \pi = 0$$

$$\frac{d}{dt} \mathcal{F} = \{ \mathcal{F}, \mathcal{H} \}$$

- means that all tendencies can be written as expressions that are *linear in the functional derivatives of total energy*

$$\mathcal{H} [v_i, \mu, \Theta] = \int \left(\frac{1}{2} g^{ij} (v_i - R_i)(v_j - R_j) + \Phi + e \left(\alpha = \frac{J}{\mu}, \theta \frac{\Theta}{\mu} \right) \right) \mu dx^1 dx^2 dx^3$$

- To do that let us first find those functional derivatives : $v_i = g_{ij} u^j + R_i, k = \frac{1}{2} g_{ij} u^i u^j$

$$\frac{d}{dt} \mathcal{H} = \int (\mu u^i \partial_t v_i + (k + \Phi) \partial_t \mu + \pi \partial_t \Theta) dx^1 dx^2 dx^3$$

The curl (vector-invariant) form and Hamiltonian formulation

$$\begin{aligned}\partial_t \mu + \partial_i \frac{\delta \mathcal{H}}{\delta v_i} &= 0 \\ \partial_t \Theta + \partial_i \left(\theta \frac{\delta \mathcal{H}}{\delta v_i} \right) &= 0 \\ \partial_t v_i + \frac{\partial_j v_i - \partial_i v_j}{\mu} \frac{\delta \mathcal{H}}{\delta v_j} + \partial_i \frac{\delta \mathcal{H}}{\delta \mu} + \theta \partial_i \frac{\delta \mathcal{H}}{\delta \Theta} &= 0\end{aligned}$$

$$\frac{d}{dt} \mathcal{F} = \{ \mathcal{F}, \mathcal{H} \}$$

- means that all tendencies can be written as expressions that are *linear in the functional derivatives of total energy*

$$\mathcal{H} [v_i, \mu, \Theta] = \int \left(\frac{1}{2} g^{ij} (v_i - R_i)(v_j - R_j) + \Phi + e \left(\alpha = \frac{J}{\mu}, \theta \frac{\Theta}{\mu} \right) \right) \mu dx^1 dx^2 dx^3$$

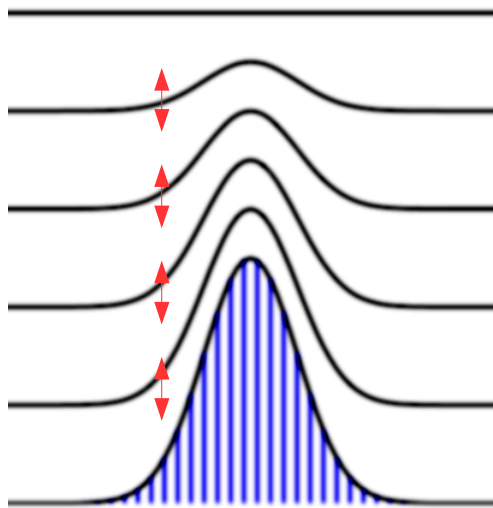
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$$\frac{d}{dt} \mathcal{H} = \int (\mu u^i \partial_t v_i + (k + \Phi) \partial_t \mu + \pi \partial_t \Theta) dx^1 dx^2 dx^3$$

Energy / vorticity conserving schemes and the curl (vector-invariant) form

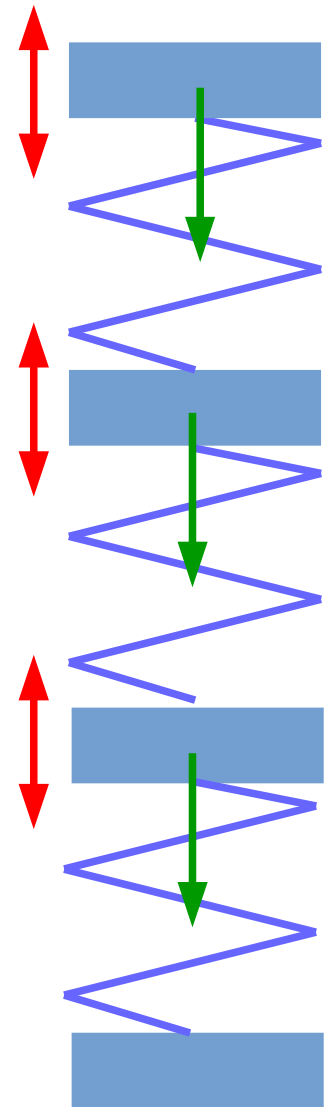
$$\partial_t \mathbf{u} + (\mathbf{f} + \nabla \times \mathbf{u}) + \nabla \left(\frac{\mathbf{u} \cdot \mathbf{u}}{2} + \Phi \right) + \frac{1}{\rho} \nabla p = 0$$

- curl-form standard starting point to devise energy/vorticity-conserving schemes
- intimately connected to Hamiltonian formulation (Salmon, 1983 ; Gassmann, 2012)
- for **compressible flow**, Hamiltonian formulation well-established if coordinate system is **Eulerian**



We want to handle **non-Eulerian** vertical coordinates :

- Isentropic/isopycnal
- based on mass / hydrostatic pressure
- natural representation of **hydrostatic adjustment**



*Hamiltonian formulation for generalized vertical coordinates :
Dubos & Tort (MWR 2014)*

Hamiltonian formulation in generalized vertical coordinates

(Dubos & Tort, MWR 2014)

$$\partial_t \mu + \partial_\eta (\mu \dot{\eta}) + \partial_i \frac{\delta \mathcal{H}}{\delta v_i} = 0,$$

$$\partial_t \Theta + \partial_\eta (\Theta \dot{\eta}) + \partial_i \left(\theta \frac{\delta \mathcal{H}}{\delta v_i} \right) = 0,$$

$$\partial_t v_i + \dot{\eta} \partial_\eta v_i + v_3 \partial_i \dot{\eta} + \frac{\partial_j v_i - \partial_i v_j}{\mu} \frac{\delta \mathcal{H}}{\delta v_j} + \partial_i \frac{\delta \mathcal{H}}{\delta \mu} + \theta \partial_i \left(\frac{\delta \mathcal{H}}{\delta \Theta} \right) = 0,$$

Integration by parts

+ invariance w.r.t. vertical coordinate

=> conservation of energy

$$\partial_t V_3 + \partial_\eta (V_3 \dot{\eta}) + \frac{\delta \mathcal{H}}{\delta \Phi} = 0,$$

$$\partial_t \Phi + \dot{\eta} \partial_\eta \Phi - \frac{\delta \mathcal{H}}{\delta V_3} = 0.$$

Isentropic / Isopycnal

$$\dot{\eta} = 0$$

Mass-based

Diagnosed from
horizontal mass flux

Eulerian

$$\partial_t \Phi = 0$$

Hamiltonian formulation in generalized vertical coordinates

(Dubos & Tort, MWR 2014)

$$\partial_t \mu + \partial_\eta (\mu \dot{\eta}) + \partial_i \frac{\delta \mathcal{H}}{\delta v_i} = 0,$$

$$\partial_t \Theta + \partial_\eta (\Theta \dot{\eta}) + \partial_i \left(\theta \frac{\delta \mathcal{H}}{\delta v_i} \right) = 0,$$

$$\partial_t v_i + \dot{\eta} \partial_\eta v_i + \cancel{v_3 \partial_\eta \eta} + \frac{\partial_j v_i - \partial_i v_j}{\mu} \frac{\delta \mathcal{H}}{\delta v_j} + \partial_i \frac{\delta \mathcal{H}}{\delta \mu} + \theta \partial_i \left(\frac{\delta \mathcal{H}}{\delta \Theta} \right) = 0,$$

Hydrostatic $v_3 \equiv \frac{\partial L}{\partial u^3} = 0$

$$\cancel{\partial_t V_3} + \partial_\eta (\cancel{V_3 \dot{\eta}}) + \frac{\delta \mathcal{H}}{\delta \Phi} = 0,$$

$$\cancel{\partial_t \Phi} + \dot{\eta} \cancel{\partial_\eta \Phi} - \frac{\delta \mathcal{H}}{\delta V_3} = 0.$$

Isentropic / Isopycnal

$$\dot{\eta} = 0$$

Mass-based

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Eulerian

$$\partial_t \Phi = 0$$

Physical space

$$\lambda, \varphi, \Phi, g$$

$$u, v, w$$

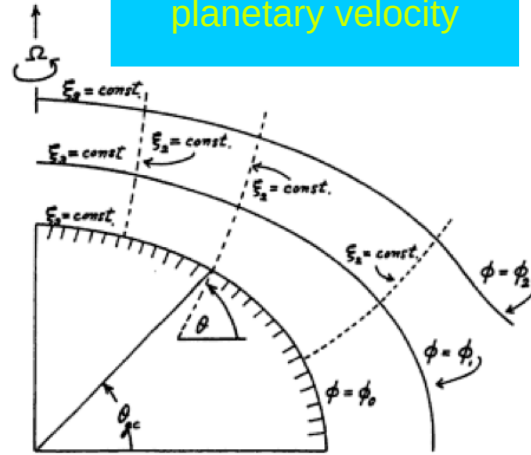
$$\alpha = 1/\rho, s, r$$

Thermodynamics

$$p, T, \chi$$

Geopotential
Coordinates

Metric, gravity,
planetary velocity



Surfaces of constant ξ_2 and constant ξ_3

Computational space $S^2 \times [0,1]$

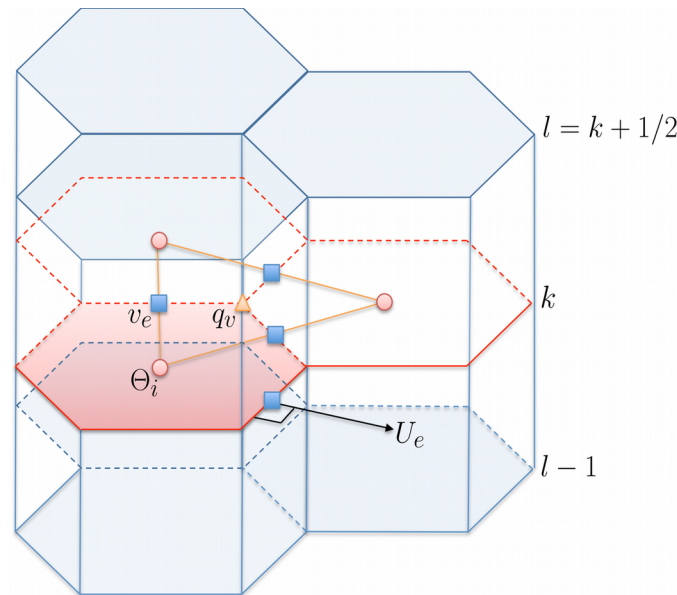
$$\mathbf{v} = \mathbf{G}(\mathbf{x}, \Phi) \cdot \dot{\mathbf{x}} + \mathbf{R}(\mathbf{x}, \Phi)$$

