

# Planetary Boundary Layer 3

Outer layer and model sensitivity

Annelize van Niekerk, Irina Sandu, Anton Beljaars

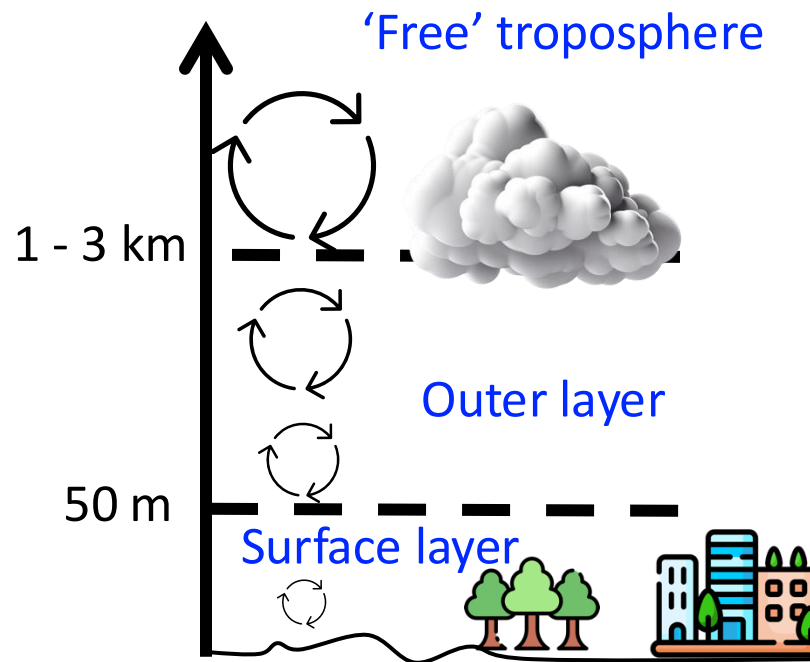
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# Contents

- K-closure in the outer layer
- Eddy-diffusivity Mass-Flux (non-local transport)
- Description of IFS scheme
- Sensitivity of forecasts to changes in diffusion scheme:
  - Model issues in stable boundary layers
  - Using prognostic TKE

# What do we need from a turbulence parametrization scheme?

- Provide turbulent fluxes of heat, momentum, moisture (and tracers) between the surface and the upper atmosphere
- Account for differences in stability and surface properties
- Provide profiles of winds and temperatures at the surface, where the model does not resolve in the vertical
- Provide turbulent mixing throughout the entire atmosphere – the mixed layer, the cloud layer and the stratosphere



Model level heights:

\_\_\_\_\_ 10 km

\_\_\_\_\_

\_\_\_\_\_

\_\_\_\_\_ 3 km

\_\_\_\_\_

\_\_\_\_\_ 500 m

\_\_\_\_\_ 40 m

\_\_\_\_\_ 20 m

\_\_\_\_\_ 10 m

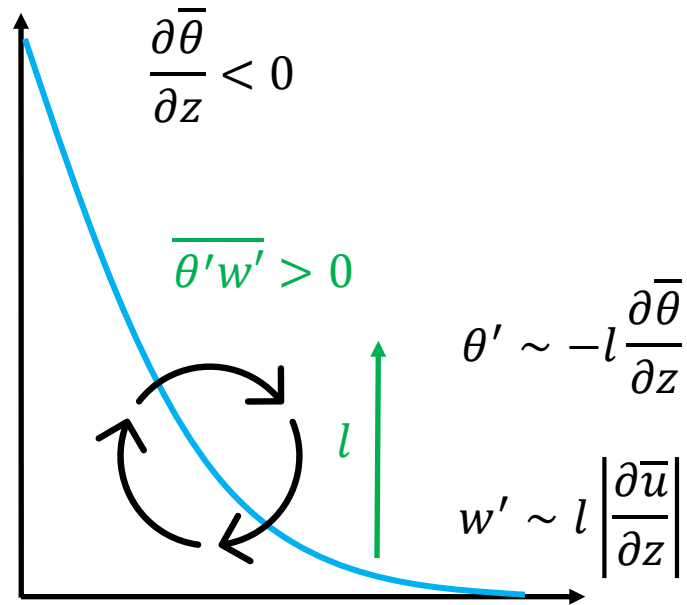
\_\_\_\_\_  $z_0$

‘Local’ turbulence closure: eddy  
diffusion above the surface

# 'Local' turbulence closure: eddy diffusion above the surface

Any quantity  $\phi$ :

$$\overline{\phi'w'} = -K_\phi \frac{\partial \bar{\phi}}{\partial z} = -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{\phi}}{\partial z}$$



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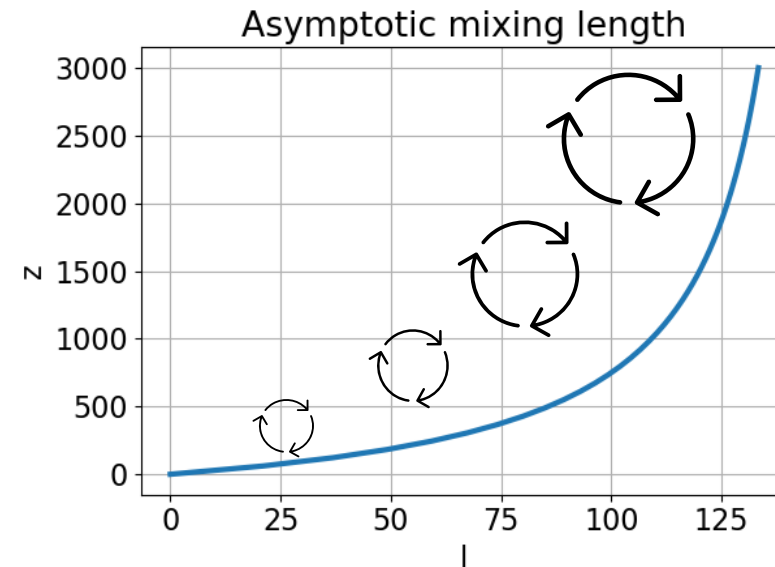
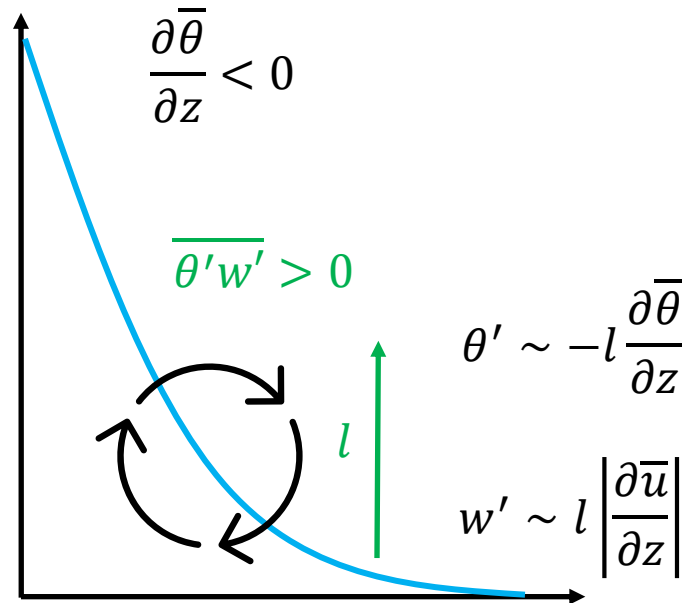
$$\overline{\phi'w'} = -K_\phi \frac{\partial \bar{\phi}}{\partial z} = -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{\phi}}{\partial z}$$

Size of eddies get larger further away from the surface:

$$l \sim \frac{\kappa z \lambda}{\kappa z + \lambda}$$

$\kappa$  = von-Karman constant

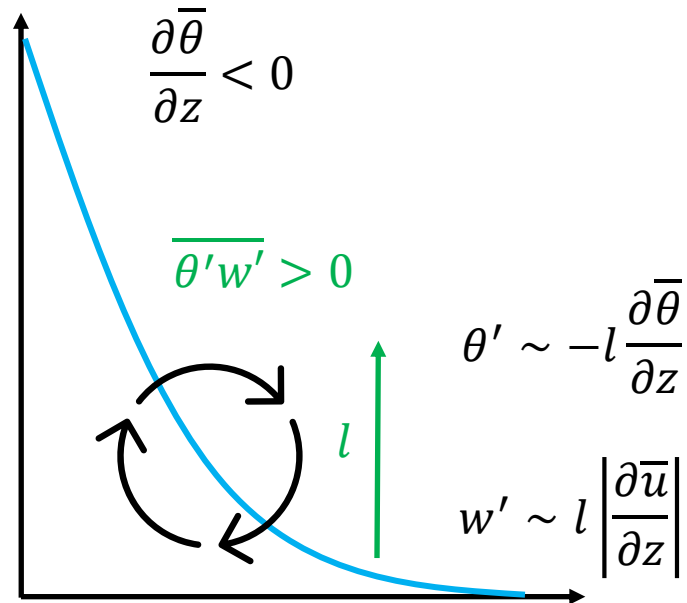
$\lambda$  = asymptotic mixing length (150 m)



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Add stability dependence:

$f_M(Ri)$ ,  $f_H(Ri)$  determined empirically and depend on  $Ri(z)$

No longer using  $\frac{z}{L}$ , since we are away from the surface

# Local similarity theory in the outer layer – making use of the surface stability functions

Any quantity  $\phi$ :

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- In stable conditions, the mid and upper boundary layer may not be in equilibrium with the surface fluxes
- Local fluxes and stability ( $Ri$ ) dominate



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$$K_\phi = \frac{l^2}{\Phi_\phi(\zeta)\Phi_M(\zeta)} \left| \frac{\partial \bar{u}}{\partial z} \right| = l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| f_\phi(Ri)$$

Use the relation

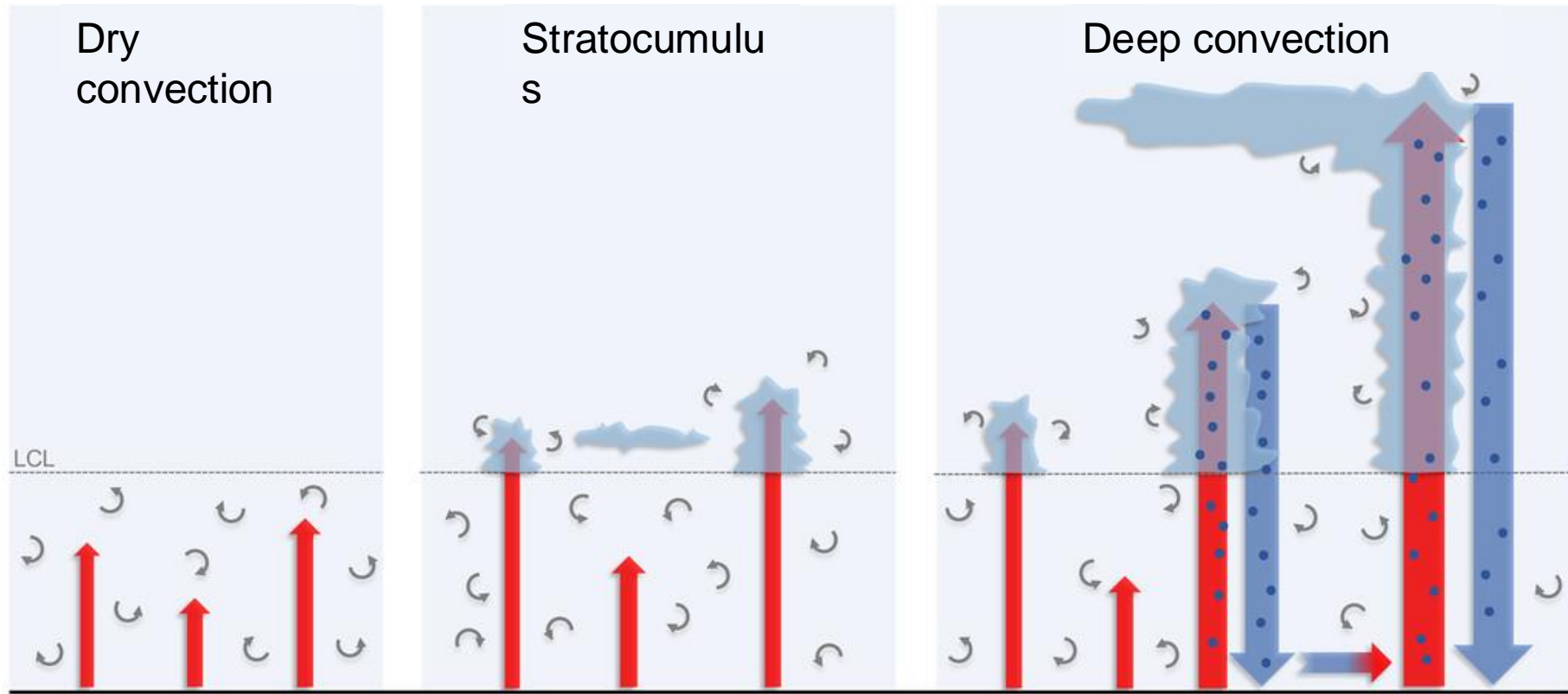
$$Ri = \zeta \frac{\Phi_H(\zeta)}{\Phi_M^2(\zeta)}$$

to convert  $\zeta = \frac{z}{L}$  to the gradient Richardson number in the outer layer

‘Non-Local’ turbulence: eddy-diffusivity mass-flux (EDMF)

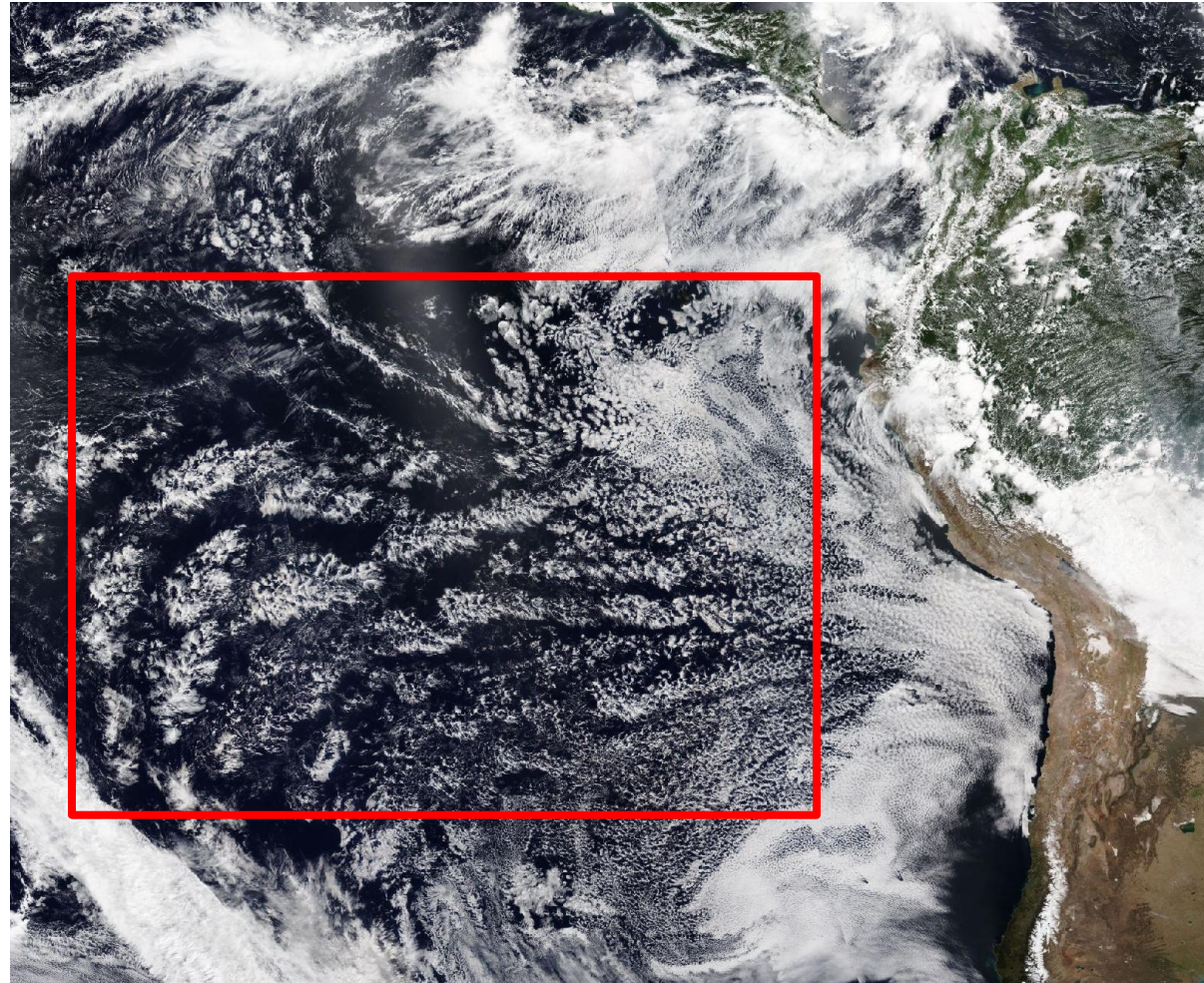
# ‘Non-Local’ turbulence: eddy-diffusivity mass-flux (EDMF)