$\mathbf{x}$

η

Horizontal mesh
Icosahedral C-grid

Vertical mesh
Lorenz



Discrete space

$$m_{ik} = \int \int \int \mu d\mathbf{x} d\eta$$

$$F_{il} = \int \int \mu \dot{\eta} d\mathbf{x}$$

$$v_{ek} = \int \mathbf{v} \cdot d\mathbf{x}$$

$$\alpha_{ik} = \alpha(p_{ik}, s_{ik}),$$

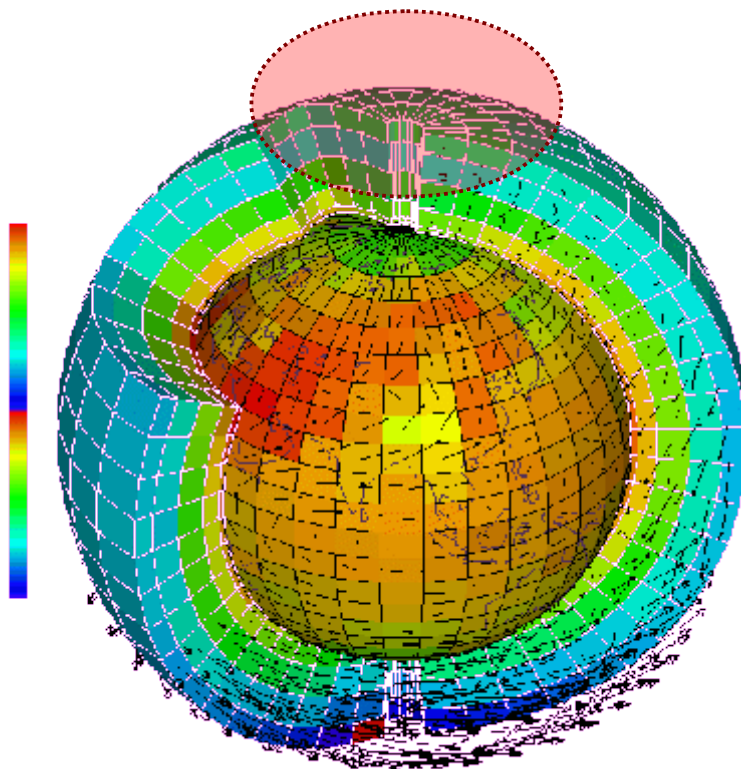
- Discrete integration by parts (Bonaventura & Ringler, 2005 ; Taylor, 2010)
- Energy- and vorticity- conserving Coriolis discretization (TRiSK : Thuburn et al., 2009 ; Ringler et al., 2010)

Energy-conserving
3D core

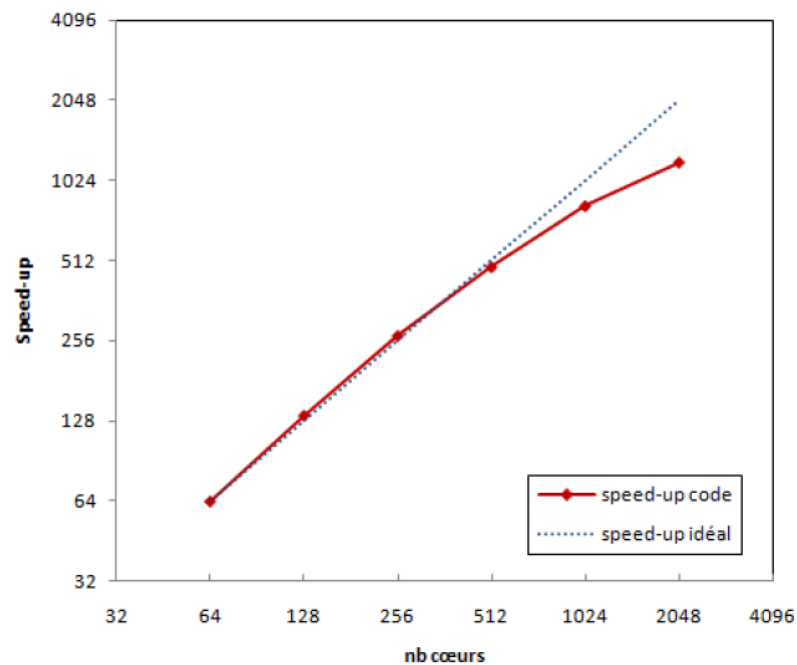
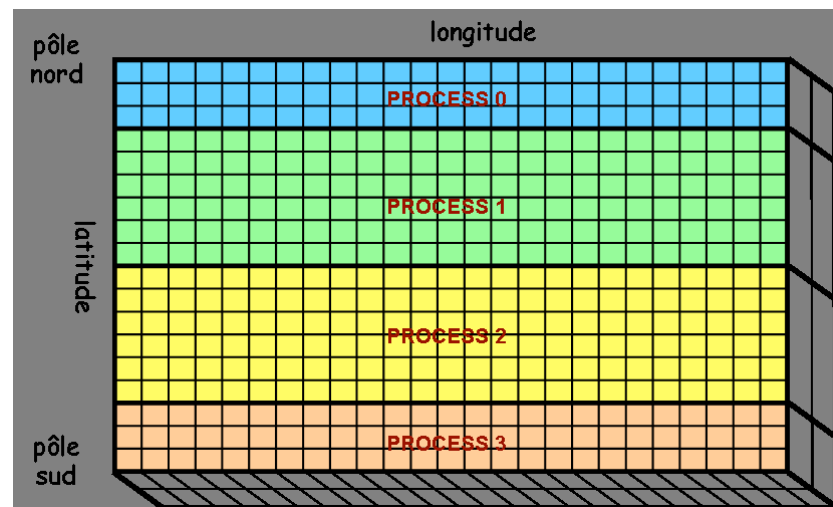
Hamiltonian structure of the equations of motion can :

- Guide the choice of prognostic variables
- Help with discrete conservation of energy
(up to time truncation errors)
- Create opportunities for *modularity / separation of concerns* :
 - Bracket (general structure) independent from specific equation set
 - Hamiltonian specific to certain equation set :
equation of state, shallow-atmosphere vs deep-atmosphere
- Implemented in DYNAMICO with FD but other building blocks possible :
 - spectral elements (horizontal)
 - compatible finite elements (GungHo)

- Hamiltonian formulations
- DYNAMICO status



LMD-Z lon-lat core



Y. Meurdesoif, 2010

1/4 degree (25km) => max. 360 MPI processes

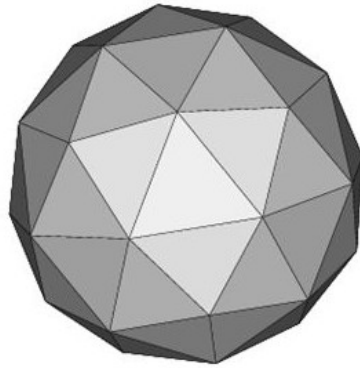
=> a couple thousand cores

Mesh partitioning for parallel computing

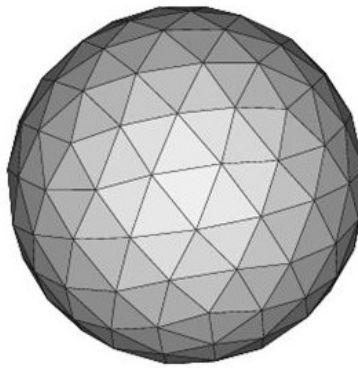
$nbp=1$



$nbp=2$



$nbp=4$

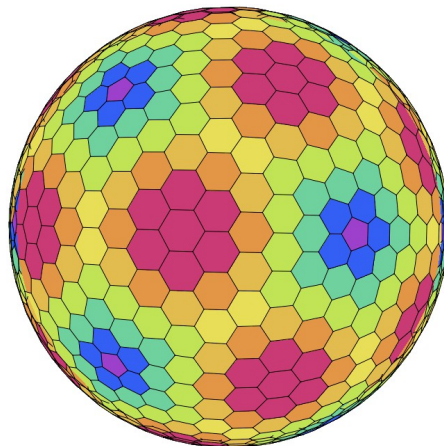


Icosahedral-triangular mesh
 $10 \cdot nbp^2 + 2$ vertices

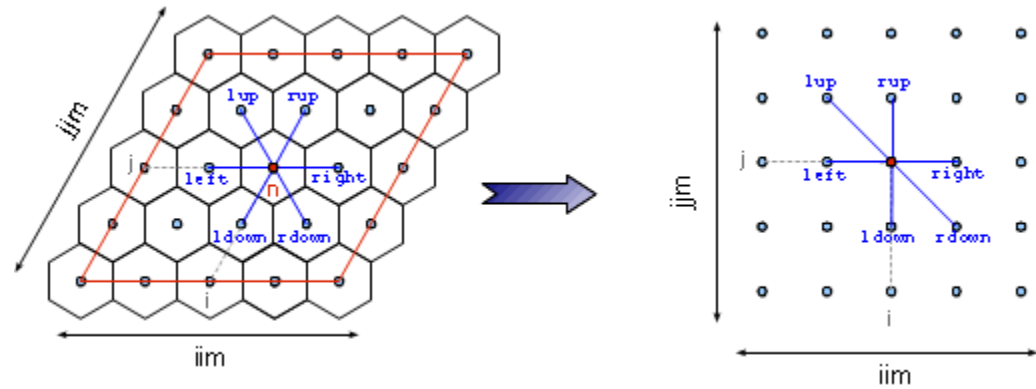
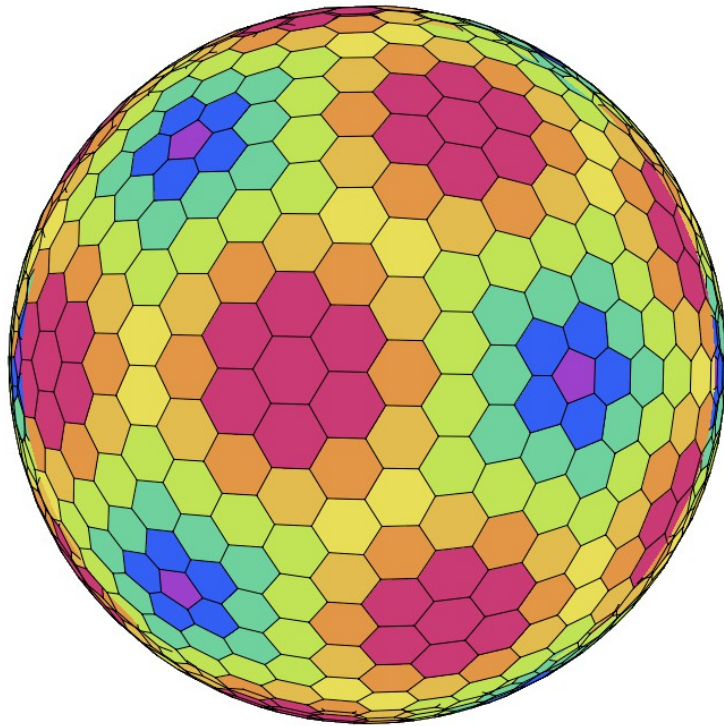
Voronoi
dual



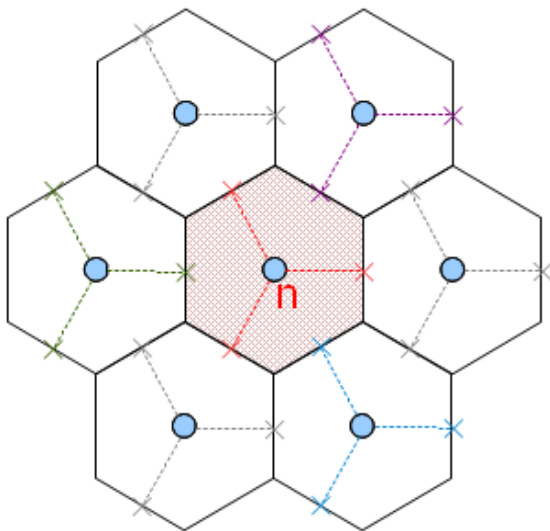
Icosahedral-hexagonal mesh
 $10 \cdot nbp^2 + 2$ cells