Local turbulent diffusion fails in convective boundary layers because it yields unrealistic zero flux in an environment with small gradients



# ‘Non-Local’ turbulence: eddy-diffusivity mass-flux (EDMF)

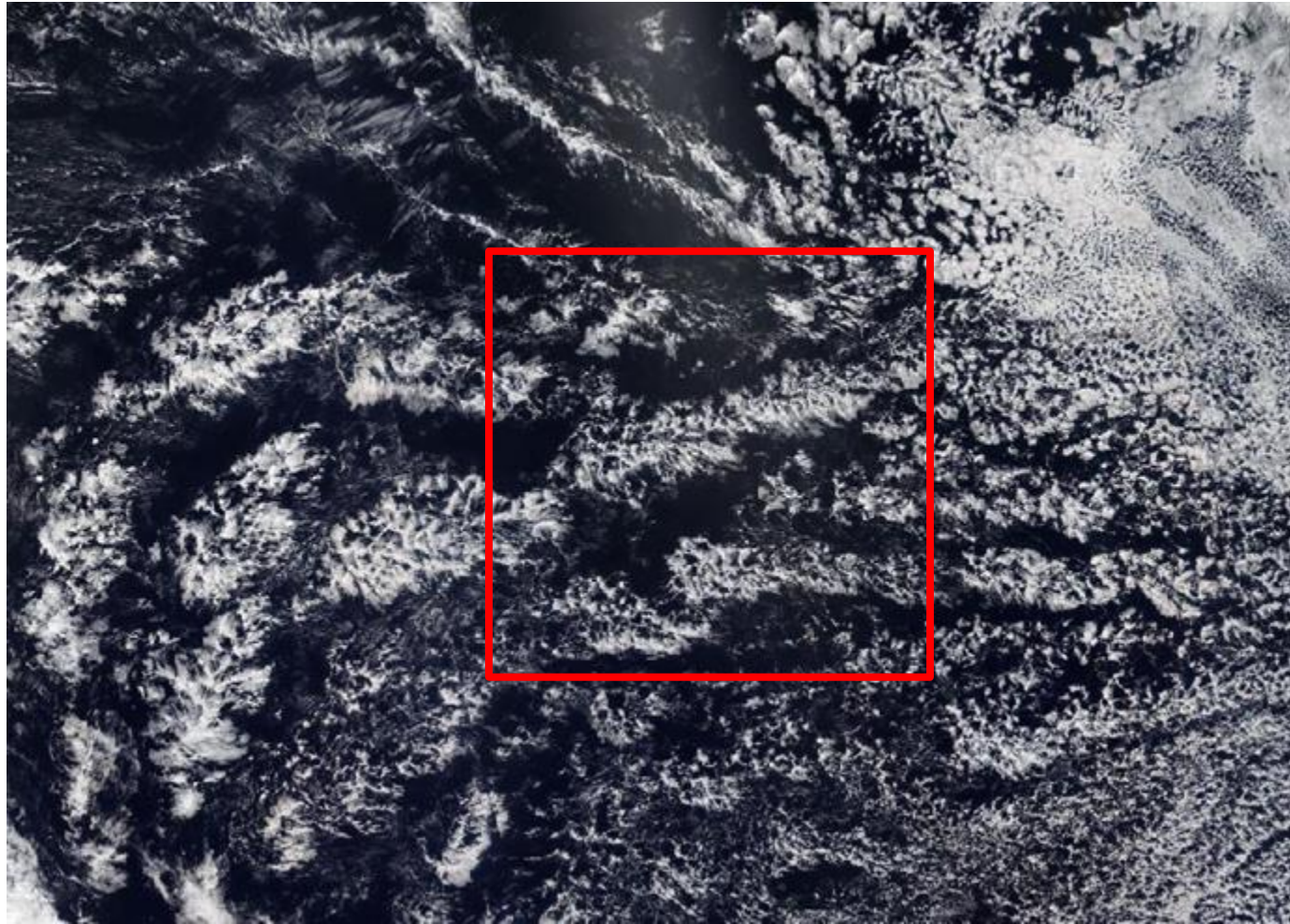
In convective boundary layers, the area of strongest updraft is typically much smaller than the environment





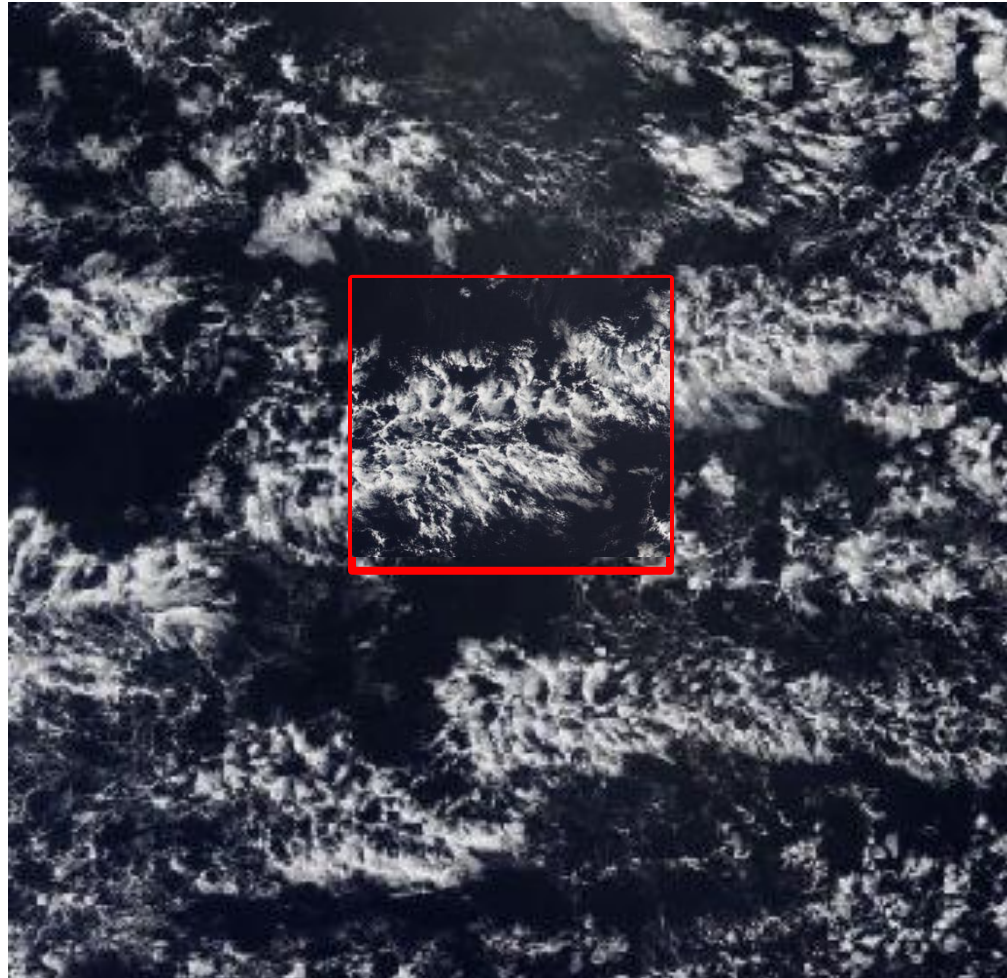
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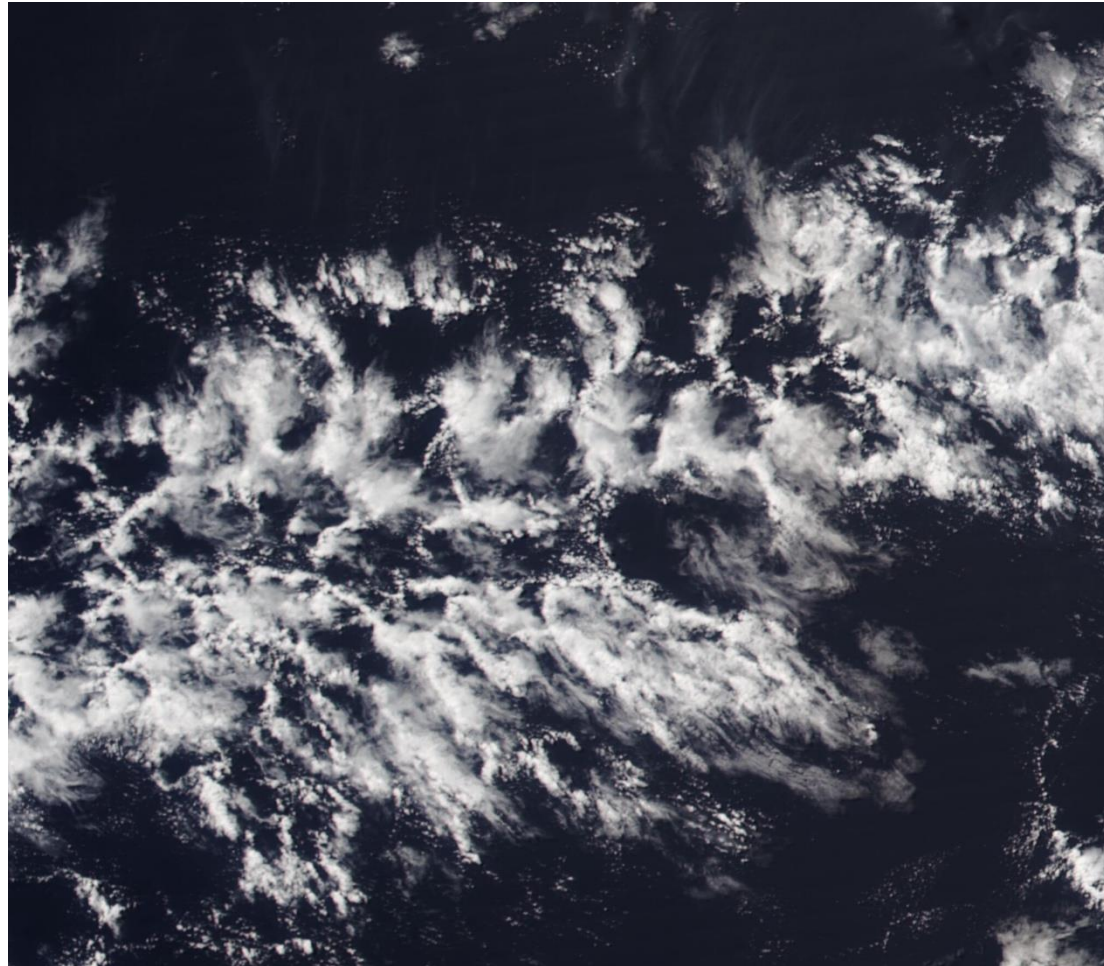
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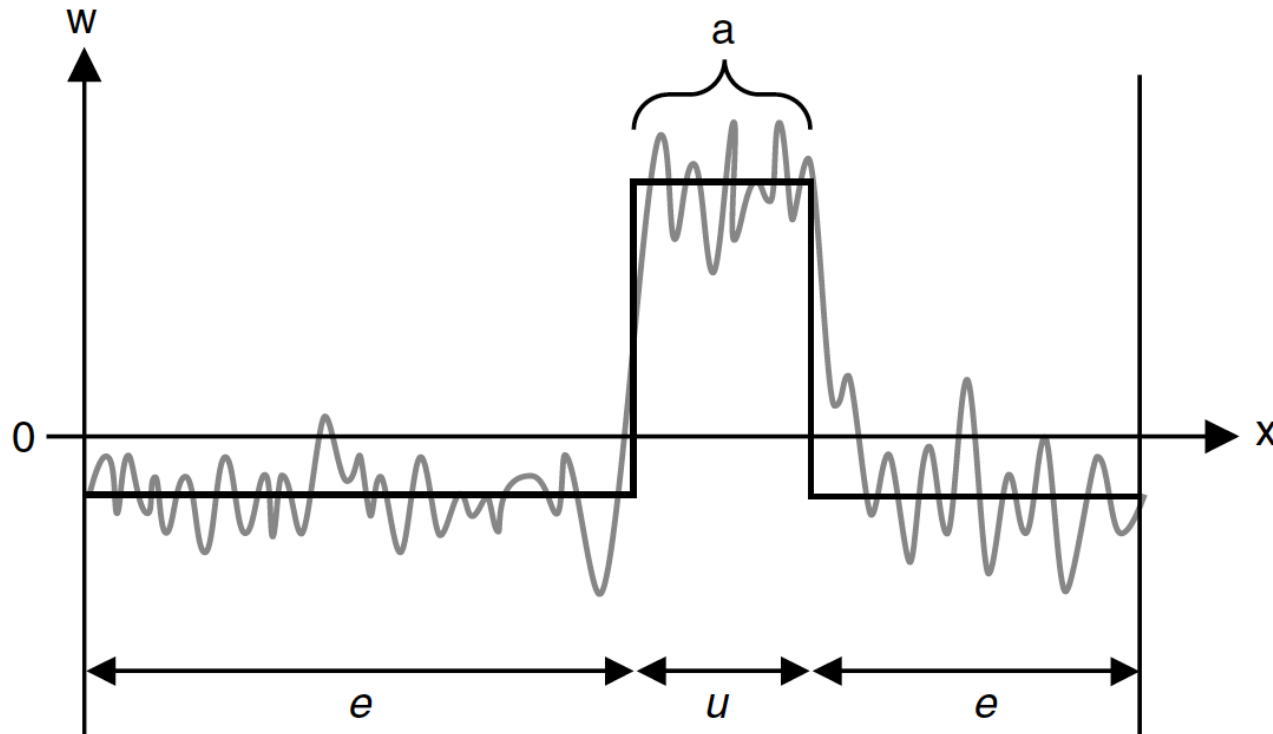
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Total turbulent flux of  $\phi$  :

$$\overline{\phi'w'} = a\overline{\phi'_uw'} + (1-a)\overline{\phi'_ew'} + a(\overline{w^u} - \overline{w})(\overline{\phi^u} - \overline{\phi^e})$$