- Easy to partition into $10 \times nsplit^2$ domains
- About $(nbp/nsplit)^2$ cells per domain = MPI process
- $Nbp/nsplit > 10$ for performance
- Ex : 25km $\sim 10^6$ cells $\sim 10\,000$ MPI processes



- Vertical direction in outer loops
- **Direct access to neighbours via constant offsets**
- No special case for pentagons (handled by metrics)

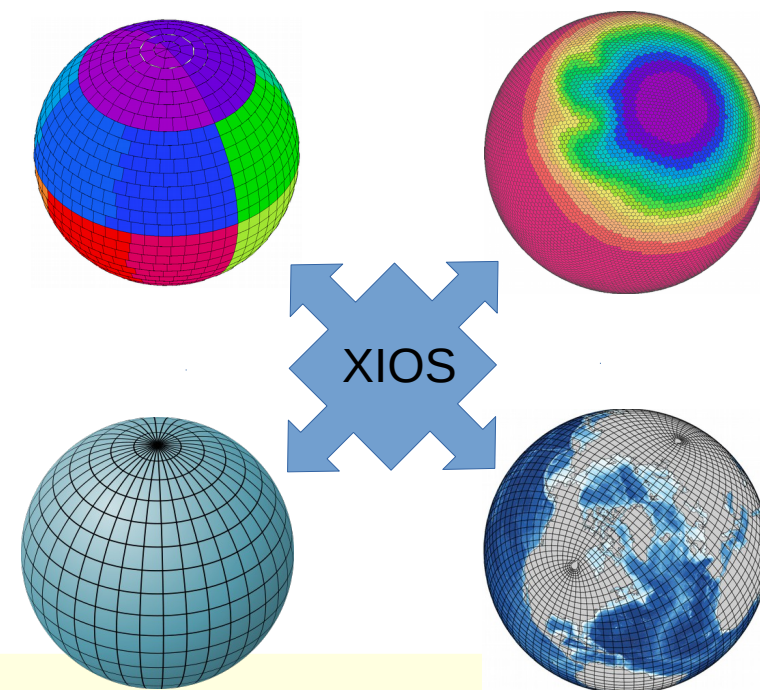
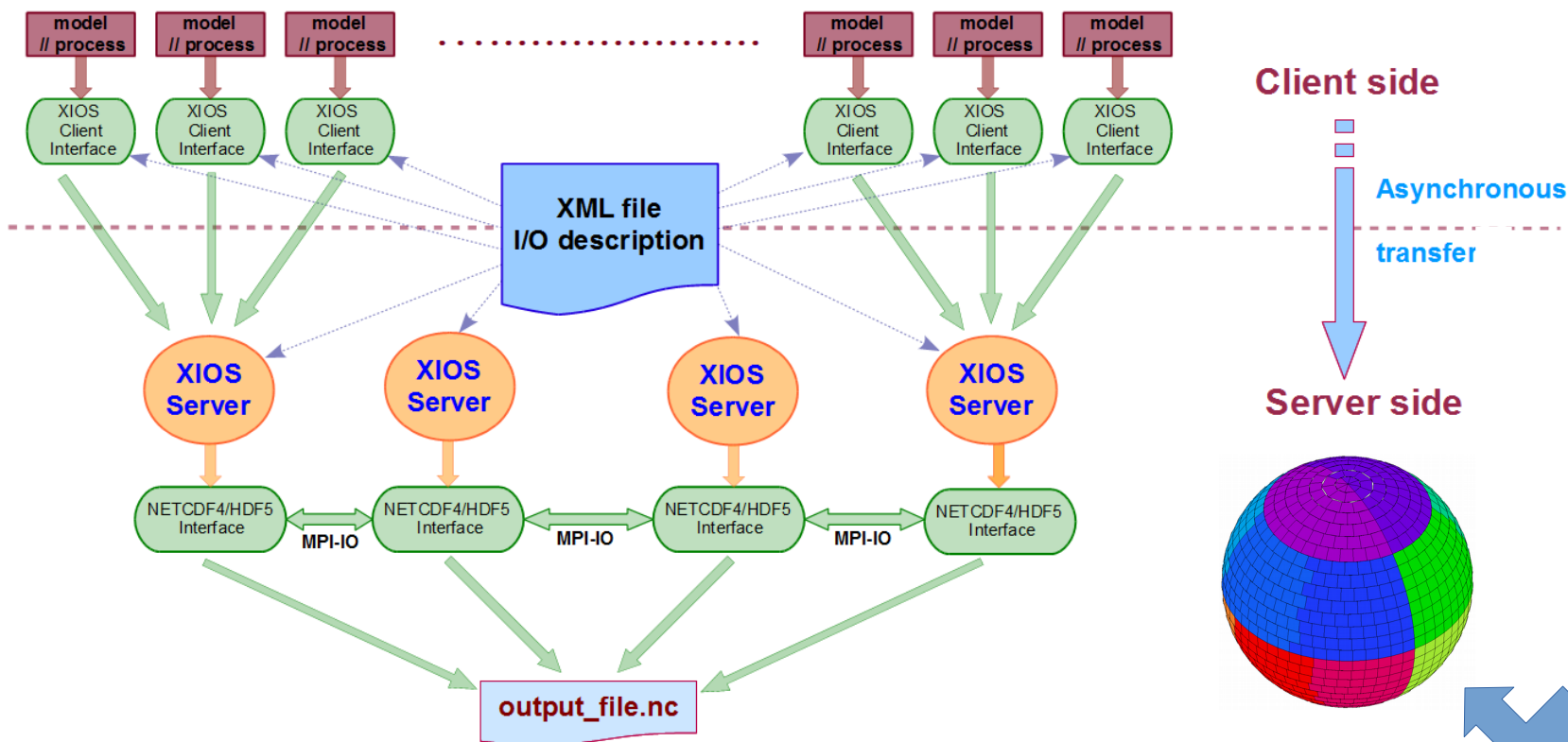


```

DO ij=ij_begin,ij_end
  ! convm = -div(mass flux), sign convention as in Ringler et al
  convm(ij,l)= -1./Ai(ij)*(ne_right*hflux(ij+u_right,l)  + &
    ne_rup*hflux(ij+u_rup,l)          + &
    ne_lup*hflux(ij+u_lup,l)          + &
    ne_left*hflux(ij+u_left,l)         + &
    ne_ldown*hflux(ij+u_ldown,l)       + &
    ne_rdown*hflux(ij+u_rdown,l))
ENDDO
  
```

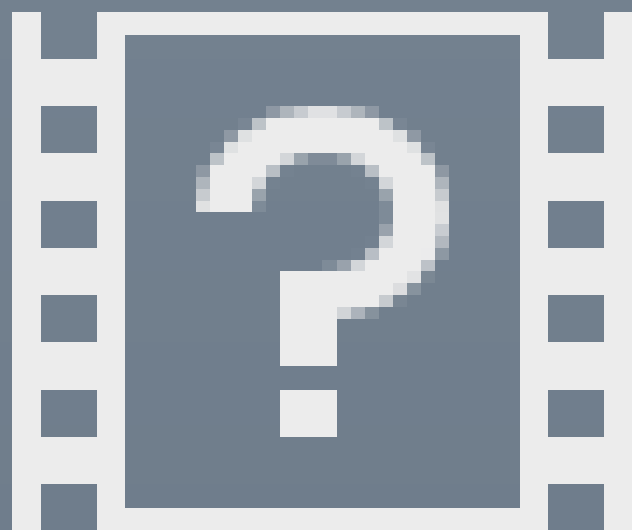

XIOS (Y. MEURDESOF) : XML I/O SERVER

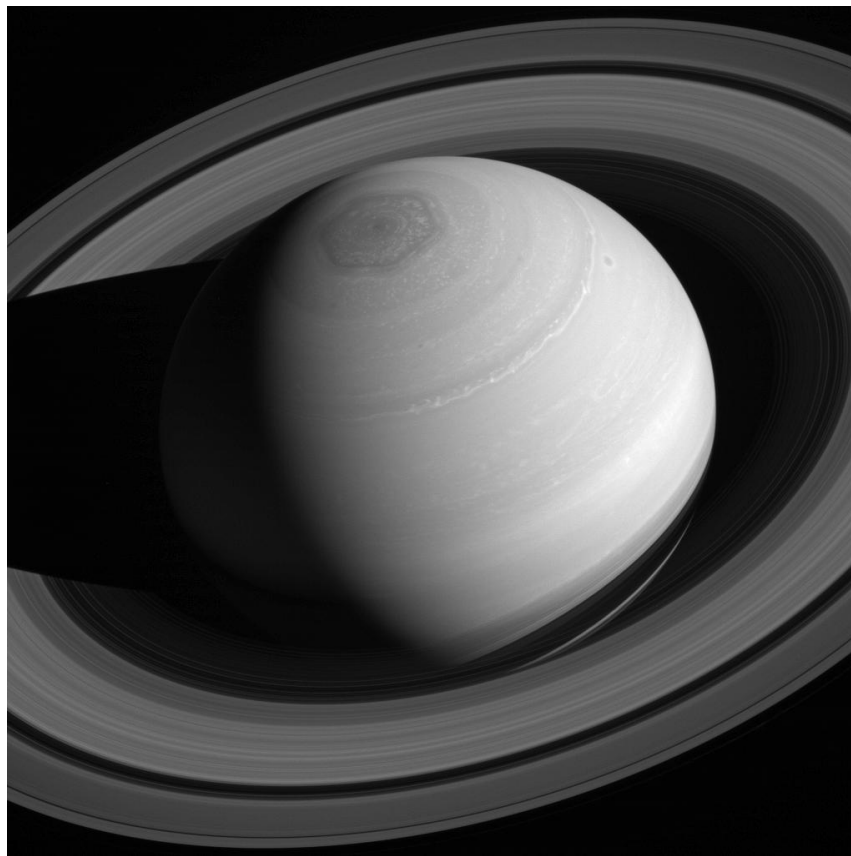
PARALLEL ASYNCHRONOUS I/O - ONLINE POST-PROCESSING LIBRARY AND SERVER



DYNAMICO

T. Dubos 22/33





Saturn GCM simulations

Grid

- Horizontal resolution: $1/2^\circ$ (+ tests $1/4^\circ$ & $1/8^\circ$)
- Vertical levels: 32 levels from 3 bars to 1 mbar

Boundary conditions

- Initial: steady-state temperature from 1D run, no winds
- Dissipation (SGS): very strong ($\mathcal{D}^+$), strong ($\mathcal{D}^=$), regular ($\mathcal{D}^-$)
- Type 1: 64 levels + sponge layer on uppermost 4 levels
- Type 2: Bottom drag $|\lambda| > 33^\circ$: $\tau = 100\text{d}$ ($\mathcal{F}^+$), 1000d ($\mathcal{F}^-$)

[Liu and Schneider JAS 2010]

Machinery

- MPI+openMP code run on Occigen cluster in CINES
- cores: 1200 ($1/2^\circ$), 9000 ($1/4^\circ$), 11520-30000 ($1/8^\circ$)
- Results are shown after 20000 Saturn days integrations

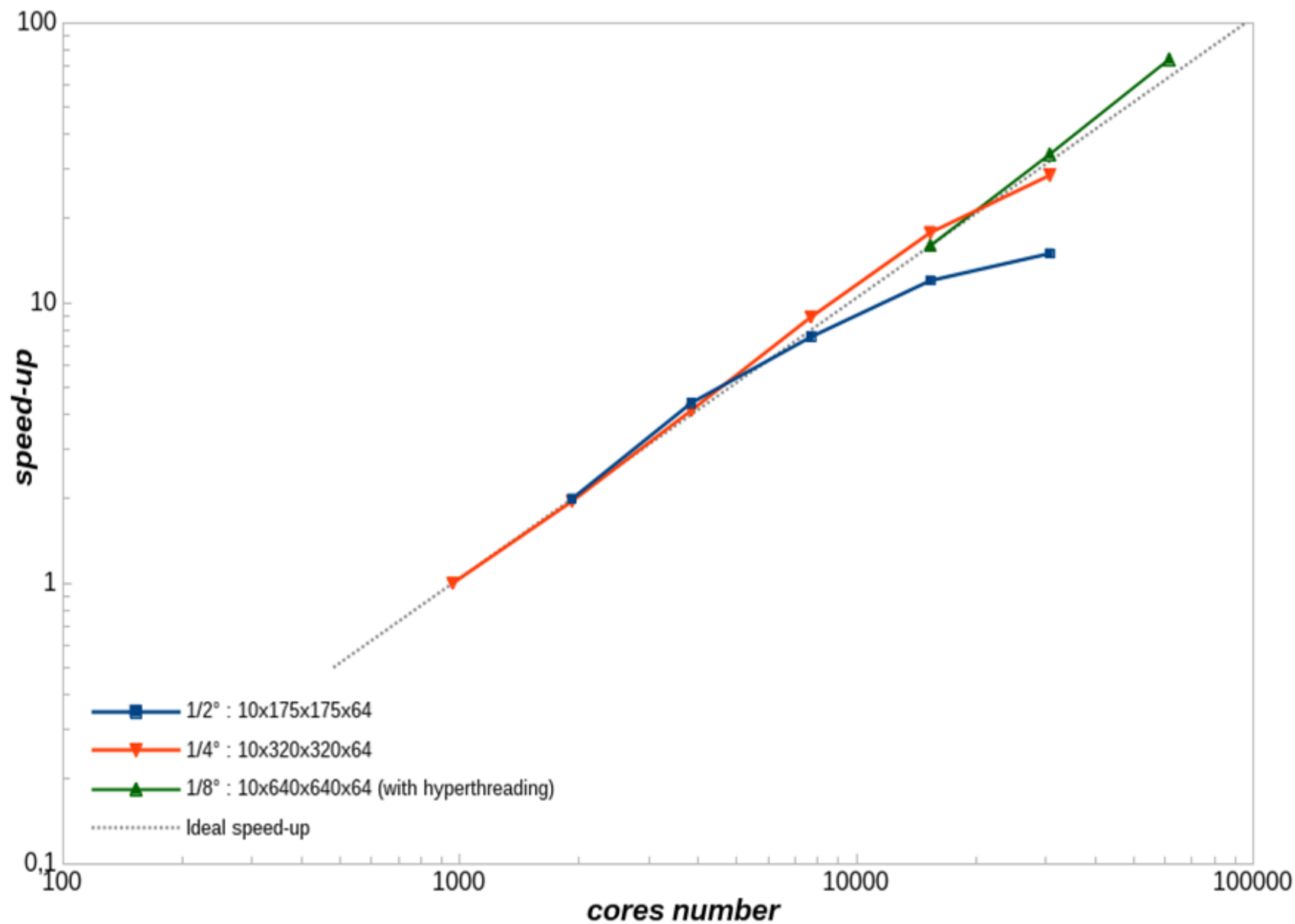
1 Sat day = 10 h ; 1 Sat year = 30 Earth yr

Physical parameterizations $\Rightarrow$ 1D computations of forcings on each grid point

☞ Radiative transfer $\Rightarrow$ Guerlet et al. Icarus 2014

- correlated- k scheme for IR and VIS heating rates [Wordsworth et al. 2010]
- gases CH_4 , C_2H_6 , C_2H_2 with optimized spectral discretization
- HITRAN 2012 database + Karkoschka and Tomasko 2010 for CH_4 around $1\mu\text{m}$
- collision-induced absorption $\text{H}_2\text{-H}_2$ and $\text{H}_2\text{-He}$ [Wordsworth et al. 2012]
- Rayleigh scattering H_2 , He
- simple two-layer aerosol model [constrained by Roman et al. 2013]
 - tropospheric haze layer 180 – 660 mbar / $\tau \sim 8$ / $r = 2\mu\text{m}$
 - stratospheric haze layer 1 – 30 mbar / $\tau \sim 0.1$ / $r = 0.1\mu\text{m}$
- free bottom surface with internal heat flux
- incoming flux: ring shadowing, oblateness

☞ Turbulent diffusion + dry convective adjustment [Hourdin et al. 1993]

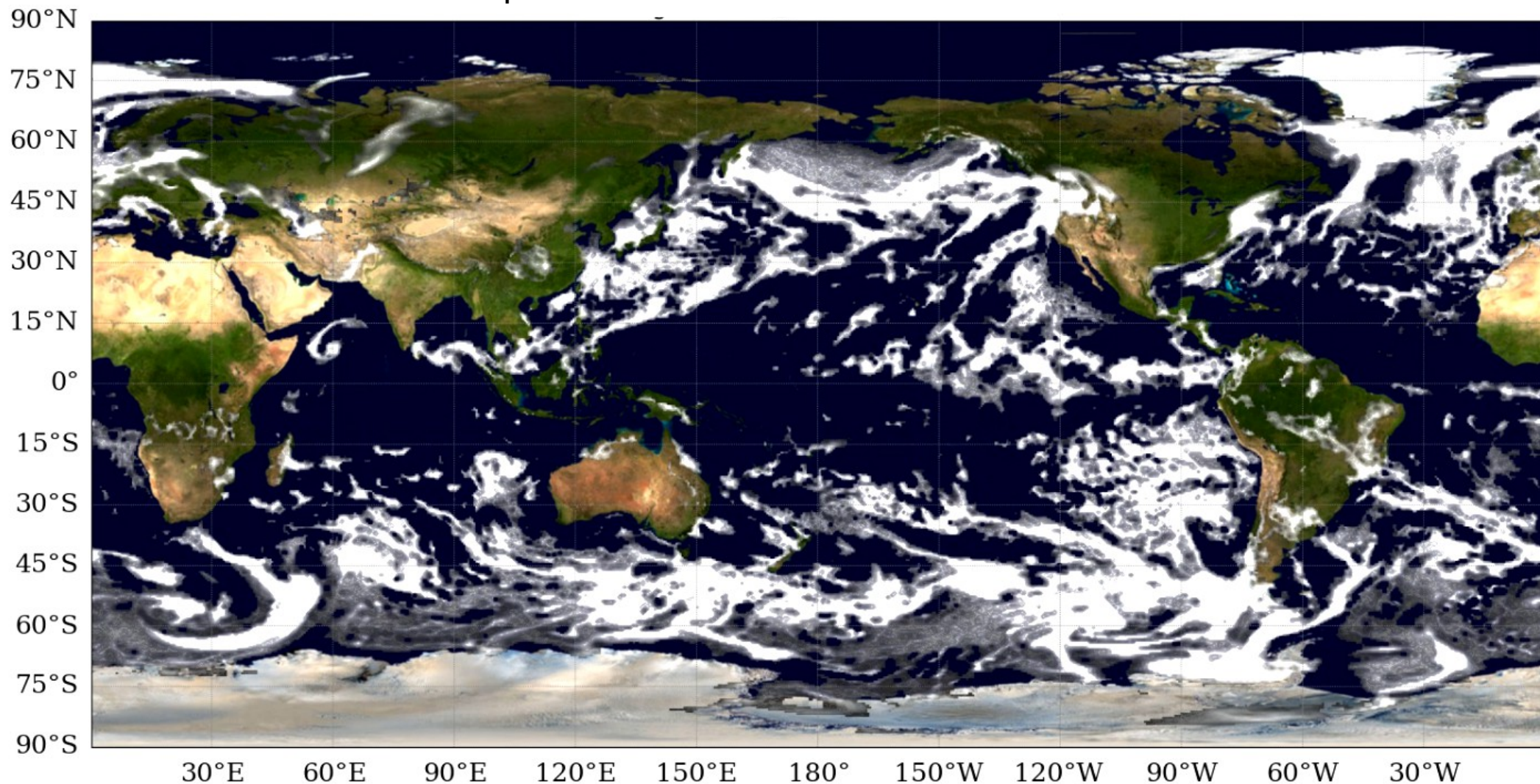


IPSL-CM7A-HR :

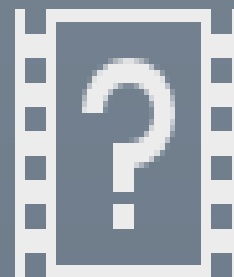
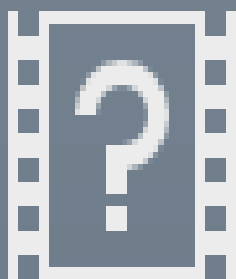
DYNAMICO-LMDZ-ORCHIDEE, 50km

5 SYPD @ 2848 cores – full physics, I/O

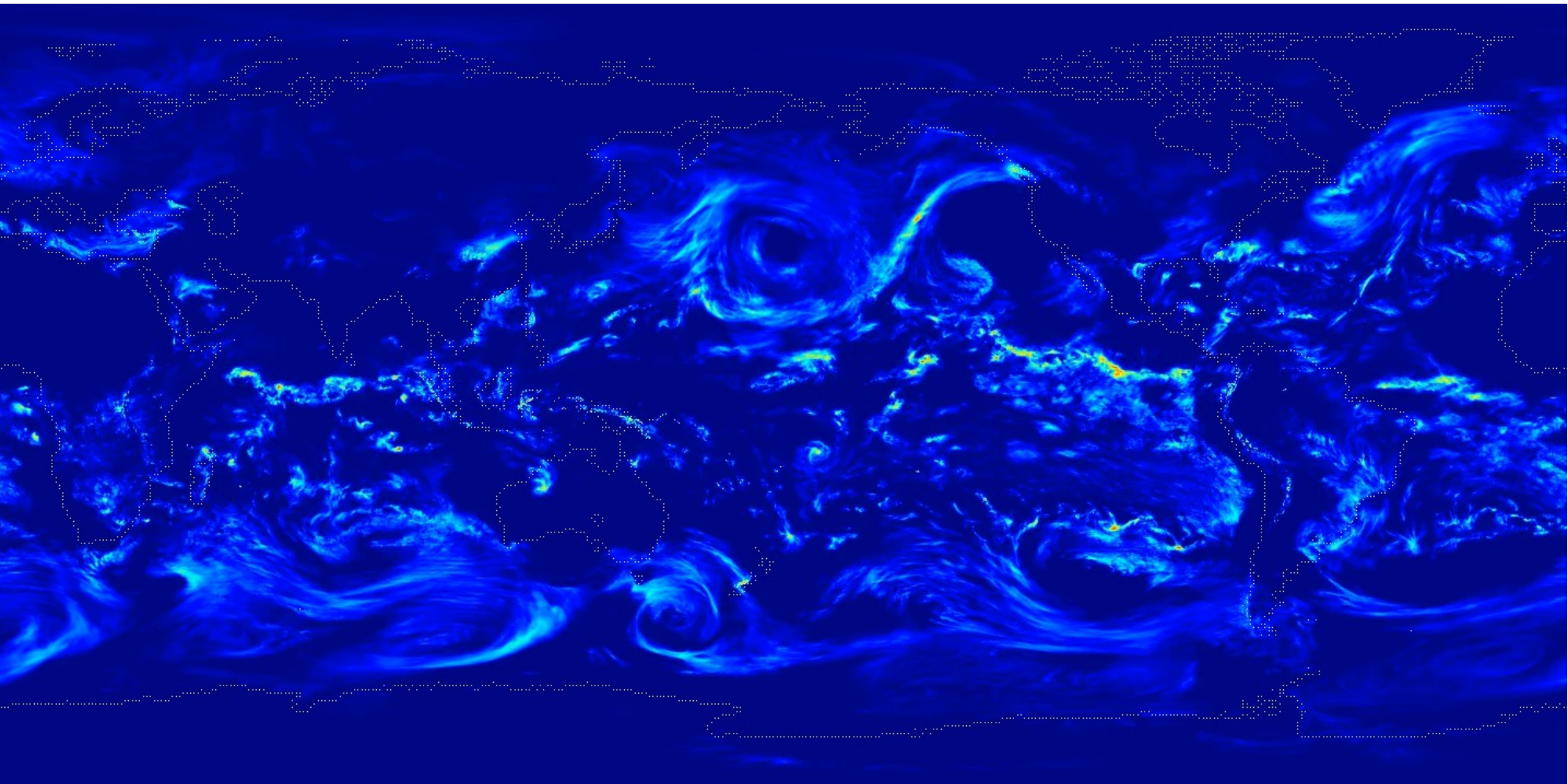
256 000 atmospheric columns x 80 vertical levels = 20 million cells