Turbulent flux within  
the strong updraft  
region

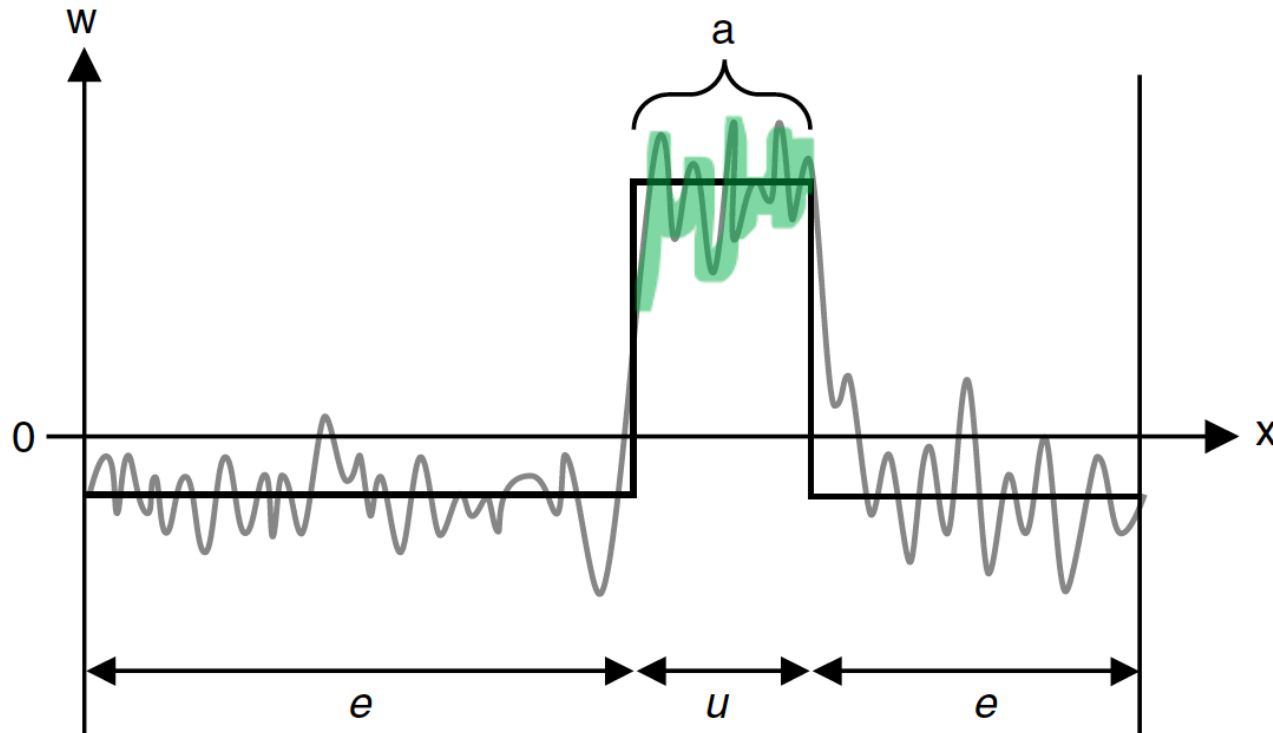


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Subcore  
flux



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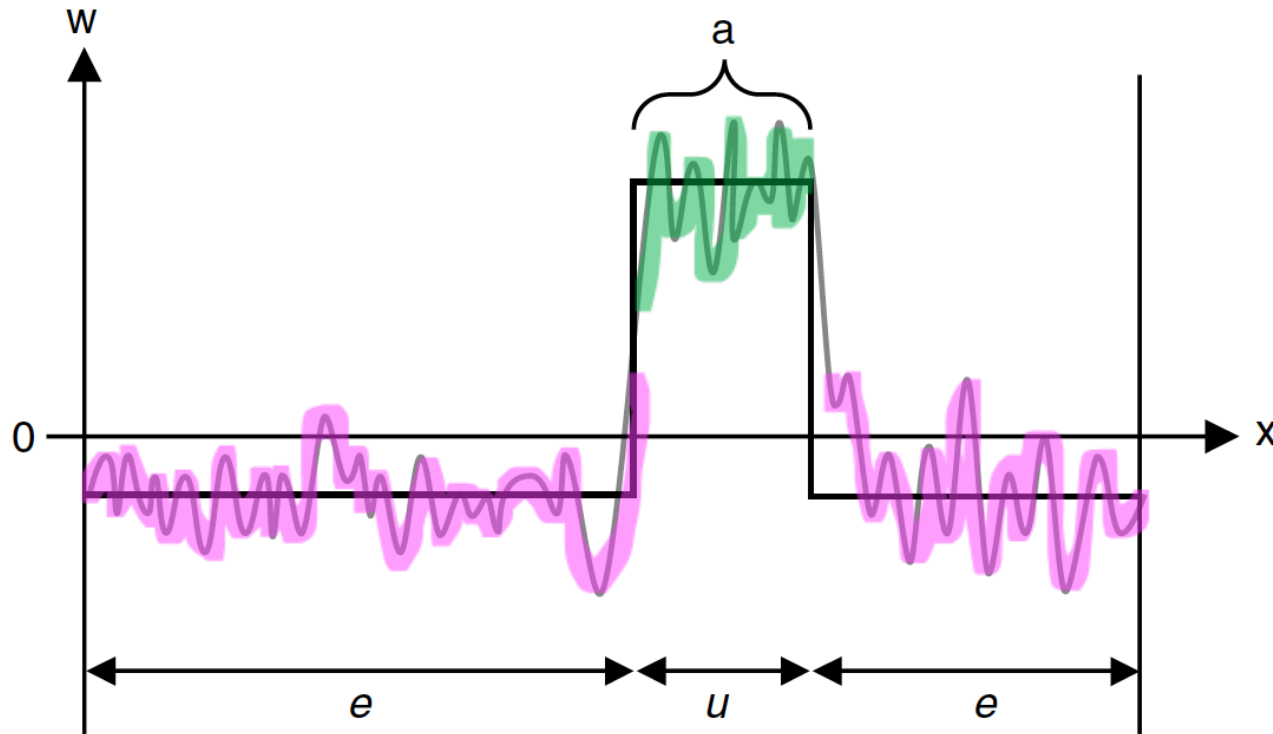
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Subcore  
flux

Environmental  
flux



Turbulent flux in the environment outside the strongest updraft

# 'Non-Local' turbulence: eddy-diffusivity mass-flux (EDMF)

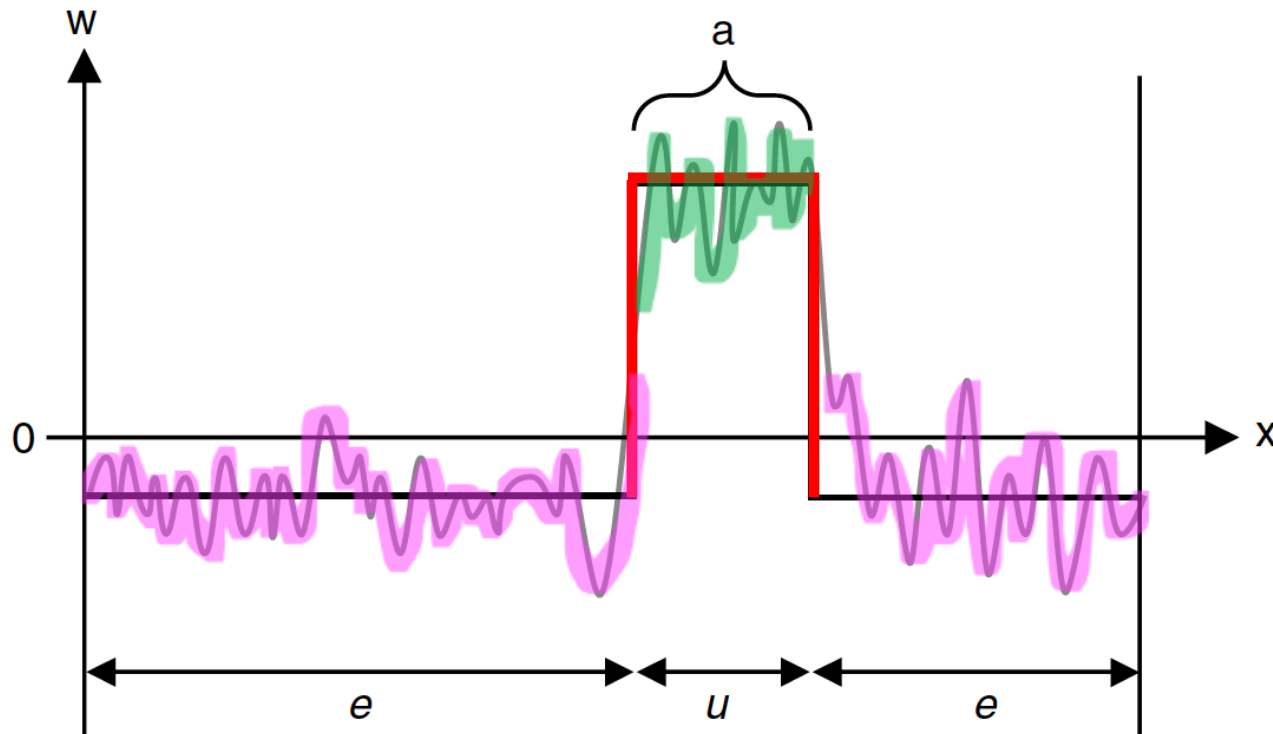
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Subcore  
flux

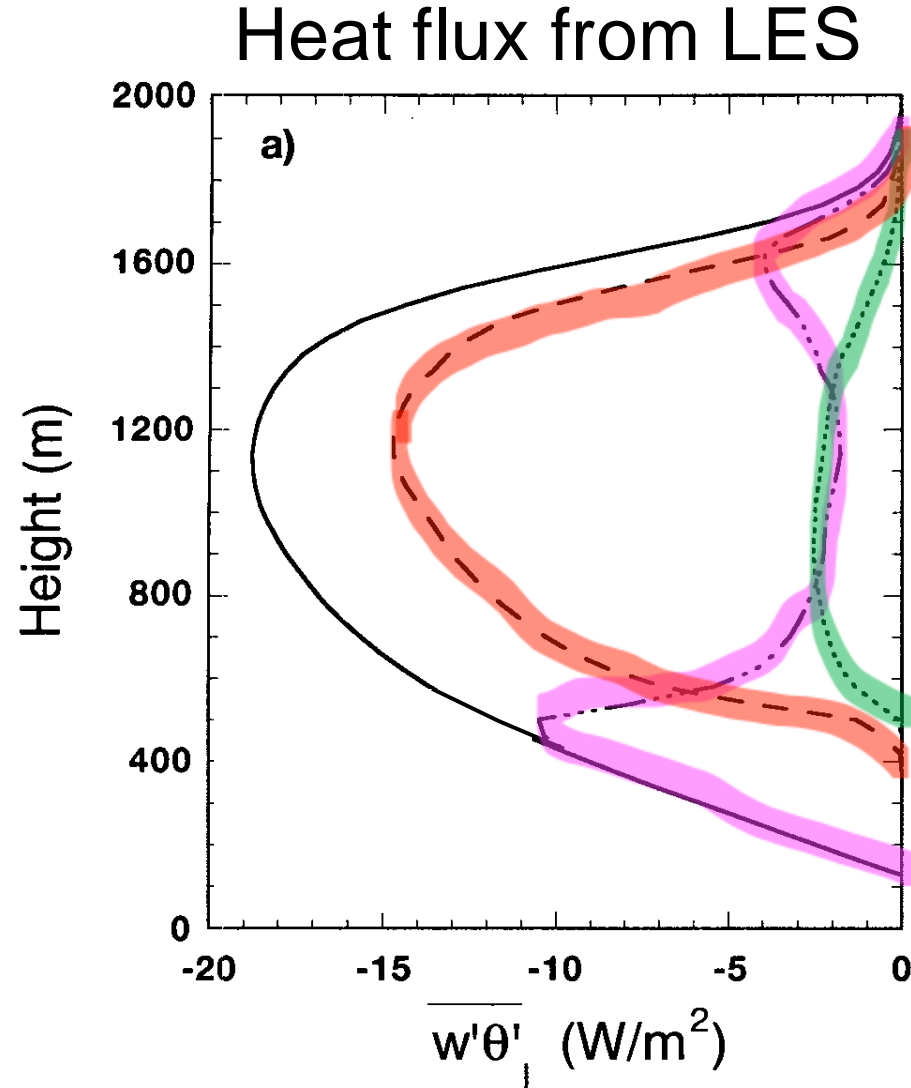
Environmental  
flux

Mass  
flux



Mean flux inside the  
strongest updraft  
region

# 'Non-Local' turbulence: eddy-diffusivity mass-flux (EDMF)



Total turbulent flux of  $\phi$  :

$$\overline{\phi'w'} = \overline{a\phi'_u w'} +$$

Subcore  
flux

$$(1 - a)\overline{\phi'_e w'} +$$

Environmental  
flux

$$a(\overline{w}^u - \overline{w})(\overline{\phi}^u - \overline{\phi}^e)$$

Mass  
flux

M-flux covers 80% of the flux for  
heat and moisture

# 'Non-Local' turbulence: eddy-diffusivity mass-flux (EDMF)

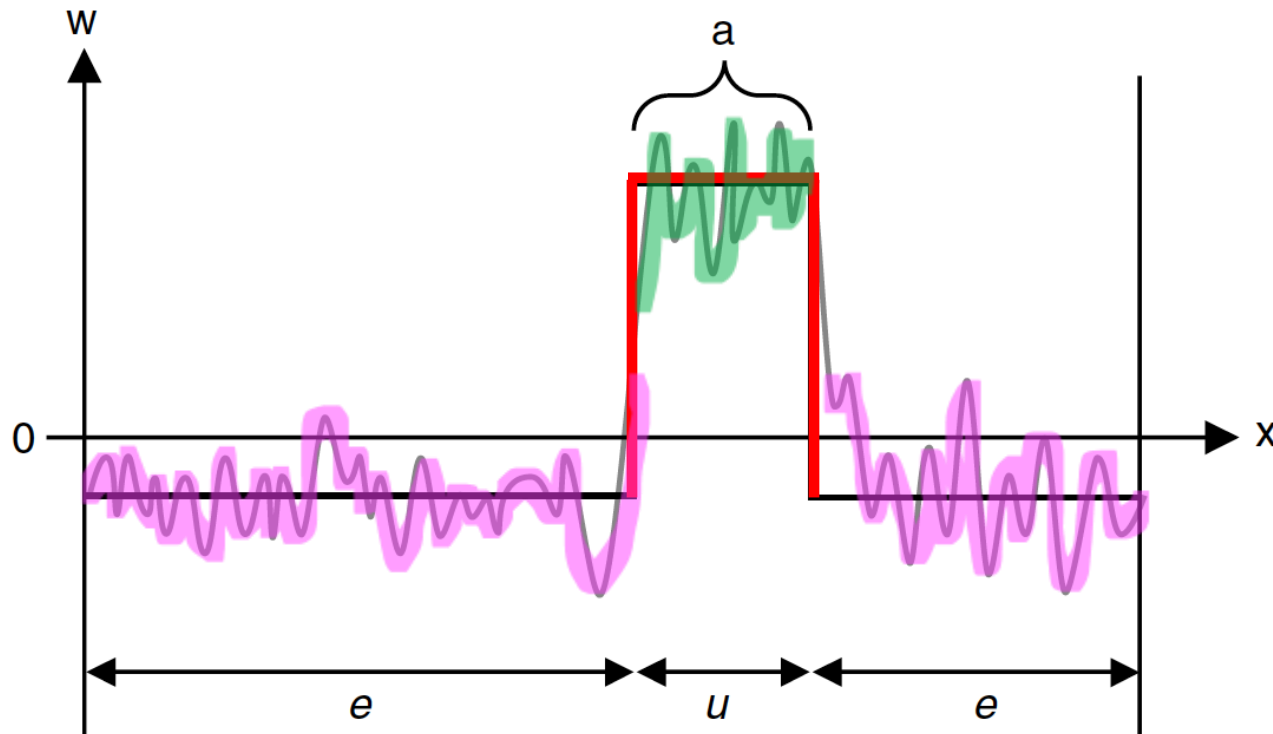
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Subcore  
flux

Environmental  
flux

Mass  
flux



**Assumptions made:**

1. Area of strongest updraft is small compared with the environment ( $a \ll 1$ ). Subcore flux is neglected

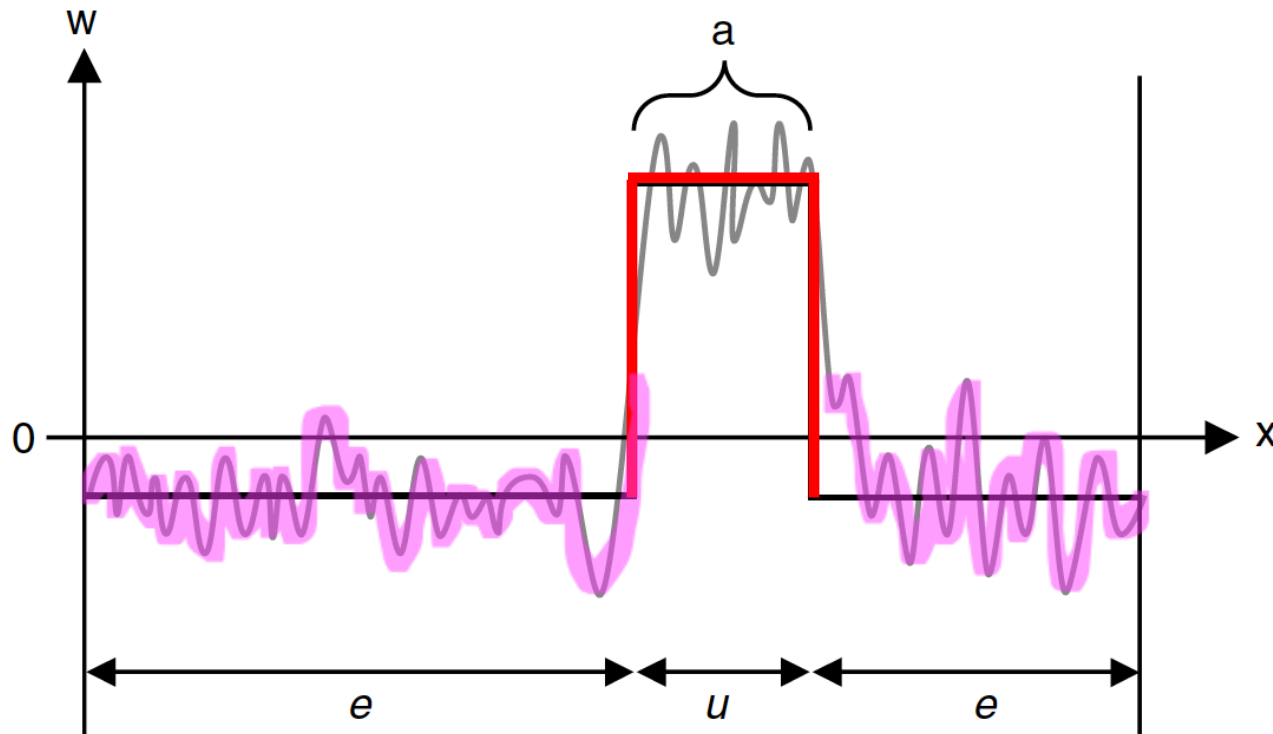
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flux

Mass  
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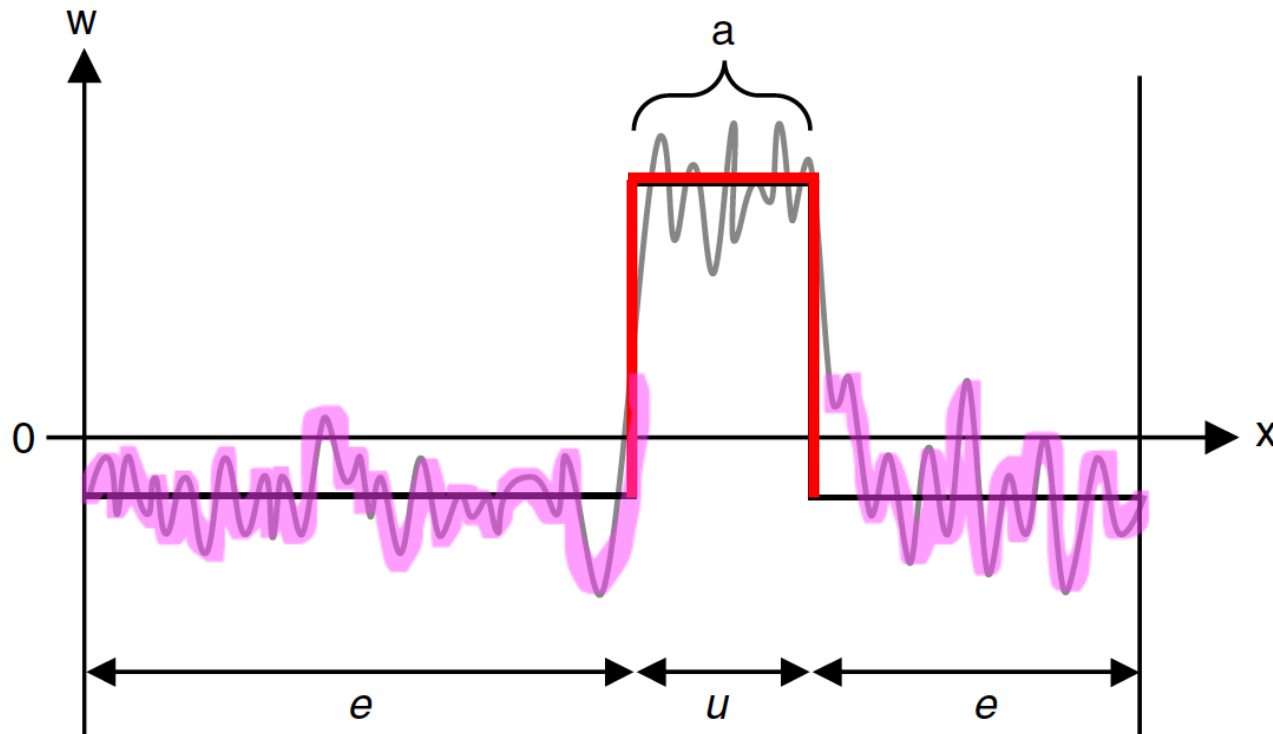
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Mass  
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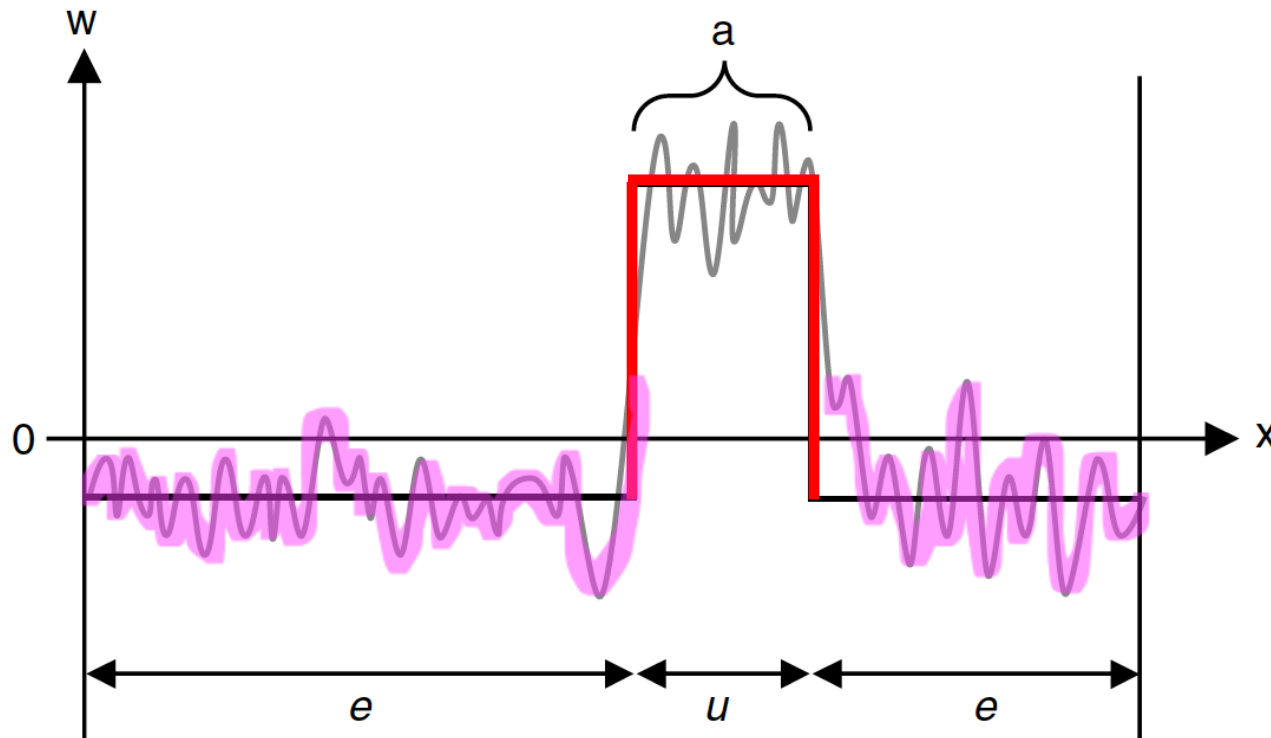
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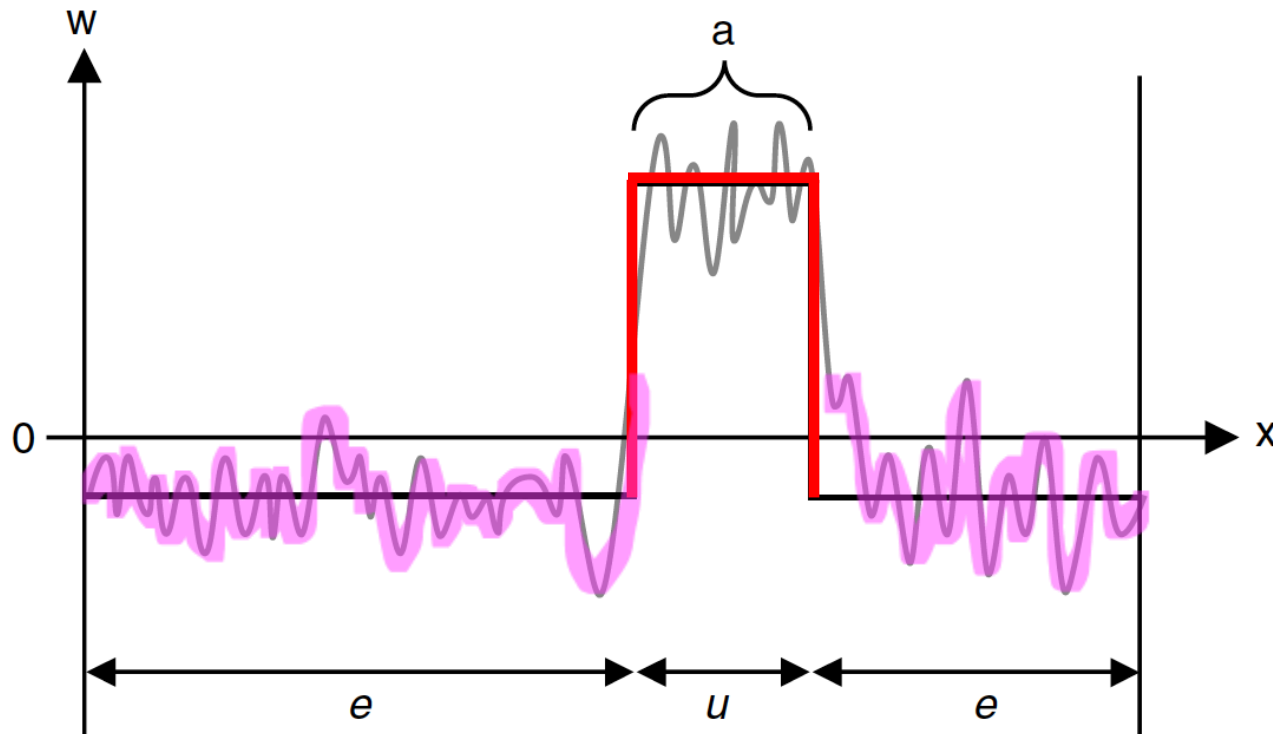
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Mass  
flux

$$M = a(\overline{w}^u - \overline{w})$$



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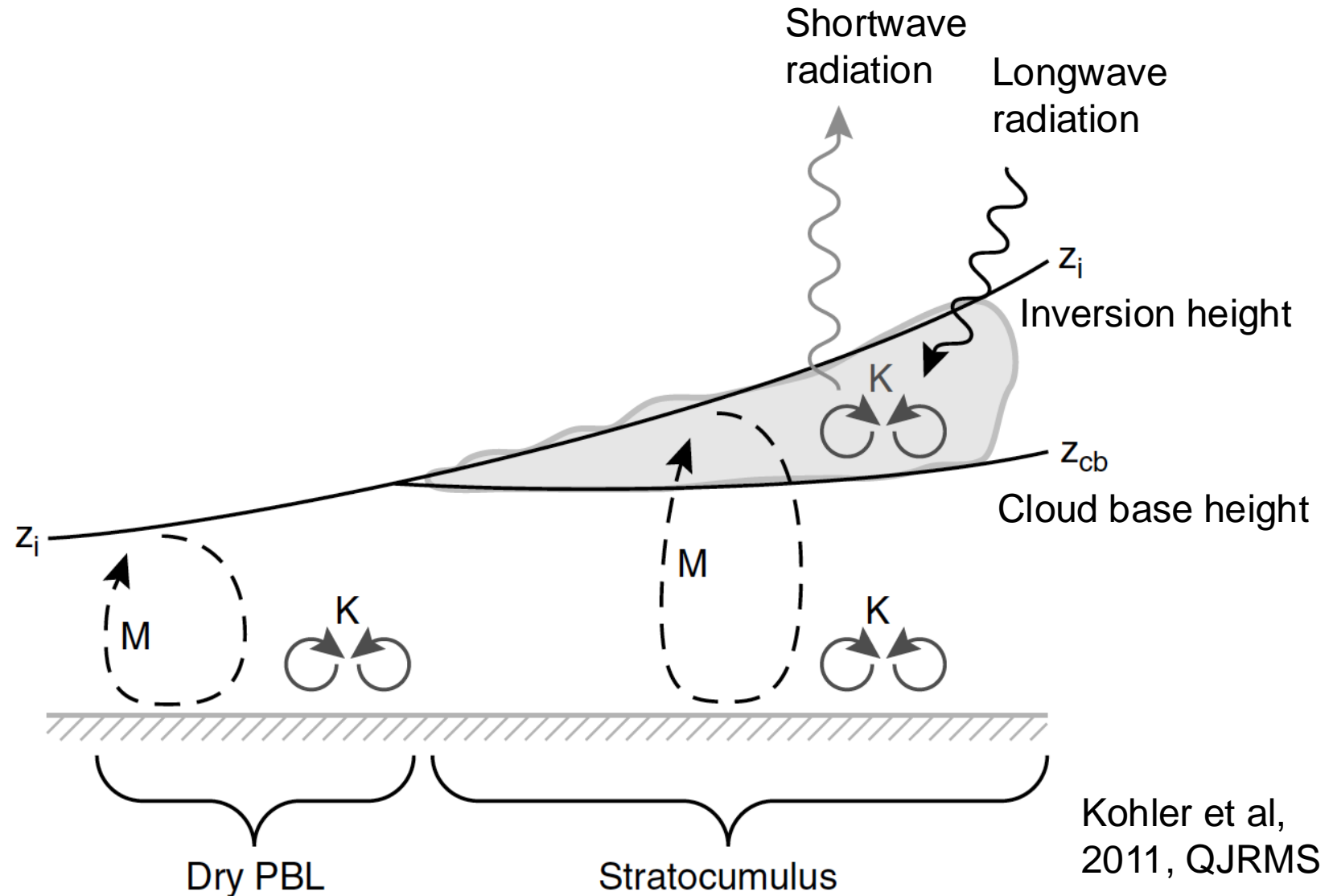
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The surface mass flux ( $M$ ) is initialised at the first model level

The mass flux profile then depends on the inversion height ( $z_i$ ) or the cloud base height ( $z_{cb}$ )

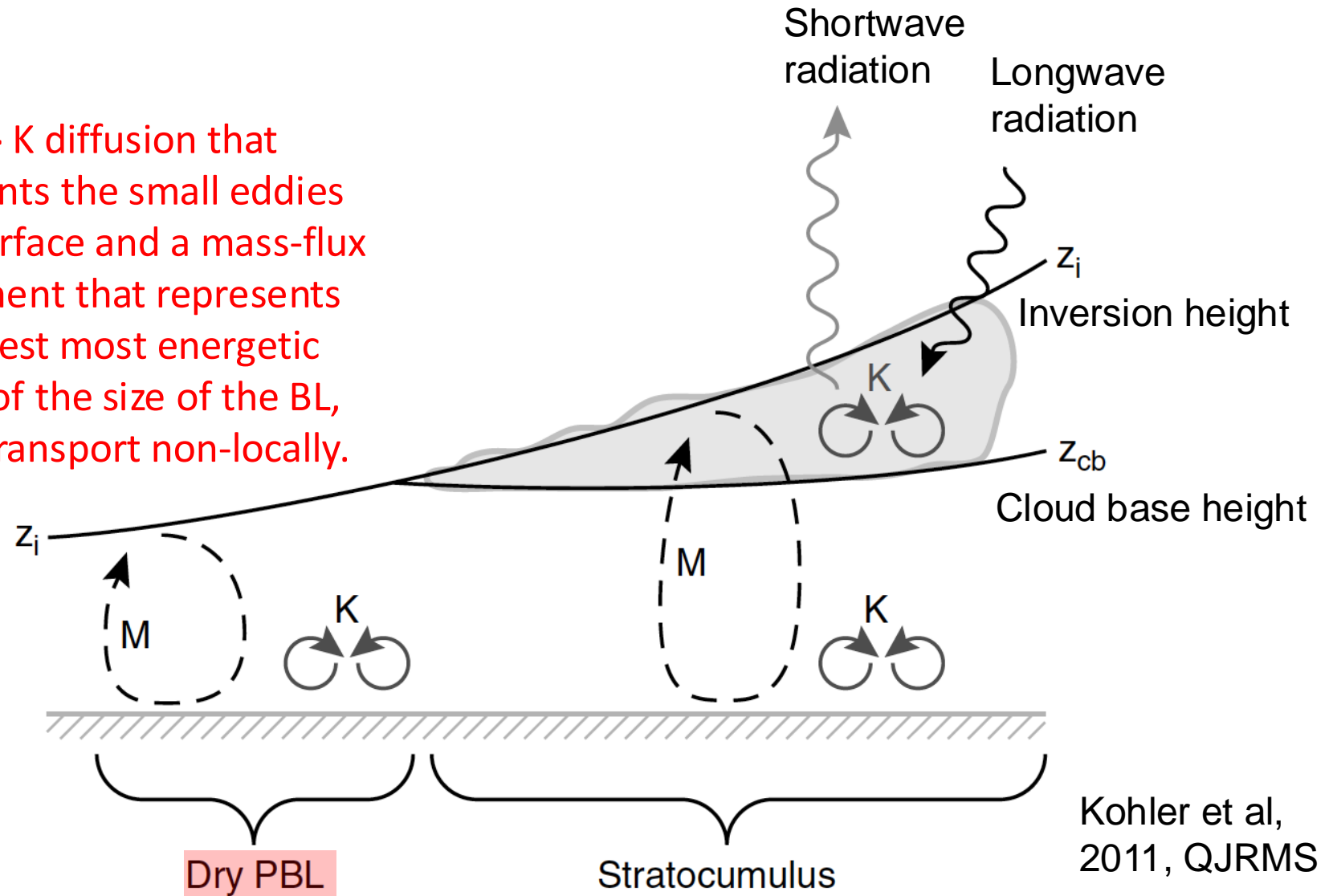


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**Dry BL** -  $K$  diffusion that represents the small eddies from surface and a mass-flux component that represents the largest most energetic eddies of the size of the BL, which transport non-locally.