Low-level cloudiness – HighResMIP experiment

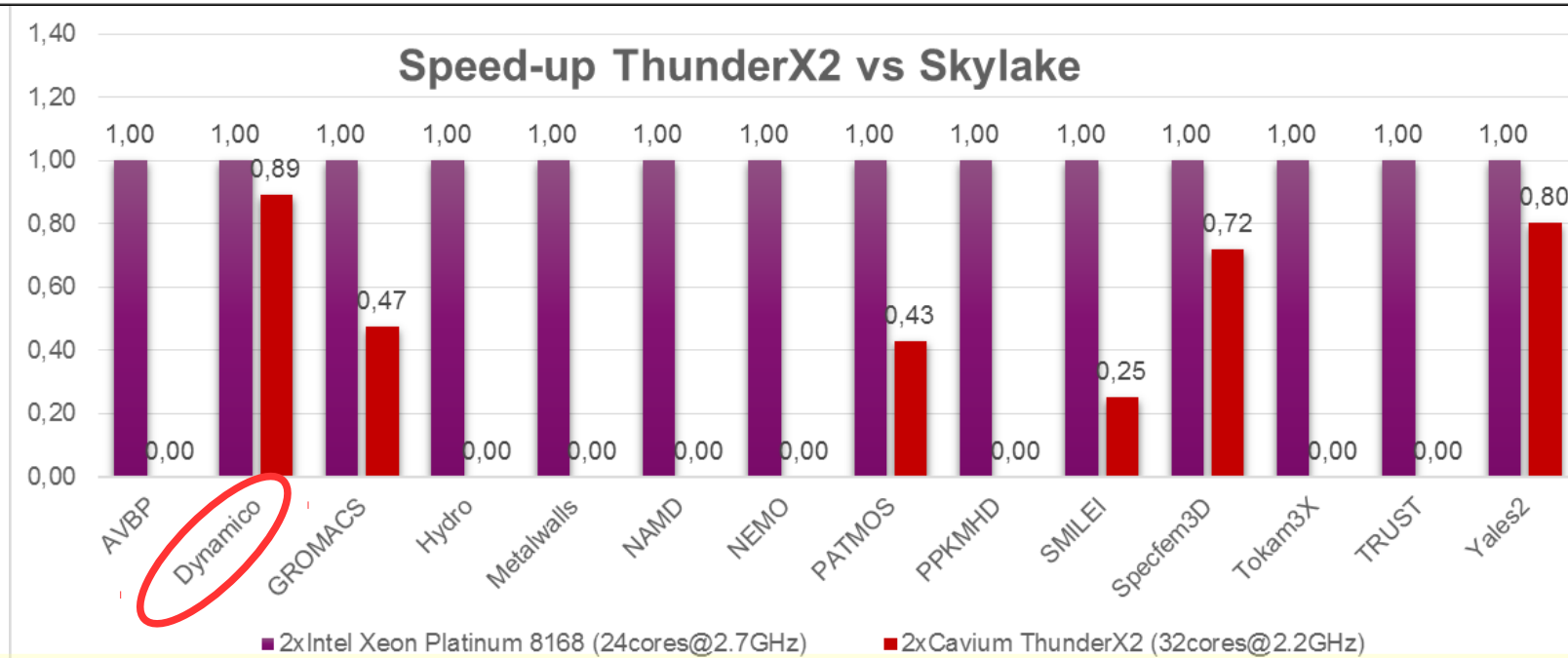
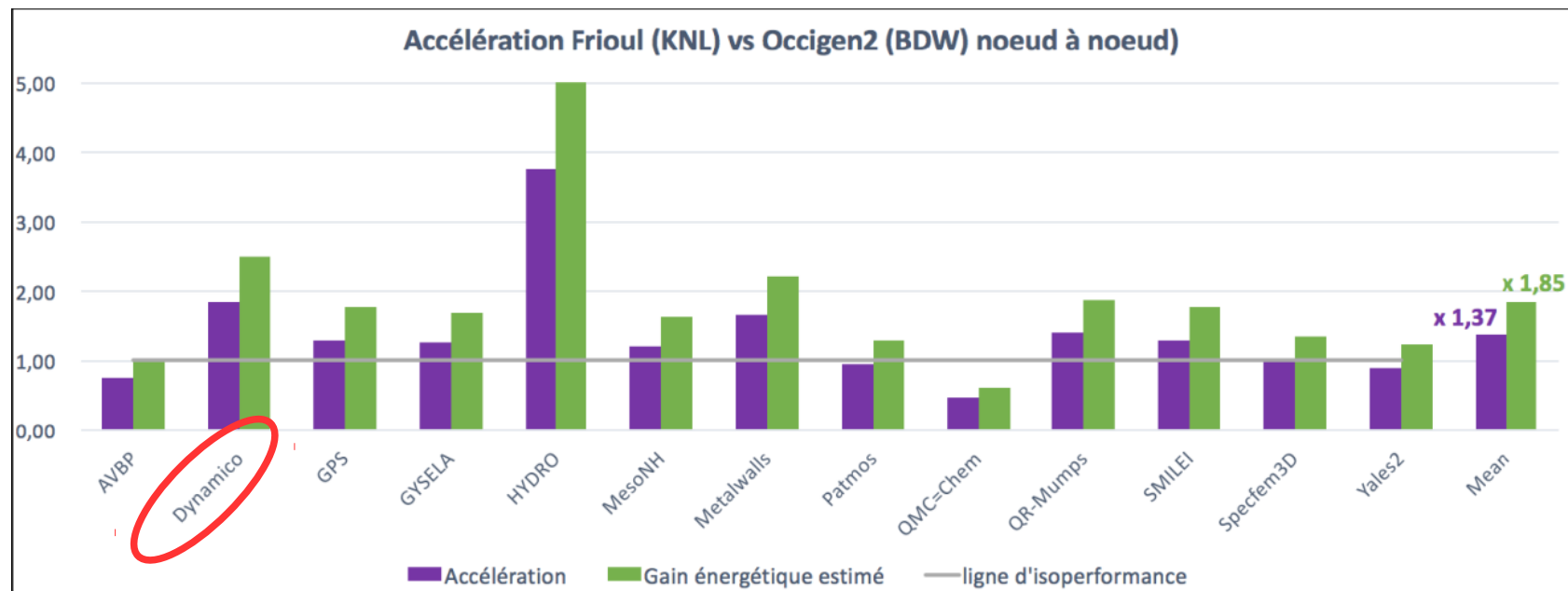


DYNAMICO-LMDZ, 25km
(1 024 000 atmospheric columns x 80 vertical levels)



Liquid water

DYNAMICO on fancy CPUs : KNL and ARM



DYNAMICO on GPUs : Nvidia V100

# columns	16 000	64 000	256 000	1 024 000
40 CPUs	68 ms	282 ms	1193 ms	5102 ms
1 GPU	26,6 ms (x 2.5)	77 ms (x3.6)	299 ms (x 4.0)	Not enough memory
2 GPUs	14,7 ms (x 4.6)	40 ms (x 7)	152 ms (x 7.9)	649 ms (x 7.9)
4 GPUs	11,5 ms (x 5.9)	24 ms (x 11.8)	75 ms (x 15.9)	304 ms (x 16.8)

```
!$acc parallel loop collapse[2]
```

```
DO 1 = ll_begin, ll_end
```

```
DO ij = ij_begin_ext, ij_end_ext
```

```
uu_right = 0.5*(rhodz[ij,l]+rhodz[ij+t_right,l])*u[ij+u_right,l]
```

```
uu_right = uu_right*le_de[ij+u_right]
```

```
hflux[ij+u_right,l] = uu_right
```

```
uu_lup = 0.5*(rhodz[ij,l]+rhodz[ij+t_lup,l])*u[ij+u_lup,l]
```

```
uu_lup = uu_lup *le_de[ij+u_lup]
```

```
hflux[ij+u_lup,l] = uu_lup
```

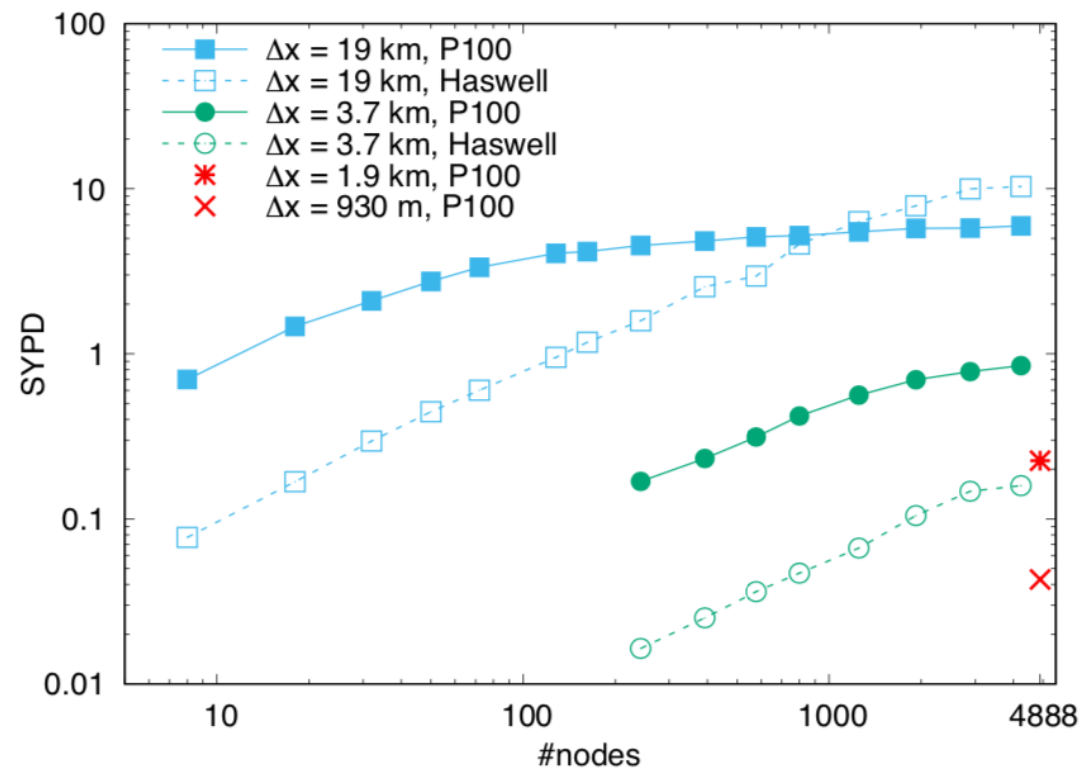
```
uu_ldown = 0.5*(rhodz[ij,l]+rhodz[ij+t_ldown,l])*u[ij+u_ldown,l]
```

```
uu_ldown = uu_ldown*le_de[ij+u_ldown]
```

```
hflux[ij+u_ldown,l] = uu_ldown
```

```
END DO
```

```
END DO
```