Kohler et al,  
2011, QJRMS

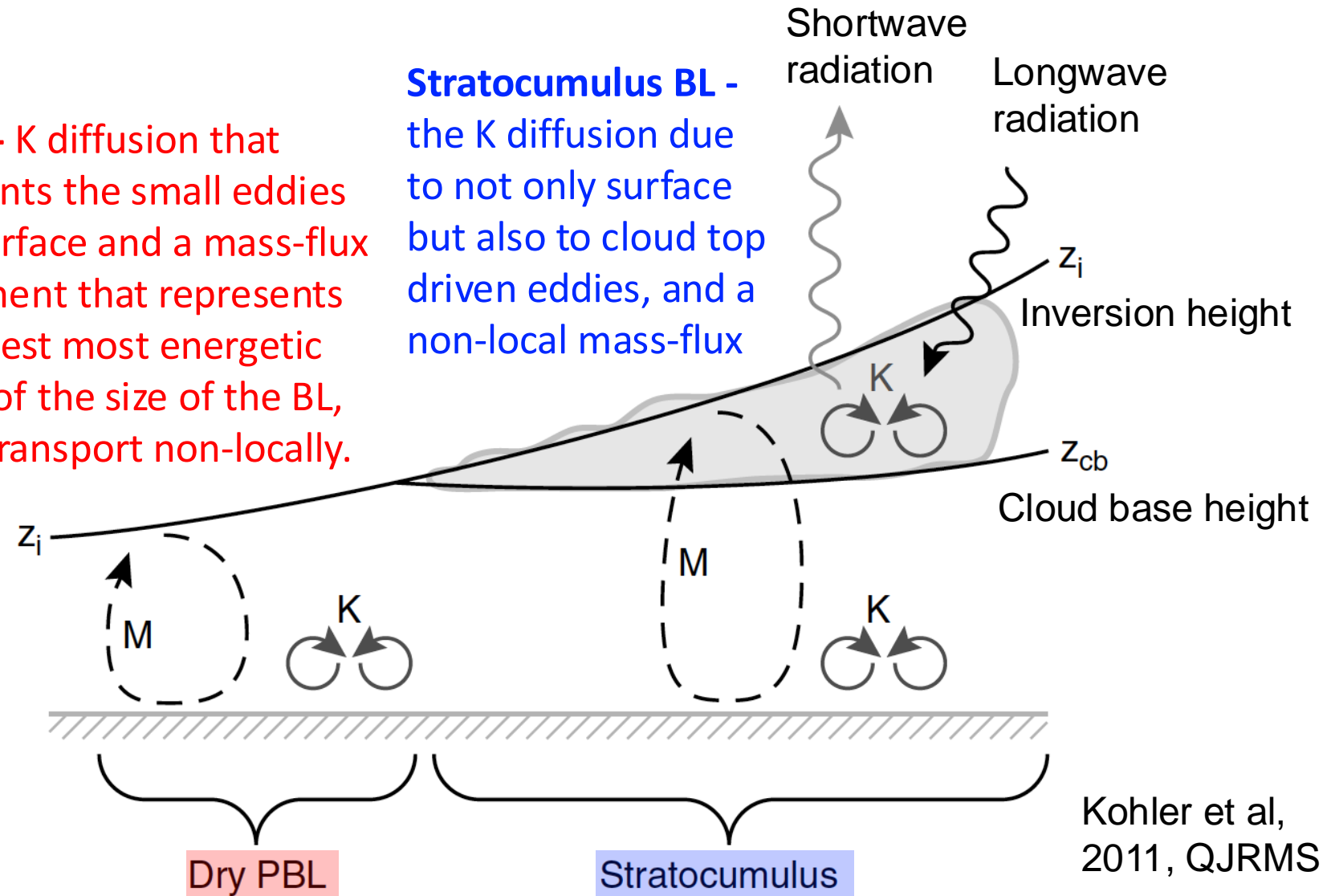
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**Dry BL** - K diffusion that represents the small eddies from surface and a mass-flux component that represents the largest most energetic eddies of the size of the BL, which transport non-locally.

**Stratocumulus BL** - the K diffusion due to not only surface but also to cloud top driven eddies, and a non-local mass-flux



# Description of the current ECMWF IFS scheme

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Stable surface layer

Unstable surface layer

Entrainment level

Lowest model level

$$\text{Surface layer: } \overline{\phi'w'_s} = C_\phi (\overline{\phi_z} - \overline{\phi_s}) |\overline{u_z}|$$

$C_\phi$  in surface layer:

$Ri > 0$  , Holtslag & Bruin (1988)

$Ri < 0$  , Dyer & Hicks (1974)

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Use the relation

$$Ri = \zeta \frac{\Phi_H(\zeta)}{\Phi_M^2(\zeta)}$$

to convert  $\zeta = \frac{z}{L}$  to the gradient Richardson number in the outer layer



# Description of the current IFS scheme

Stable surface layer

Unstable surface layer

$K_\phi$  above surface:

Dyer & Hicks,  $Ri < 0$

Outer layer:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$$

Entrainment level

EDMF:

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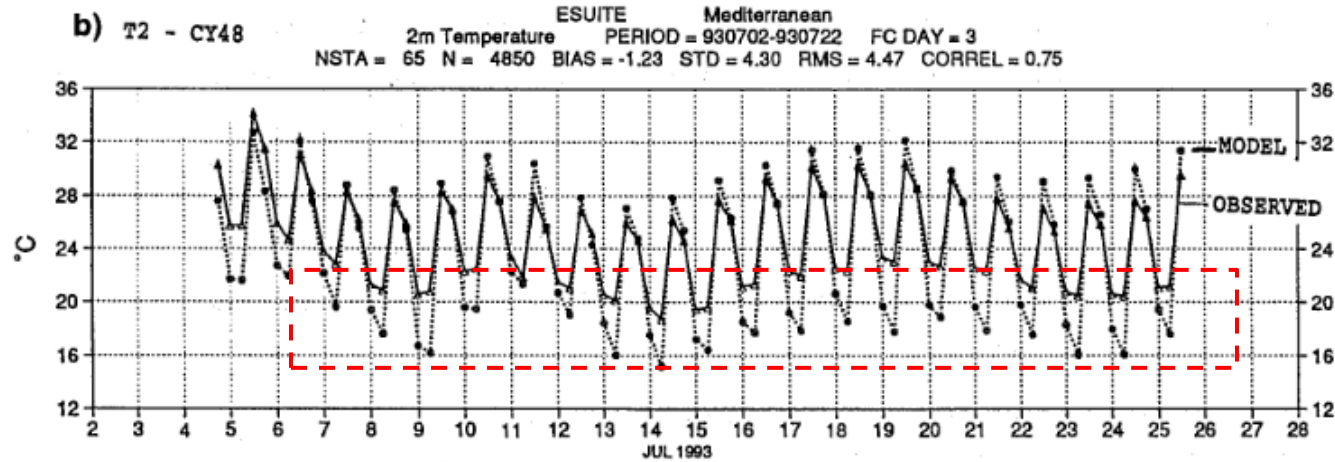
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# Observations can only take us so far....

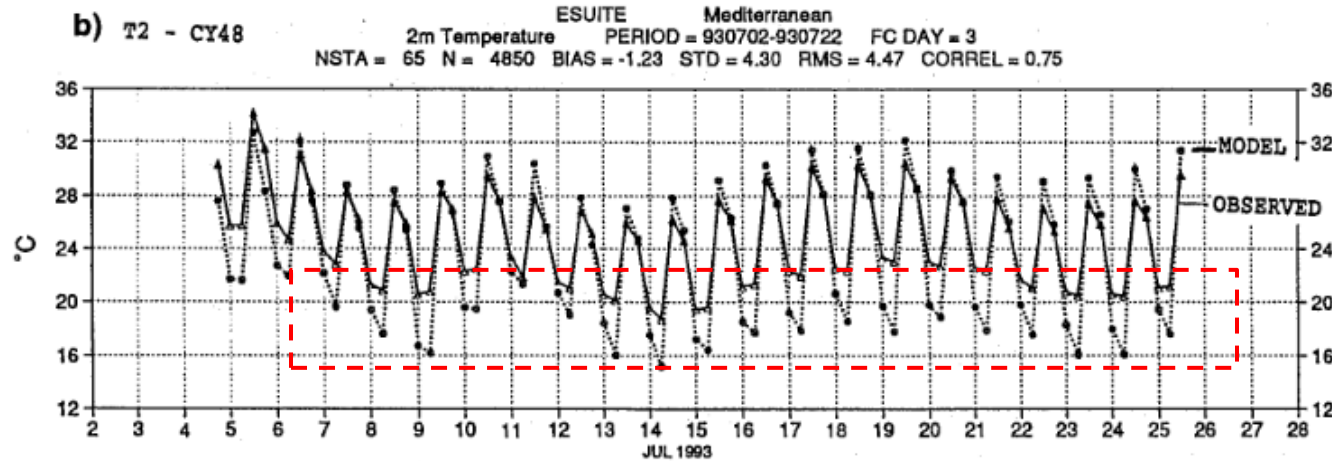
Nighttime (stable regime) temperatures were too cold



Beljaars, 1991

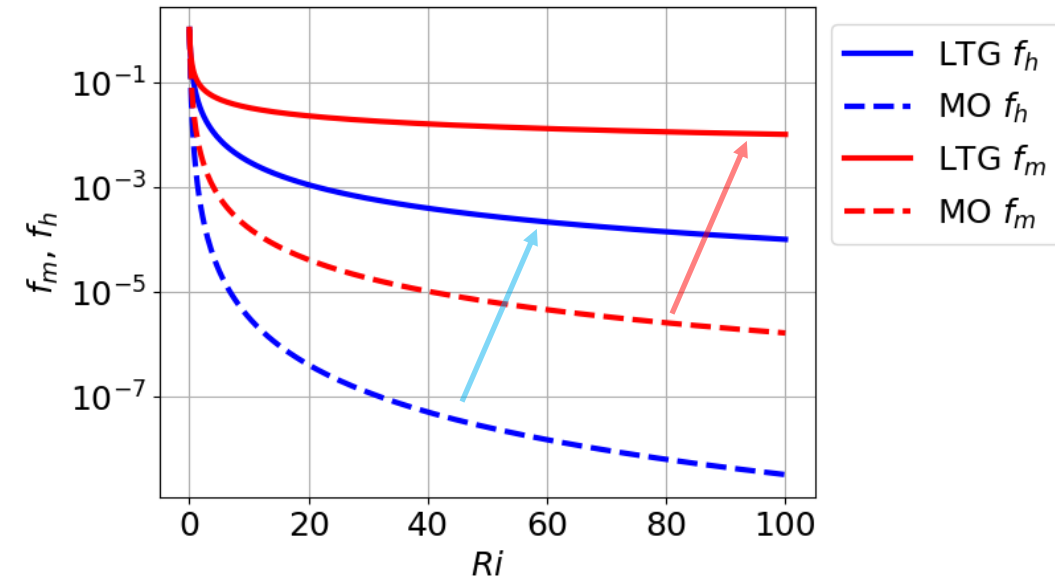
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Mixing was increased in stable BLs

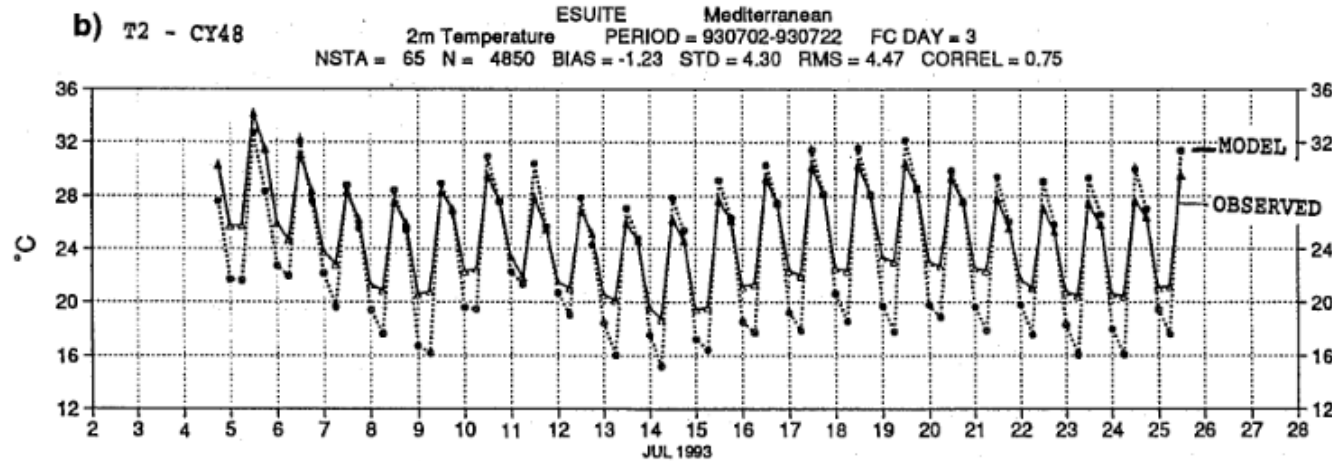


$$\overline{\theta'w'} \sim -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{\theta}}{\partial z} f_H(Ri)$$

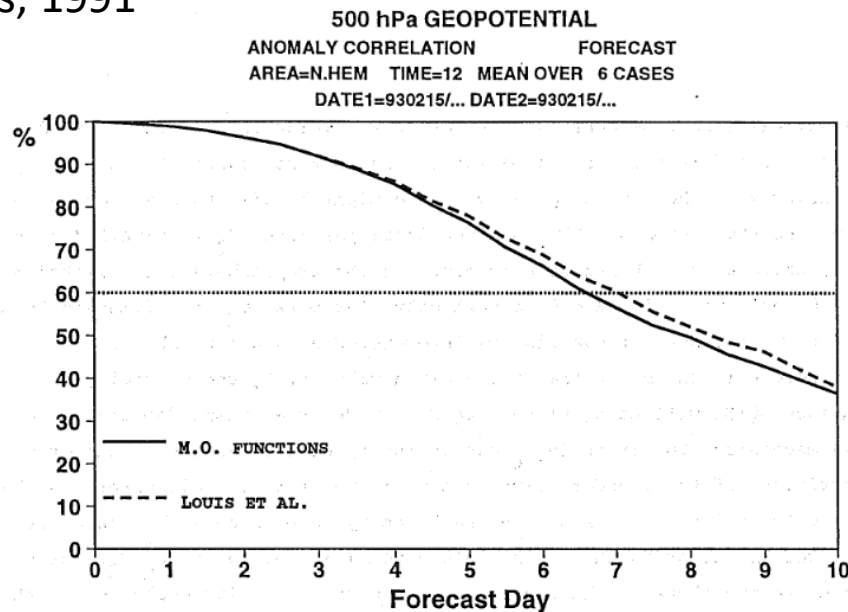
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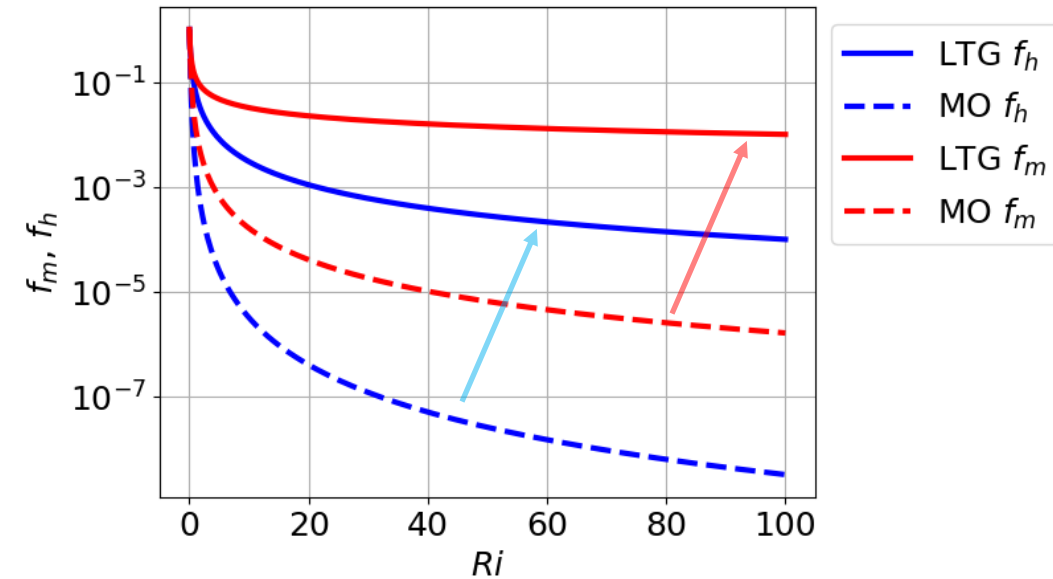