Fuhrer et al. 2017 <https://doi.org/10.5194/gmd-2017-230>

Prospects for higher-resolution *climate* modelling

- CMIP requires a throughput of x10000 (30SYPD)
- Some climate modelling still doable with x1000 (3SYPD)
- Ability to attain x1000 depends on maximum stable time step (numerics) and walltime needed to perform one time step (implementation)
- Assuming a large enough machine, reducing walltime is a **strong scaling problem**
- For DYNAMICO, dt (in sec) is about $2.5 \cdot dx$ (in km)
=> ms/time step required to reach 3SYPD :
 - 25km 60 ms per full time step
 - 8km 20 ms per full time step
 - 1km 2.5 ms per full time step

CMIP6 physics (79 vertical levels) cost 2-3 ms per column per call
(24 SYPD with 96 calls per day, 36 columns per core)

If physics are called every 5 dynamics time steps :

100 columns/core	=> $20 + (100 \cdot 2/5) = 60$ ms	=> 3 SYPD at 25 km with 10 000 cores
30 columns/core	=> $10 + (30 \cdot 2/5) = 22$ ms	=> 3 SYPD at 8 km with 300 000 cores ?
10 columns/core	=> $3 + (10 \cdot 2/5) = 7$ ms	=> 1 SYPD at 1 km with 50 million cores ??

Ongoing (or rather unfinished ?) work not covered in this talk :

- Coupling with ocean (NEMO) => IPSL-CM
- Fully unstructured re-implementation
=> variable resolution (à la MPAS)
- Limited-area
- Non-hydrostatic dynamics
- Deep-atmosphere / non-traditional dynamics
- Fancy equation of state (Venus)

Summary

Hamiltonian and variational structures

- help consider different equation sets in a unified way and clarify basic concepts : kinematics vs dynamics, velocity vs momentum
- separate the problem of conserving energy from other numerical properties : accuracy, dispersion, ...

DYNAMICO : now mature as part of IPSL-CM and for planets

Separation of concerns :

- currently addressed as a software design problem
- claim / feeling : also and maybe primarily a higher-level (physics, numerics) design goal
- work needed by physicists, mathematicians to split models into smaller pieces with minimal interaction (e.g. transport, thermodynamics, geometry...)

EXTRA

Hydrostatic primitive equations : continuous

shallow atmosphere traditional approximation

$\mathbf{v} = \textcolor{red}{a}^2 \dot{\mathbf{x}} + \textcolor{green}{\mathbf{R}} \quad \mathbf{R} = \textcolor{green}{a}^2 \cos \phi \Omega \mathbf{e}_\lambda$

potential energy mass-weighted potential temperature

$$H[\mu, \Theta, \Phi, \mathbf{v}] = \int \left[\left(\Phi + e \left(\frac{\textcolor{red}{a}^2 \partial_\eta \Phi}{\textcolor{red}{g} \mu}, \frac{\Theta}{\mu} \right) + \frac{1}{2} \textcolor{red}{a}^{-2} (\mathbf{v} - \textcolor{green}{\mathbf{R}})^2 \right) \mu d^2 \mathbf{x} d\eta \right]$$

Internal energy specific volume horizontal kinetic energy K_H

Bernoulli Exner Hydrostatic balance mass flux

$$\frac{dH}{dt} = \int \left[(\Phi + K_H) \partial_t \mu + \pi \partial_t \Theta + \left(\mu + \frac{a^2}{g} \partial_\eta p \right) \partial_t \Phi + \mu \dot{\mathbf{x}} \cdot \partial_t \mathbf{v} \right] d^2 \mathbf{x} d\eta$$

$\frac{\delta H}{\delta \mu}$ $\frac{\delta H}{\delta \Theta}$ $\frac{\delta H}{\delta \Phi}$ $\frac{\delta H}{\delta \mathbf{v}}$

All terms in the equations of motion can be expressed as functional derivatives of total energy

Hydrostatic primitive equations : discrete

(Lagrangian vertical coordinate – extra terms for vertical transport when mass-based)

mass-weighted
potential temperature

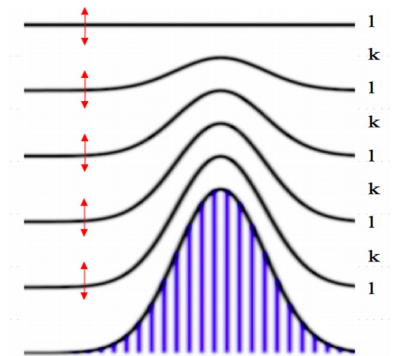
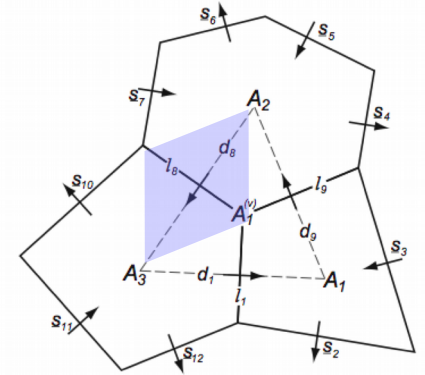
$$H = K(m_{ik}, v_{ek}) + P(m_{ik}, \Theta_{ik}, \Phi_{il})$$

$$K = \sum_{ike} m_{ik} \frac{A_{ie}}{A_i} \left(\frac{v_{ek} - R_e}{ad_e} \right)^2 \quad \leftarrow \text{normal velocity}$$

$$P_{HPE} = \sum_{ik} m_{ik} \left[\overline{\Phi}_i^k + e \left(\alpha_{ik} = \frac{A_i \delta_k (a^2 \Phi_i / g)}{m_{ik}}, \theta_{ik} = \frac{\Theta_{ik}}{m_{ik}} \right) \right]$$

TRISK

shallow-atmosphere metric
planetary velocity



$$\partial_t m_{ik} + \delta_i \frac{\partial H}{\partial v_{ek}} = 0 \quad \partial_t \Theta_{ik} + \delta_i \left(\theta_{ek}^* \frac{\partial H}{\partial v_{ek}} \right) = 0$$

$$\partial_t v_{ek} + \left(\frac{\delta_v v_k}{\overline{m}_k} \frac{\partial H}{\partial v_{ek}} \right)^\perp + \delta_e \frac{\partial H}{\partial m_{ik}} + \theta_{ek}^* \delta_e \frac{\partial H}{\partial \Theta_{ik}} = 0$$

SW pot. Vort.

Bernoulli function

Exner function

Discrete energy budget : Lagrangian vertical coordinate

$$\begin{array}{ccc}
 \text{mass flux} & \text{centered/upwind} & \\
 \partial_t m_{ik} + \delta_i \frac{\partial H}{\partial v_{ek}} = 0 & \partial_t \Theta_{ik} + \delta_i \left(\theta_{ek}^* \frac{\partial H}{\partial v_{ek}} \right) = 0 & \\
 \\
 \partial_t v_{ek} + \left(\frac{\delta_v v_k}{\overline{m_k^v}} \frac{\partial H}{\partial v_{ek}} \right)^\perp + \delta_e \frac{\partial H}{\partial m_{ik}} + \theta_{ek}^* \delta_e \frac{\partial H}{\partial \Theta_{ik}} = 0 & & \\
 \text{SW pot. Vort.} & \text{Bernoulli function} & \text{Exner function}
 \end{array}$$

discrete div
 $\updownarrow$ discrete integration by parts
 discrete grad

$$\frac{dH}{dt} = \sum \left(\frac{\partial H}{\partial m_{ik}} \partial_t m_{ik} + \frac{\partial H}{\partial \Theta_{ik}} \partial_t \Theta_{ik} + \frac{\partial H}{\partial v_{ek}} \partial_t v_{ek} + \frac{\partial H}{\partial \Phi_{il}} \partial_t \Phi_{il} \right) = 0$$

$\uparrow$
 Hydrostatic balance

$$\Rightarrow p_{ik} \Rightarrow \alpha_{ik} \Rightarrow \Phi_{il}$$

Non-Lagrangian vertical coordinate : also possible to cancel additional contributions from vertical transport (Tort et al., QJRMS 2015)

What is new in NH (fully compressible) compared to hydrostatic ?

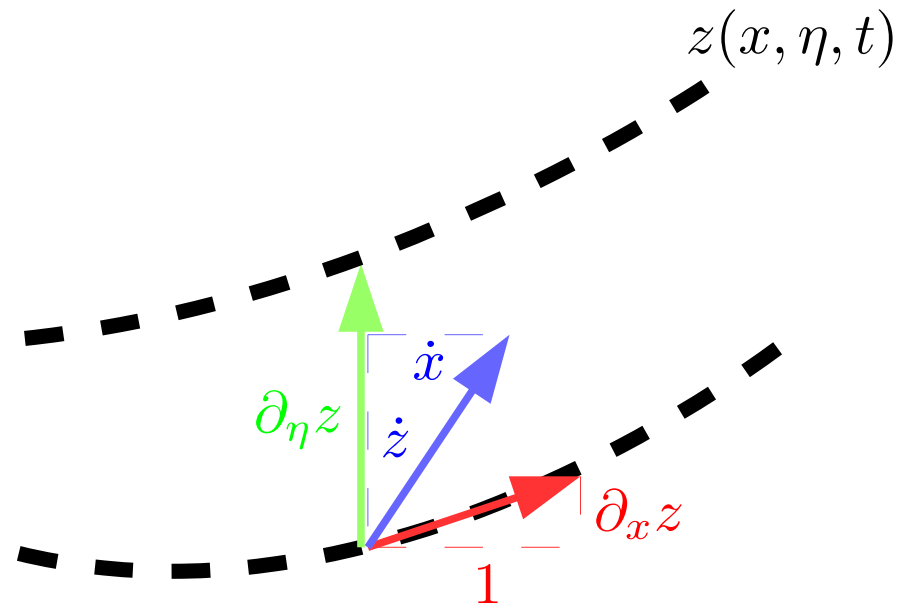
- More *energy* : vertical kinetic energy K_v not neglected any more
- More *prognostic* variables : geopotential, vertical momentum

Example : Cartesian (x,z) slice, non-rotating

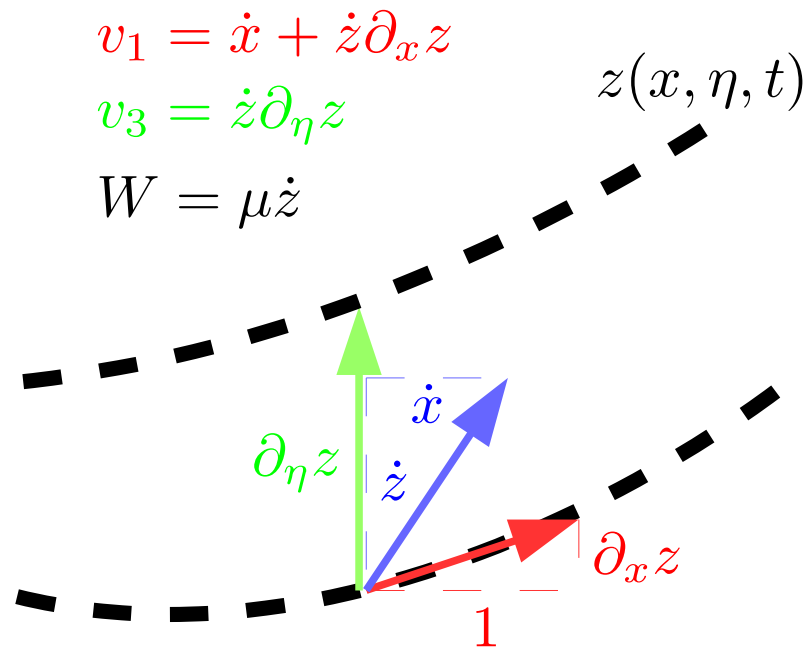
$$v_1 = \dot{x} + \dot{z} \partial_x z$$

$$v_3 = \dot{z} \partial_\eta z$$

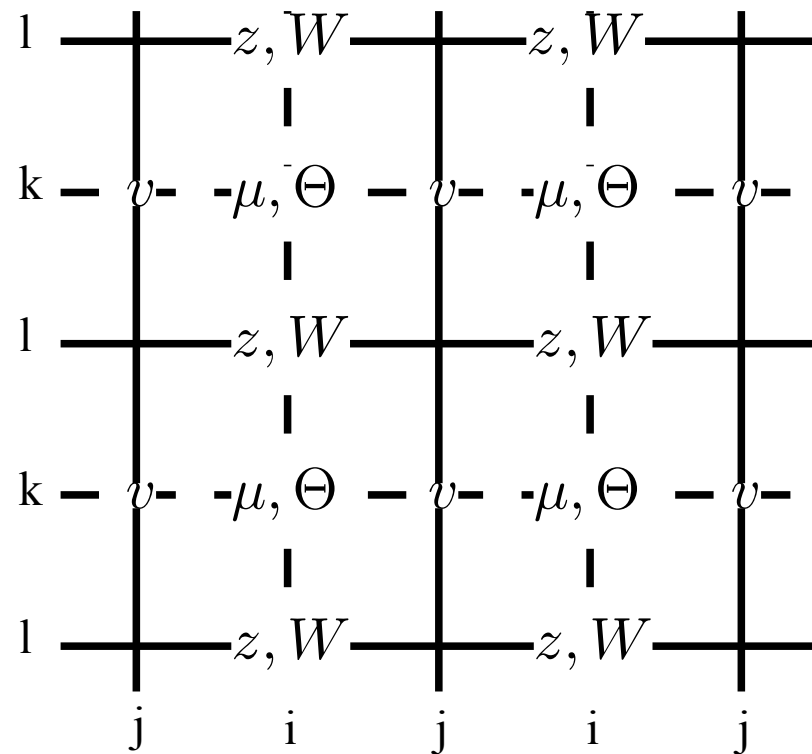
$$W = \mu \dot{z}$$



$$\mathcal{H} = \int \left[\frac{\left(v_1 - \frac{W}{\mu} \partial_x z \right)^2 + \left(\frac{W}{\mu} \right)^2}{2} + e \left(\frac{\partial_\eta z}{\mu}, \frac{\Theta}{\mu} \right) + gz \right] \mu dx d\eta$$



C-grid Lorentz staggering



$$K_H = \frac{1}{2} \sum \left(v_{jk}^2 \overline{\mu_i^j} - 2 \overline{W_l^j} \overline{v_j^l} D_j z_l + \frac{(\overline{D_j z})^2 \overline{W_{il}^2}}{\overline{\mu_i^l}} \right) \simeq \frac{1}{2} \int \left(v - \partial_x z \frac{W}{\mu} \right)^2 \mu dx d\eta$$

$$\frac{\partial K_H}{\partial v_{jk}} = \overline{\mu_i^j} v_{jk} - \overline{D_j z_l \overline{W_l^j}}$$

« horizontal » mass flux

$$\frac{\partial K_H}{\partial z_{il}} = D_i \left(\overline{W_l^j} \overline{v_j^l} - D_j z_l \frac{\overline{W_{il}^2}}{\overline{\mu_i^l}} \right)$$

« horizontal » advection of vertical momentum

slow

fast

$$H = K_H[\mu, \Phi, \mathbf{v}, W] + K_V[\mu, \Phi, W] + P[\mu, \Theta, \Phi]$$

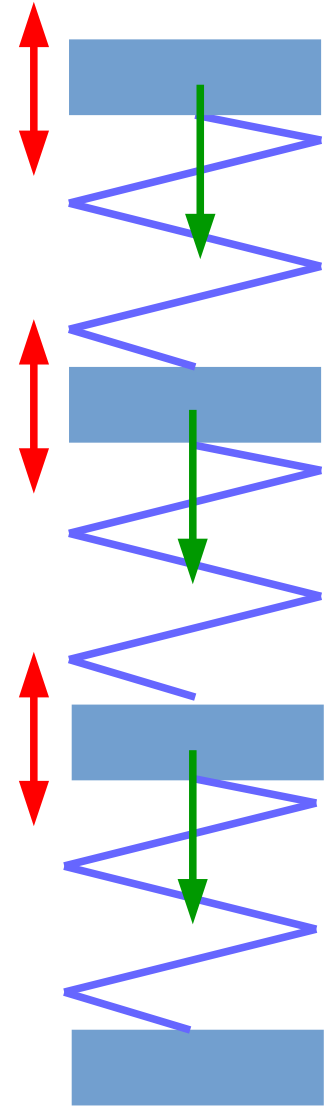
$$\partial_t \mu + \partial_\eta (\mu \dot{\eta}) + \partial_{\mathbf{x}} \cdot \frac{\delta H}{\delta \mathbf{v}} = 0$$

$$\partial_t \Theta + \partial_\eta (\theta \mu \dot{\eta}) + \partial_{\mathbf{x}} \cdot \theta \frac{\delta H}{\delta \mathbf{v}} = 0$$

$$\begin{aligned} \partial_t \mathbf{v} + \dot{\eta} \partial_\eta \mathbf{v} + v_3 \partial_{\mathbf{x}} \dot{\eta} + \frac{\partial_{\mathbf{x}} \times \mathbf{v}}{\mu} \times \frac{\delta H}{\delta \mathbf{v}} + \partial_{\mathbf{x}} \frac{\delta K_H}{\delta \mu} \\ + \partial_{\mathbf{x}} \frac{\delta(K_V + P)}{\delta \mu} + \theta \partial_{\mathbf{x}} \frac{\delta P}{\delta \Theta} = 0 \end{aligned}$$

$$\partial_t \Phi + \dot{\eta} \partial_\eta \Phi - \frac{\delta K_H}{\delta W} - \frac{\delta K_V}{\delta W} = 0$$

$$\partial_t W + \partial_\eta (\dot{\eta} W) + \frac{\delta K_H}{\delta \Phi} + \frac{\delta(K_V + P)}{\delta \Phi} = 0$$



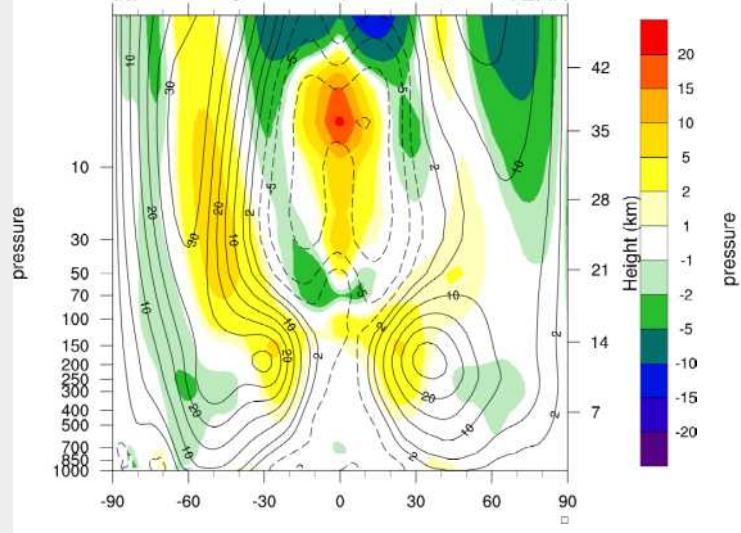
- The implicit problem only couples vertical position and vertical momentum
- eliminate W and obtain a scalar tridiagonal implicit problem for Φ

Zonal wind

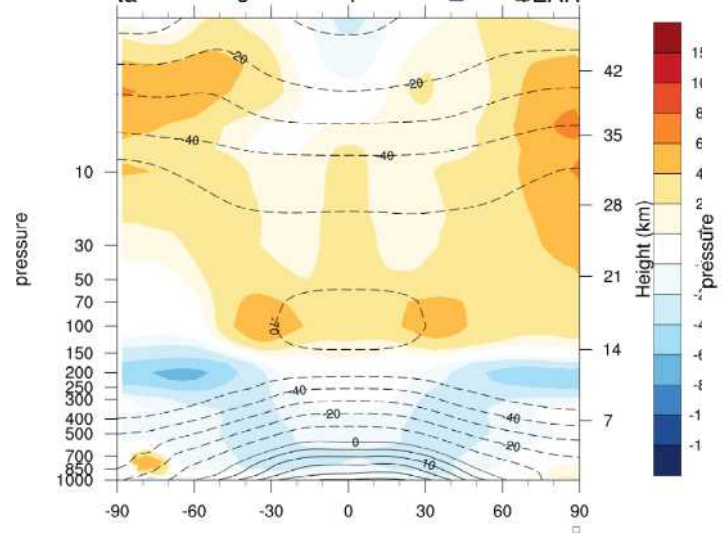
Temperature

Relative humidity

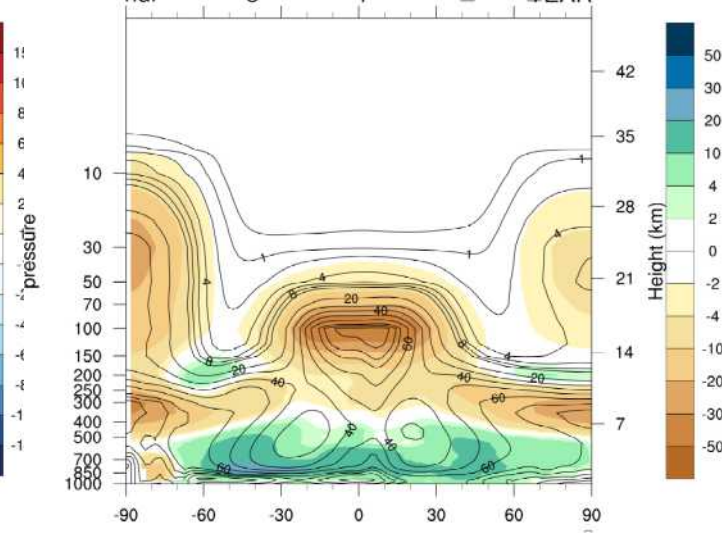
CM6.1-ATM-LR-highresSST-present.2_1990-1999 - OBS