Beljaars, 1991



This was a change  
predominantly  
motivated by forecast  
scores, and not direct  
measurements

Mixing was increased in stable BLs



$$\overline{\theta'w'} \sim -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{\theta}}{\partial z} f_H(Ri)$$

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# Description of the current IFS scheme

'Free' atmosphere:  $\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$

$K_\phi$  above surface:

Stable surface layer

Unstable surface layer

Outer layer:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$$

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Dyer & Hicks,  $Ri < 0$

Louis (1982),  $Ri > 0$

Entrainment level

EDMF:

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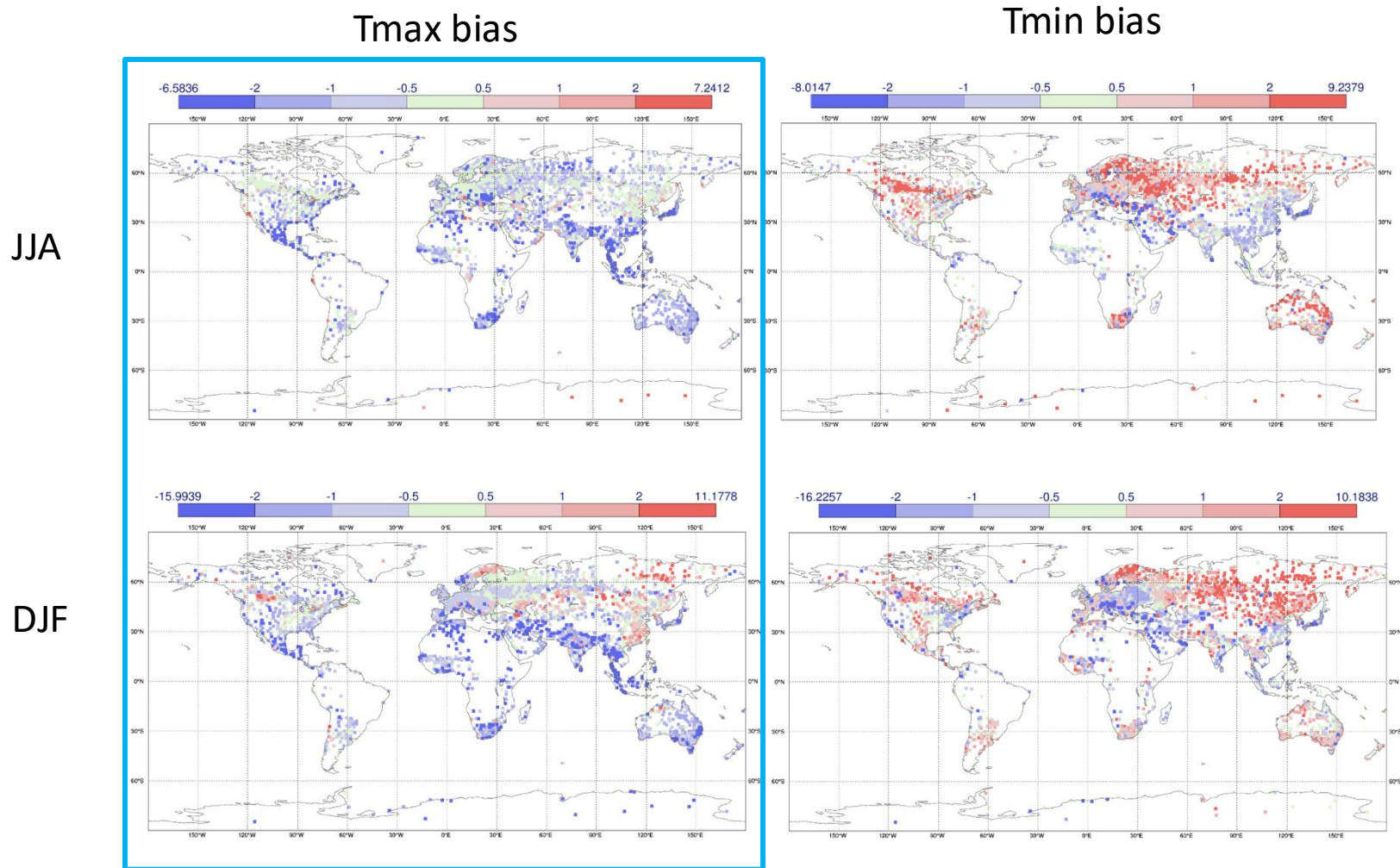
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# Current near-surface model issues in stable boundary layers

# Near surface errors in stable conditions – 2m T



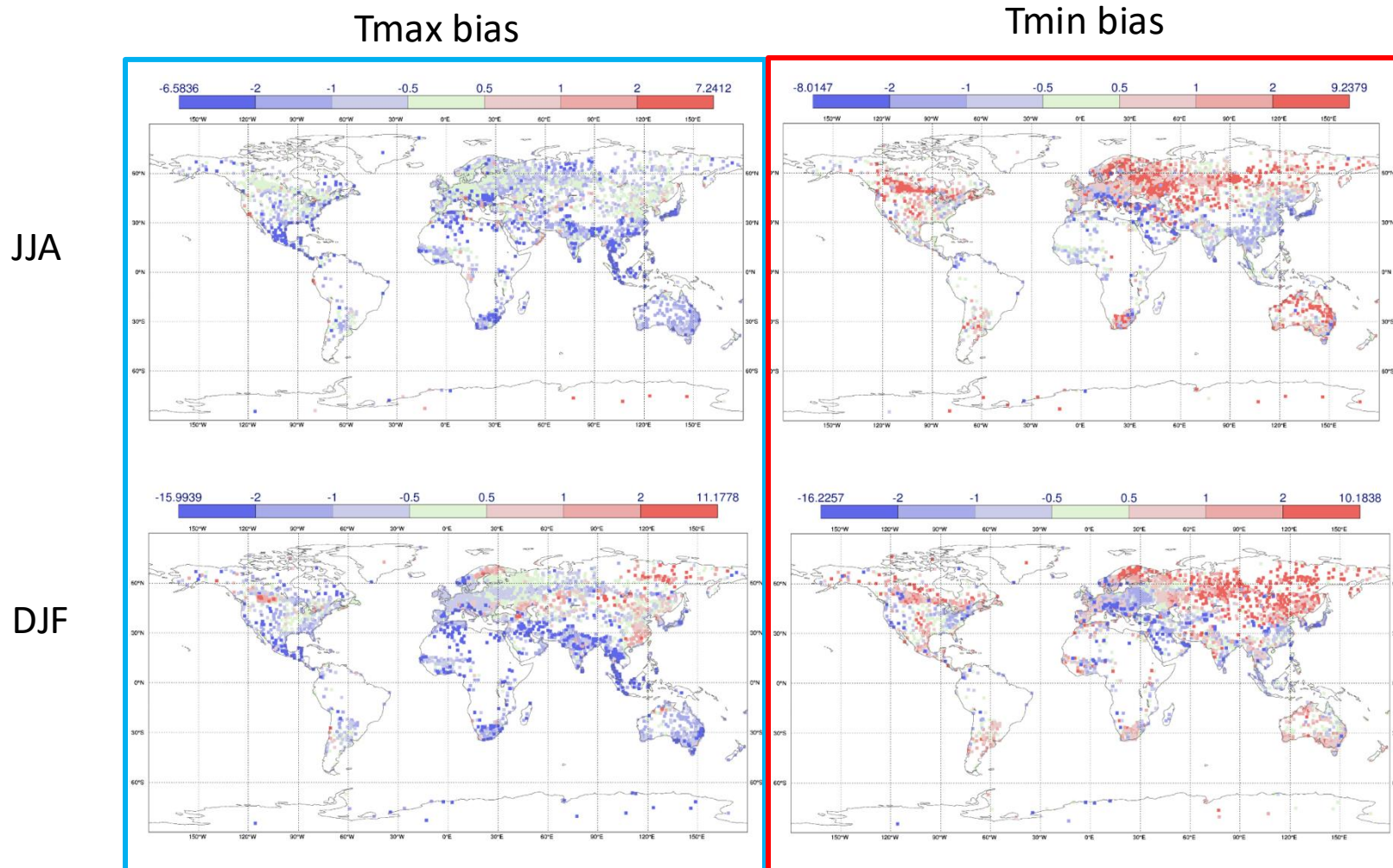
Plots show mean error of maximum (Tmax) and minimum (Tmin) 2m temperatures over 2018/2019 compared with SYNOP observations

Maximum (daytime) temperatures too cold

Sandu et al, 2020, ECMWF Tech memo 875



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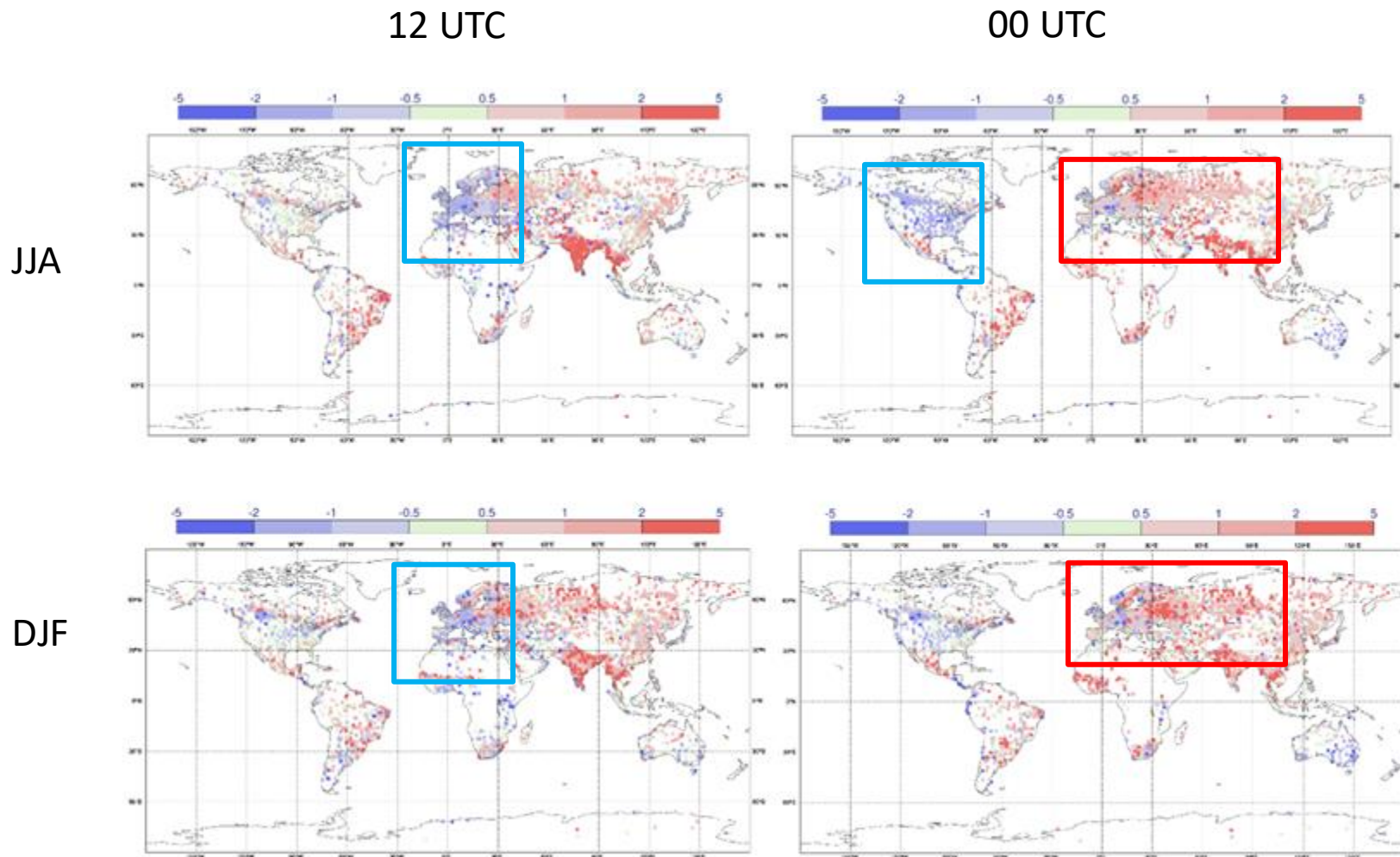
Minimum (nighttime) temperatures too warm

Mirrors the seasonal cycle

Sandu et al, 2020, ECMWF Tech memo 875



# Near surface errors in stable conditions – 10 UV



Plots show mean error of  
**10m wind** over 2018/2019  
compared with SYNOP  
observations

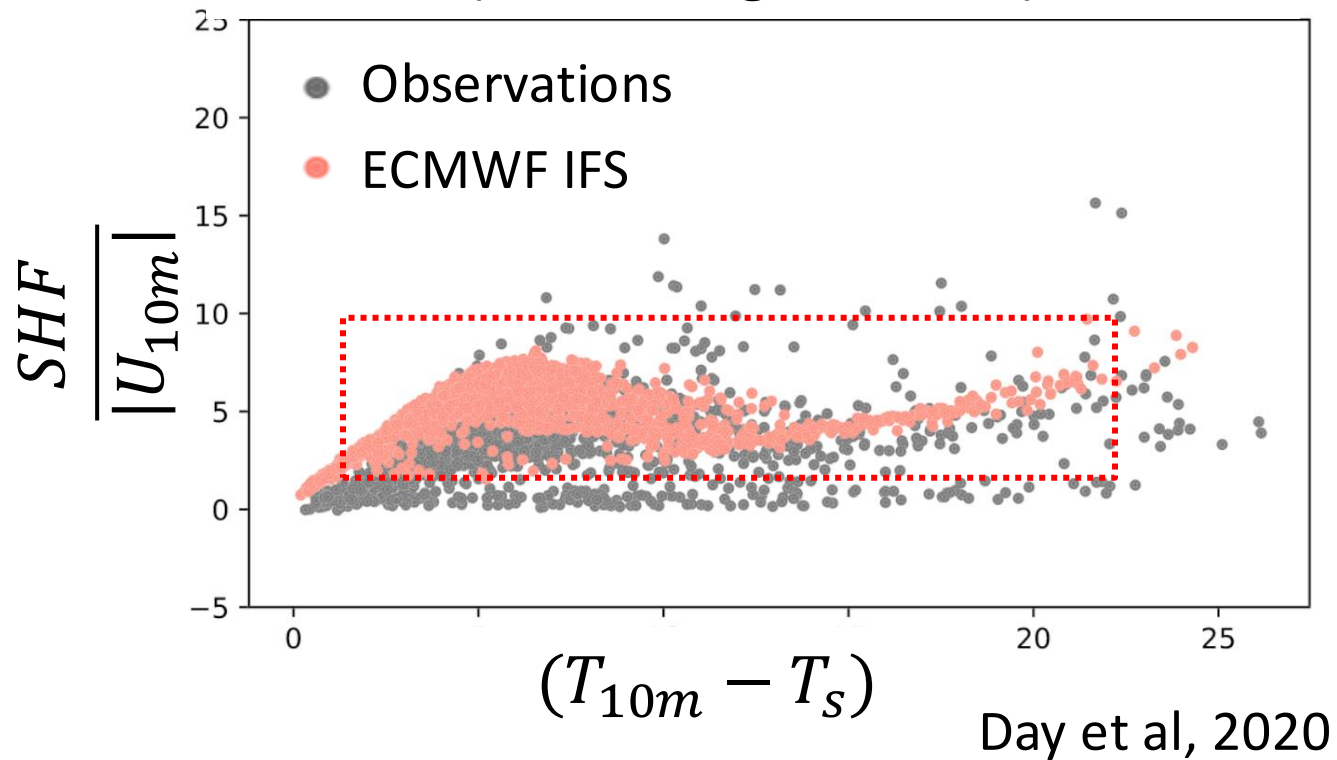
Daytime winds generally too  
weak

Nighttime / wintertime  
winds generally too strong

Sandu et al, 2020, ECMWF Tech memo 875

# Too much turbulent mixing in stable conditions

Surface sensible heat flux  
(Summit, greenland)