CM6.1-ATM-LR-highresSST-present.2_1990-1999 - OBS



CM6.1-ATM-LR-highresSST-present.2_1990-1999 - OBS

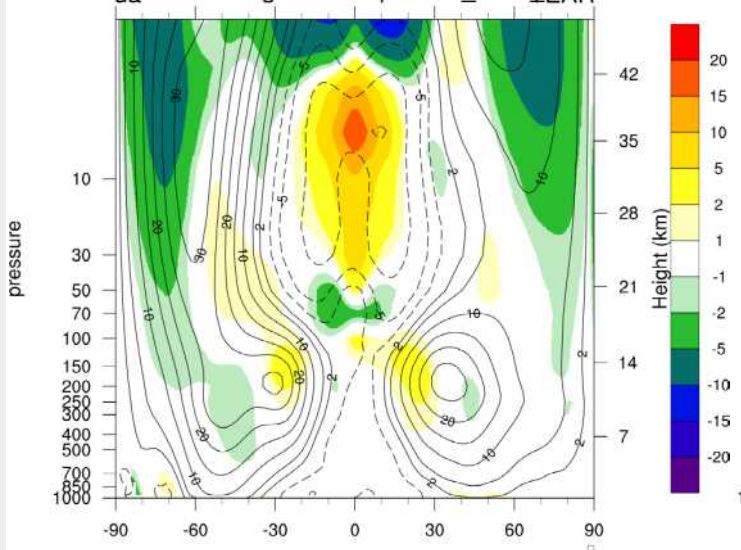


IPSL-CM6.1-LR

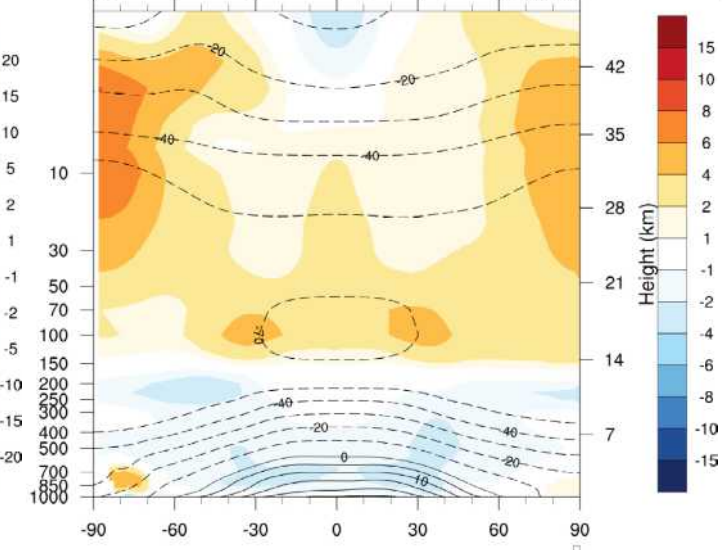
Contours : OBS (ERA)

Colors : bias = model-OBS

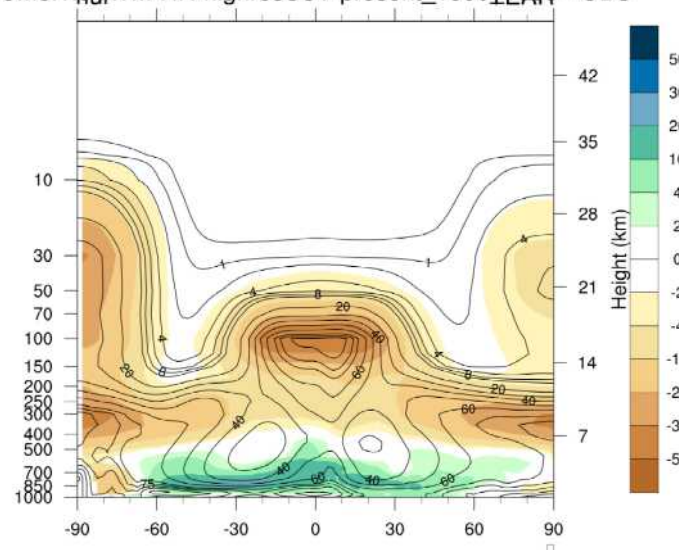
CM6.1.5-ATM-HR-highresSST-present_1990-1999 - OBS



CM6.1.5-ATM-HR-highresSST-present_1990-1999 - OBS



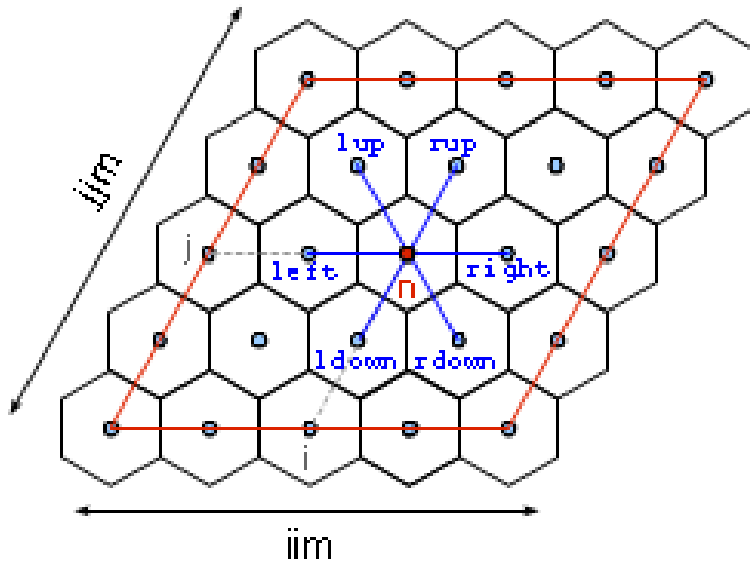
CM6.1.5-ATM-HR-highresSST-present_1990-1999 - OBS



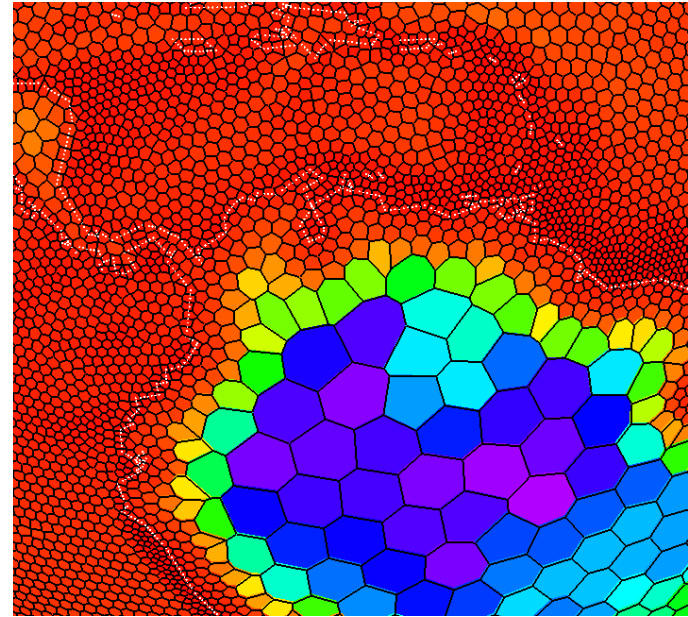
IPSL-CM6.1-HR

=> some reduction of mean biases at higher resolution

Unstructured-mesh capability



Structured mesh



Unstructured mesh

- Quasi-uniform resolution
- Zoom possible but limited
- Regular data access / compute pattern

- Variable resolution
- Very flexible zoom capability
- Irregular data access / compute pattern
- Needs scale-aware physics

Status :

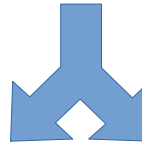
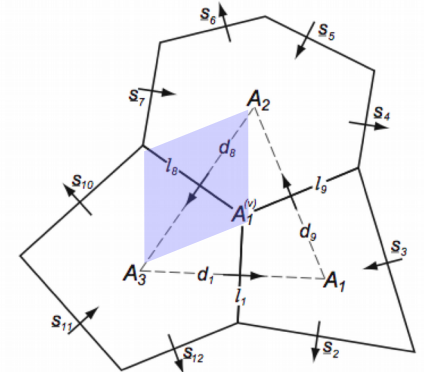
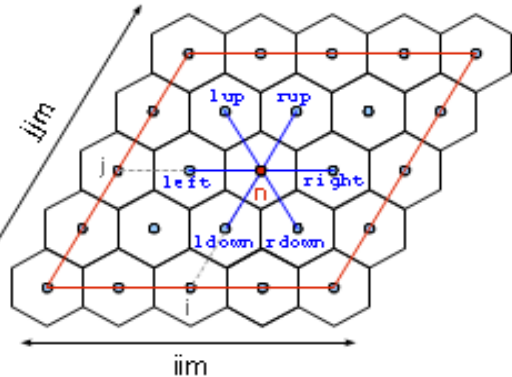
- Prototype available in the « devel » branch
- MPI/OpenMP parallelism
- More work ahead (performance, transport scheme, ...)

DySL : Dynamico-specific « language »

```

49  KERNEL('div')
50    FORALL_CELLS_EXT()
51      ON_PRIMAL
52        div_ij=0.
53        FORALL_EDGES
54          div_ij = div_ij + SIGN*LE_DE*u(EDGE)
55        END_BLOCK
56        divu(CELL) = div_ij / AI
57      END_BLOCK
58    END_BLOCK
59  END_BLOCK

```



```

1  DO l = ll_begin, ll_end
2  !DIR$ SIMD
3    DO ij=ij_begin_ext, ij_end_ext
4      div_ij=0.
5      div_ij = div_ij + ne_rup*le_de(ij+u_rup)*u(ij+u_rup,l)
6      div_ij = div_ij + ne_lup*le_de(ij+u_lup)*u(ij+u_lup,l)
7      div_ij = div_ij + ne_left*le_de(ij+u_left)*u(ij+u_left,l)
8      div_ij = div_ij + ne_ldown*le_de(ij+u_ldown)*u(ij+u_ldown,l)
9      div_ij = div_ij + ne_rdown*le_de(ij+u_rdown)*u(ij+u_rdown,l)
10     div_ij = div_ij + ne_right*le_de(ij+u_right)*u(ij+u_right,l)
11     divu(ij,l) = div_ij / Ai(ij)
12   END DO
13 END DO
14

```

```

2  DO ij = 1, primal_num
3    DO l = 1, llm
4      div_ij=0.
5      DO iedge = 1, primal_deg(ij)
6        edge = primal_edge(iedge,ij)
7        div_ij = div_ij + primal_ne(iedge,ij)*le_de(edge)*u(l,edge)
8      END DO
9      divu(l,ij) = div_ij / Ai(ij)
10    END DO
11  END DO

```

- Code generator purely text-based (no grammar / parser / syntax tree)
- Macro substitution (cpp) and inlining (jinja2)
- Generates *human-readable* code chunks to be #included in hand-written Fortran code
- Lightweight, low exit cost