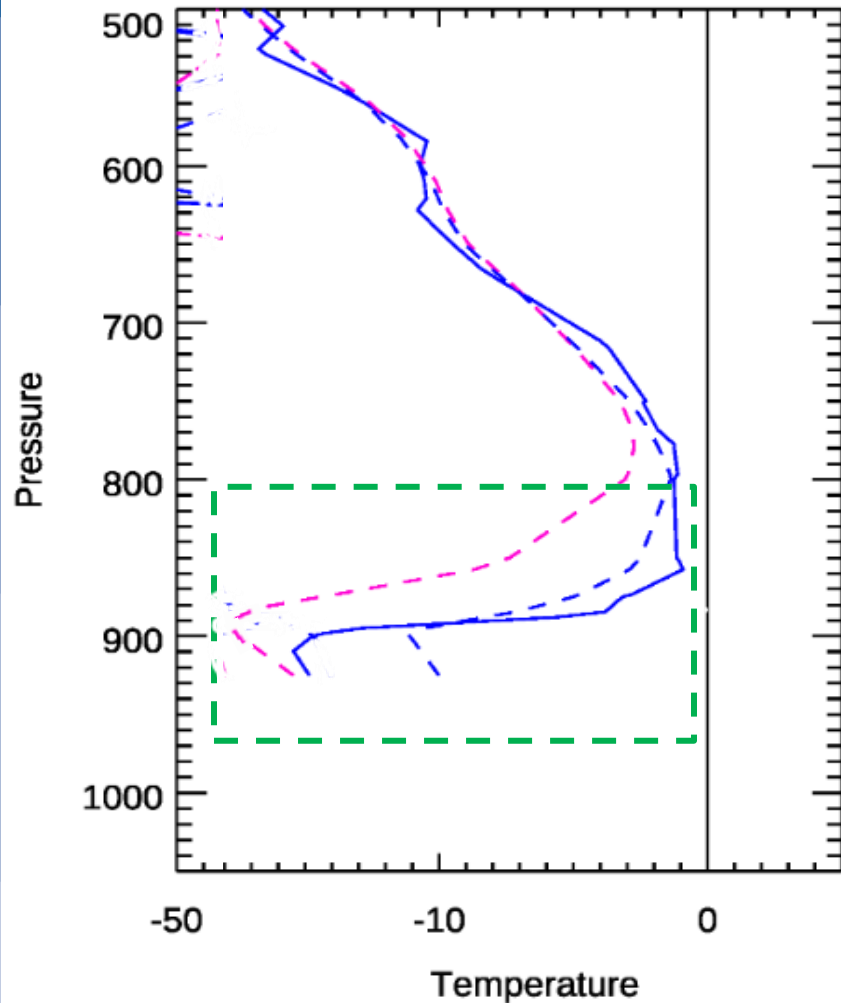


$$SHF = \rho C_H |U_{10m}| (T_{10m} - T_s)$$

Comparison between **model** and **observations** suggest too much turbulent mixing in stable conditions

# Assimilation of 2m temperatures : Issues in stable conditions

Model comparison  
with radiosonde



Edmonton (Canada) radiosonde : **Solid Blue**

- Strong near-surface inversion

Background model: **Dashed blue**

- Several degrees too warm

Analysis: **Dashed magenta**

- Deep increment above inversion

To avoid deep increments:

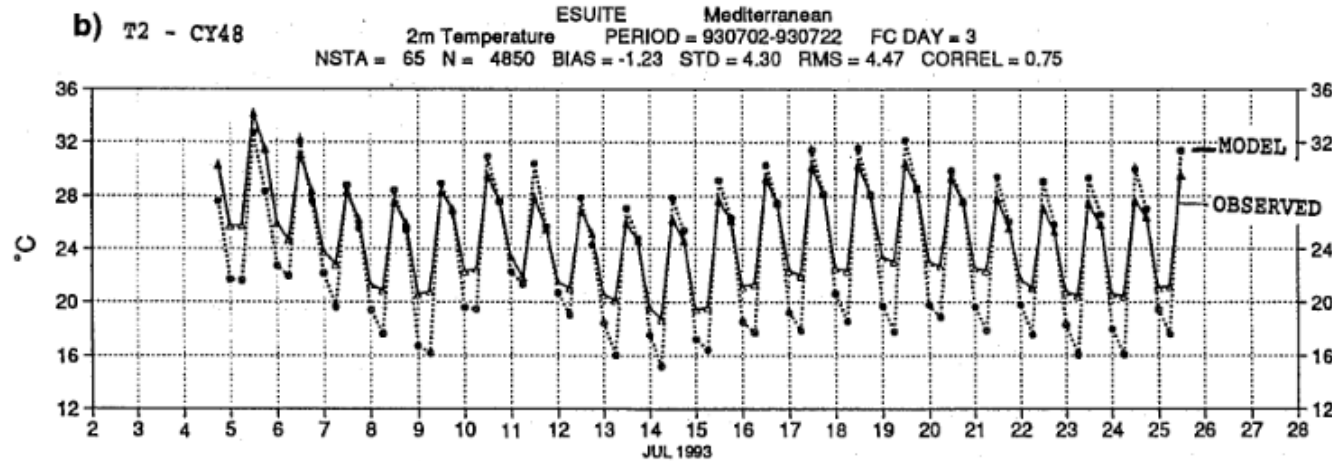
Large differences of the lapse rate adjusted temperature from the background T2m temperature field are given a lower weighting.

Temperature differences of more than 7.5 K are not used.

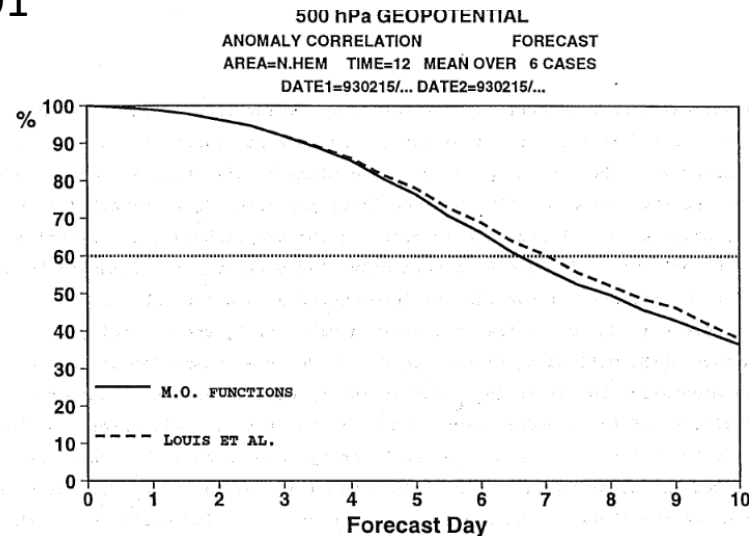
Observations are limited to the first six hours of the 12-hour 4D-Var window to produce more localised increments.

# Remember when there was not enough mixing?

Nighttime (stable regime) temperatures were too cold

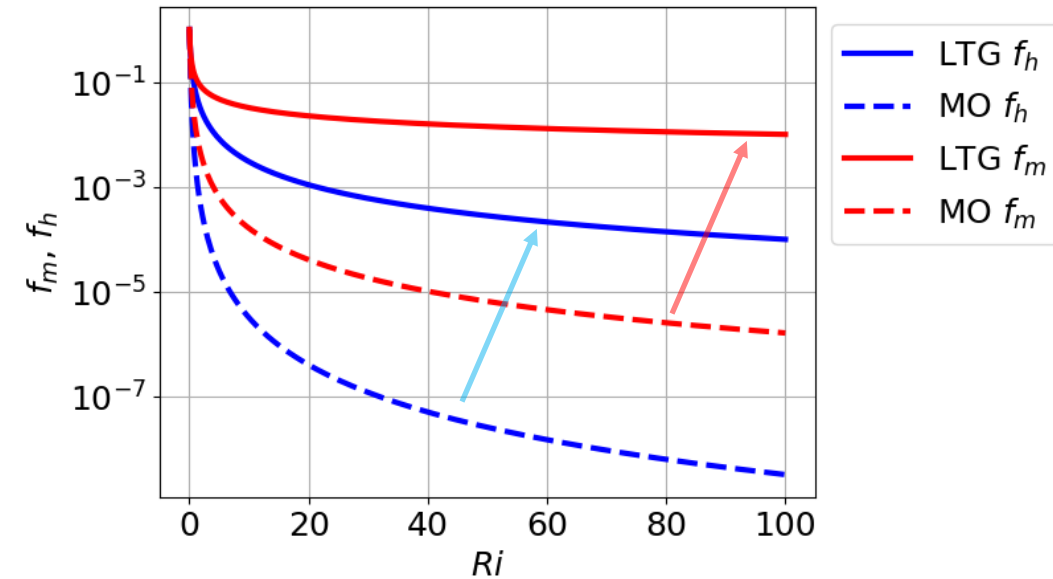


Beljaars, 1991



This was a change  
predominantly  
motivated by forecast  
scores, and not  
measurement

Mixing was increased in stable BLs



$$\overline{\theta'w'} \sim -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{\theta}}{\partial z} f_H(Ri)$$

$$\overline{u'w'} \sim -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{u}}{\partial z} f_M(Ri)$$

# Impact of changing empirical functions

# Impact of changing functions

'Free' atmosphere:  $\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$

$K_\phi$  above surface:

Stable surface layer

Unstable surface layer

Outer layer:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$$

Outer layer:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$$

Dyer & Hicks,  $Ri < 0$

Louis (1982),  $Ri > 0$

Entrainment level

EDMF:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z} + M(\overline{\phi^u} - \overline{\phi^e})$$

Lowest model level

Surface layer:  $\overline{\phi'w'_s} = C_\phi (\overline{\phi_z} - \overline{\phi_s}) |\overline{u_z}|$

$C_\phi$  in surface layer:

$Ri > 0$ , Holtslag & Bruin (1988)

$Ri < 0$ , Dyer & Hicks (1974)

# Impact of changing functions (outer layer)

'Free' atmosphere:  $\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$

$K_\phi$  above surface:

Stable surface layer

Unstable surface layer

Outer layer:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$$

Outer layer:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z}$$

Dyer & Hicks,  $Ri < 0$

Hogström (1988),  $Ri > 0$

Entrainment level

EDMF:

$$\overline{\phi'w'} = -K_\phi \frac{\partial \phi}{\partial z} + M(\overline{\phi}^u - \overline{\phi}^e)$$

Lowest model level

Surface layer:  $\overline{\phi'w'_s} = C_\phi (\overline{\phi}_z - \overline{\phi}_s) |\overline{u}_z|$

$C_\phi$  in surface layer:

$Ri > 0$ , Holtslag & Bruin (1988)

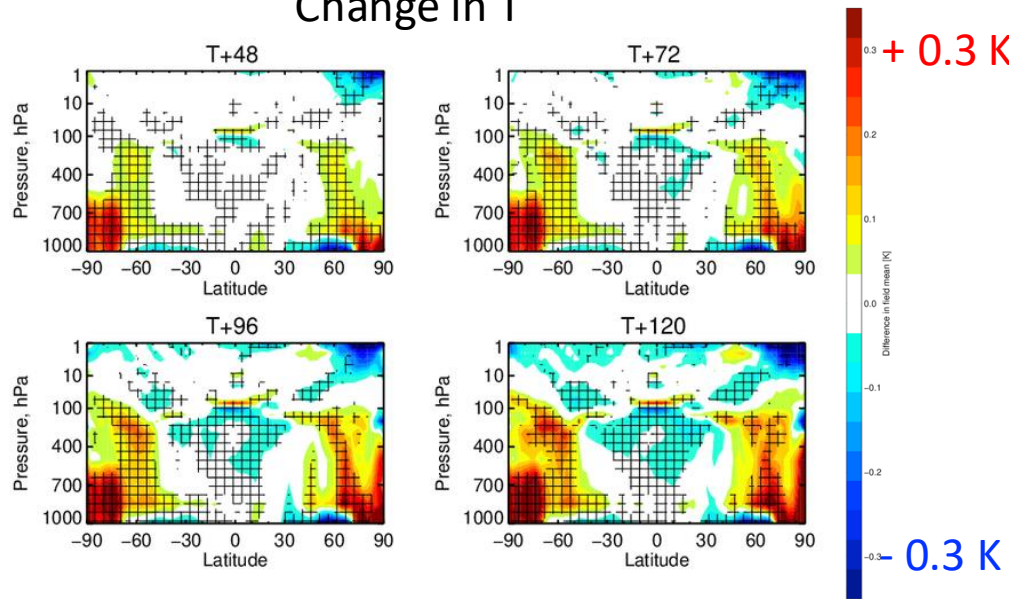
$Ri < 0$ , Dyer & Hicks (1974)



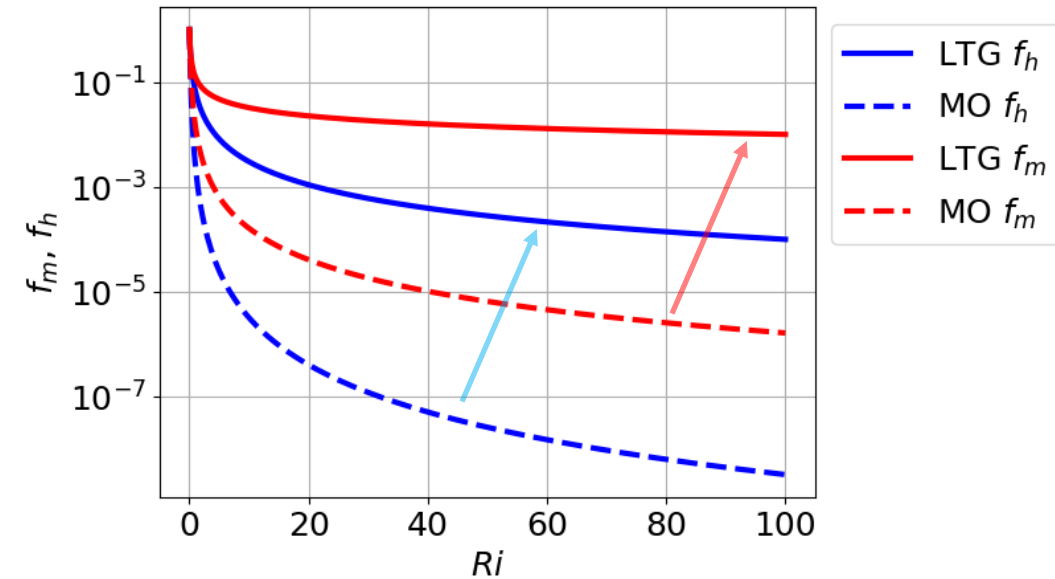
# Changing these back has a large impact on the forecast...

Change in T

Cooling near the surface and heating above – less mixing



Mixing reduced again in stable BLs



$$\overline{\theta'w'} \sim -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{\theta}}{\partial z} f_H(Ri)$$

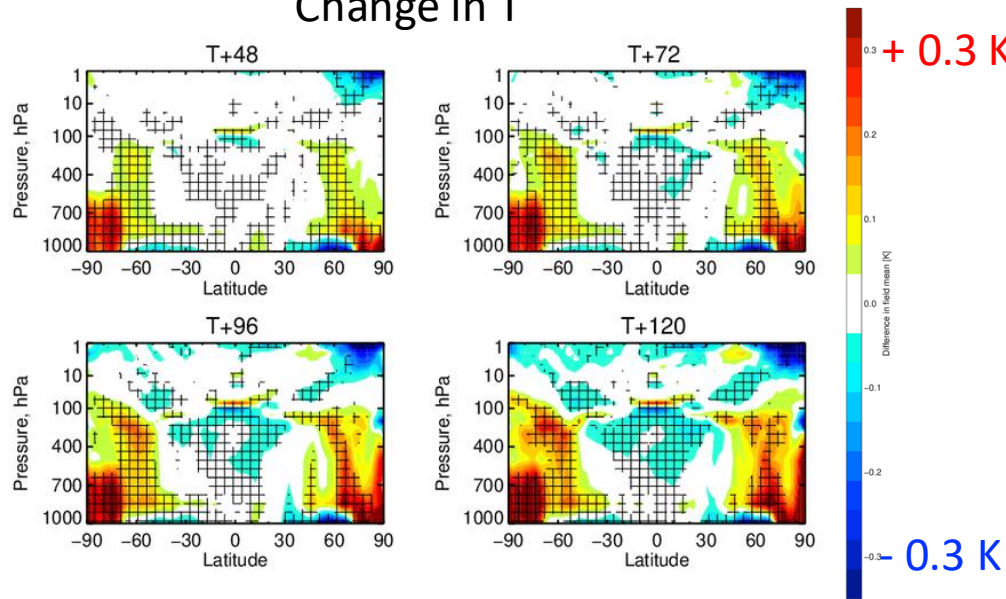
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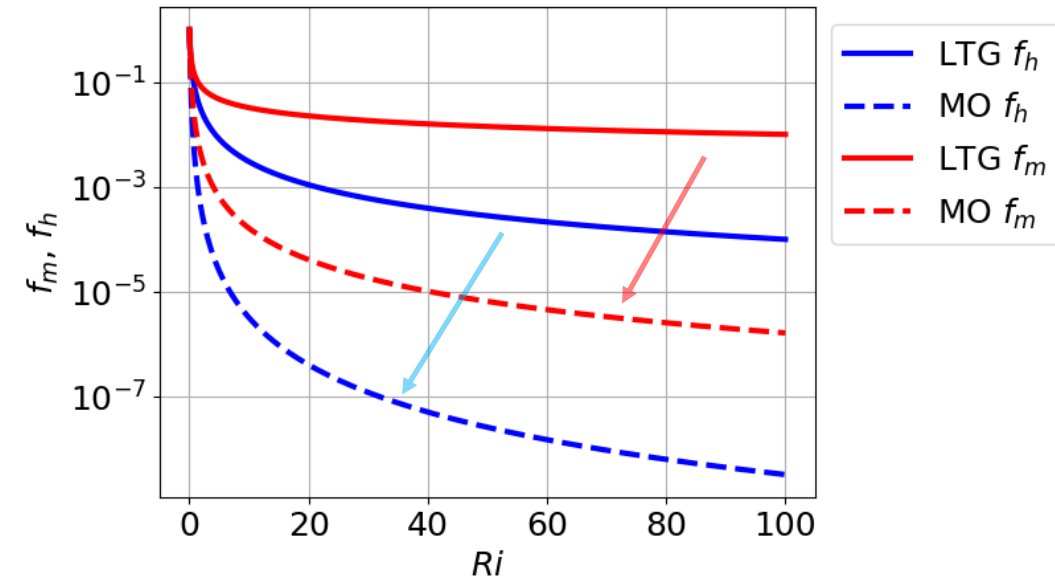
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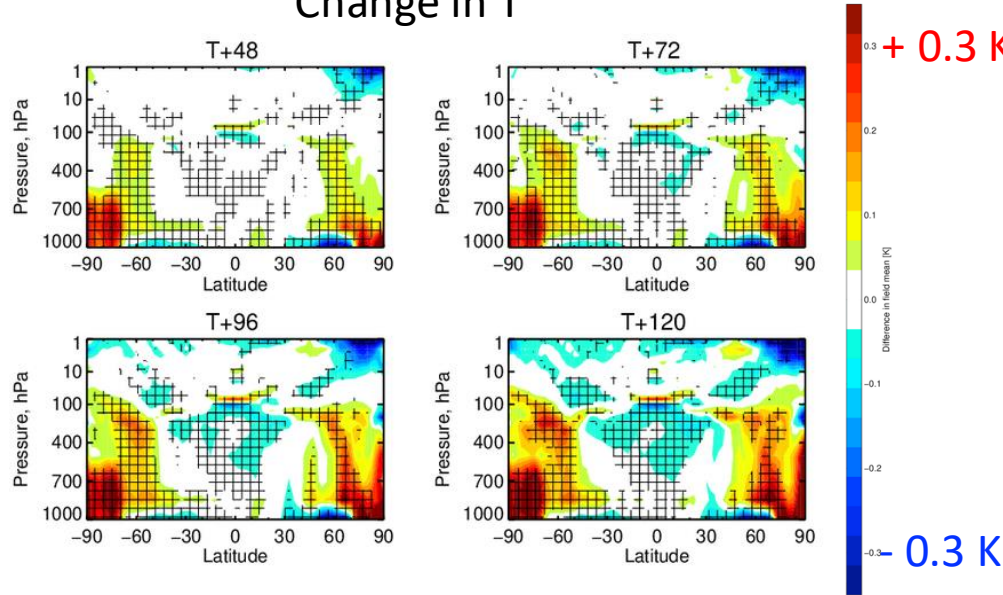
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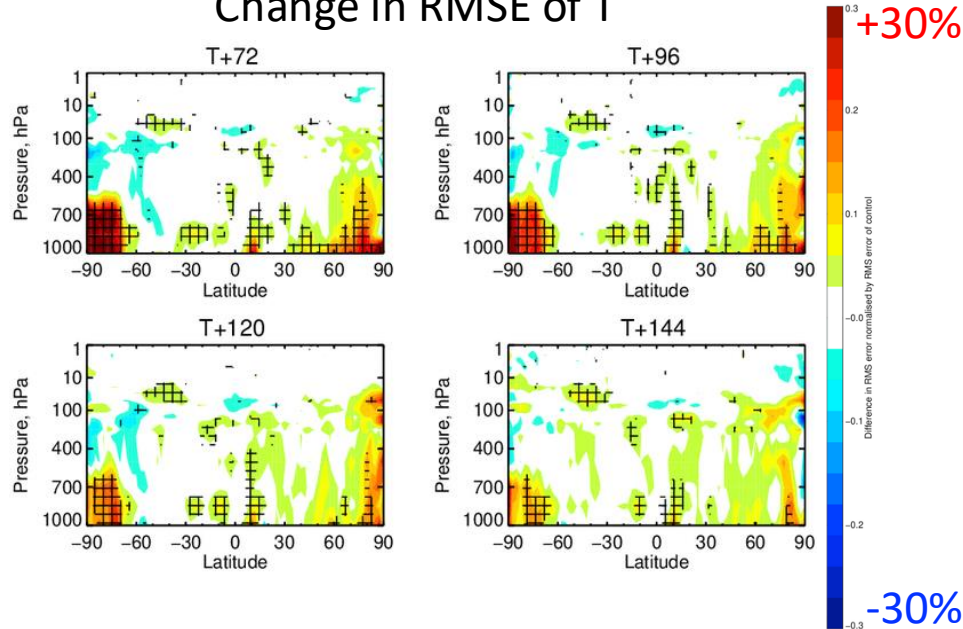
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Cooling near the surface and heating above – less mixing

Change in T

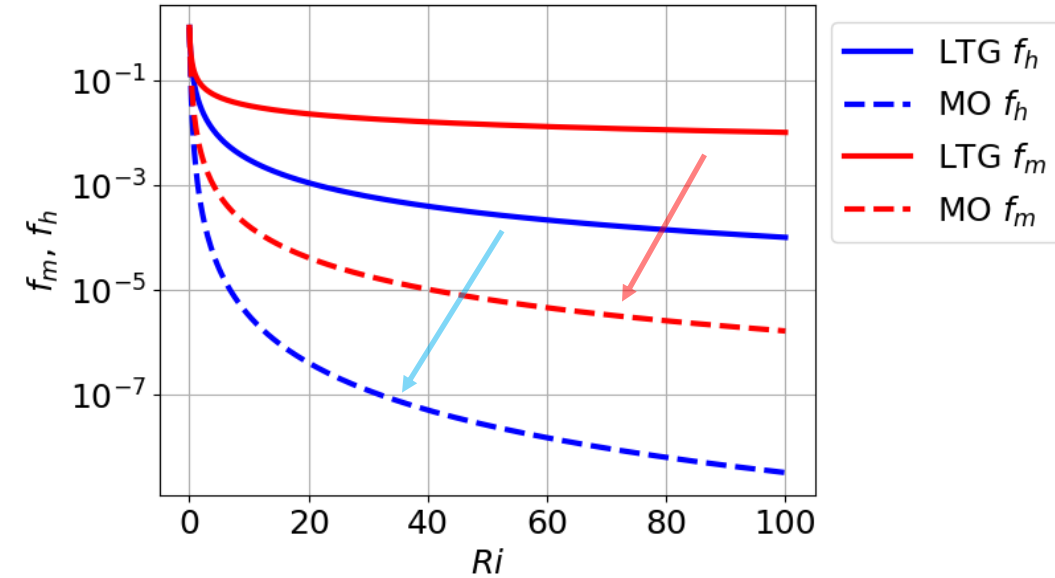


Change in RMSE of T



Large degradation of temperature forecast

Mixing reduced again in stable BLs



$$\overline{\theta'w'} \sim -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| \frac{\partial \bar{\theta}}{\partial z} f_H(Ri)$$

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# Empirical stability functions

- SHEBA (very stable)



Mix well to form:

Richardson number:

$$Ri = \frac{g}{\theta} \frac{\frac{\partial \theta}{\partial z}}{\frac{\partial U}{\partial z}}$$

Dimensionless wind shear:

$$\Phi_M = \frac{\kappa z}{u_*} \frac{\partial U}{\partial z}$$

Dimensionless temperature gradient:

$$\Phi_H = \frac{\kappa z}{\theta_*} \frac{\partial \theta}{\partial z}$$

Dimensionless height:

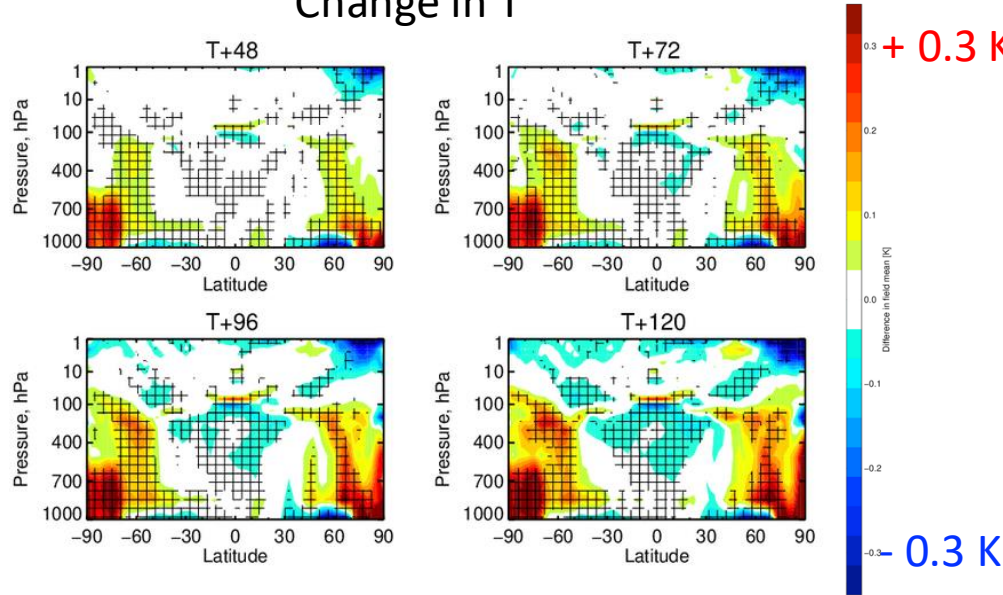
$$\zeta = \frac{z}{L} = z \frac{\overline{\kappa g \theta' w'}}{\theta u_*^3}$$



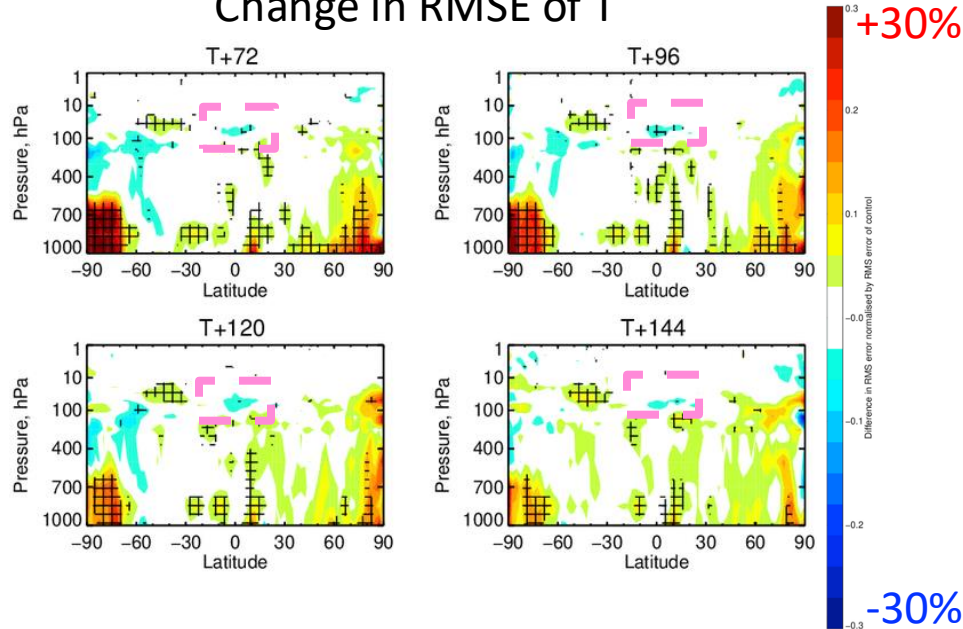
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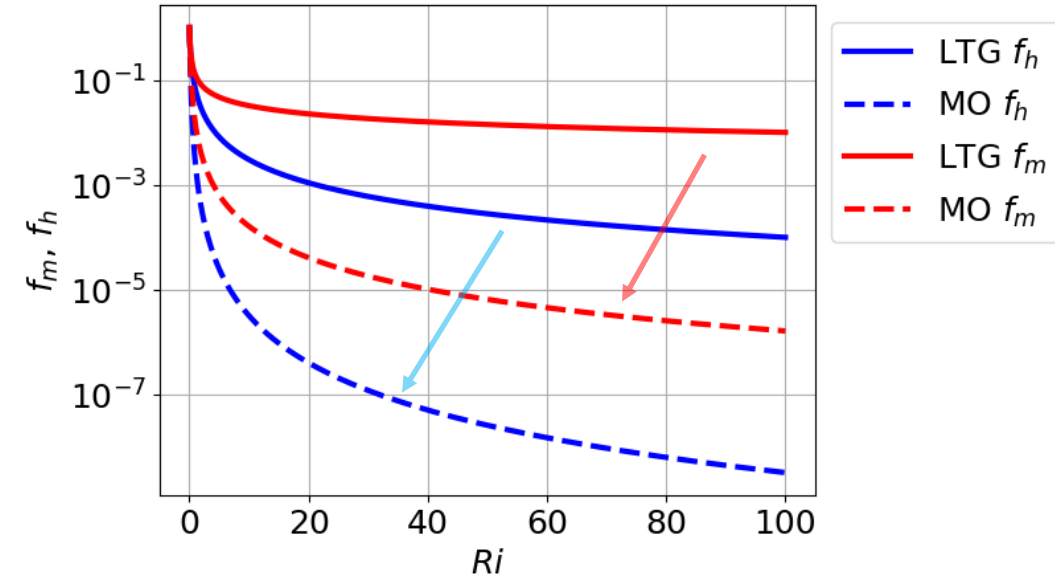


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# Impact of changing functions (stratosphere)

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Louis (1982),  $Ri > 0$

Entrainment level

Above Tropopause:

EDMF:

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Hogström (1988),  $Ri > 0$

Lowest model level

$$\text{Surface layer: } \overline{\phi'w'_s} = C_\phi (\overline{\phi_z} - \overline{\phi_s}) |\overline{u_z}|$$

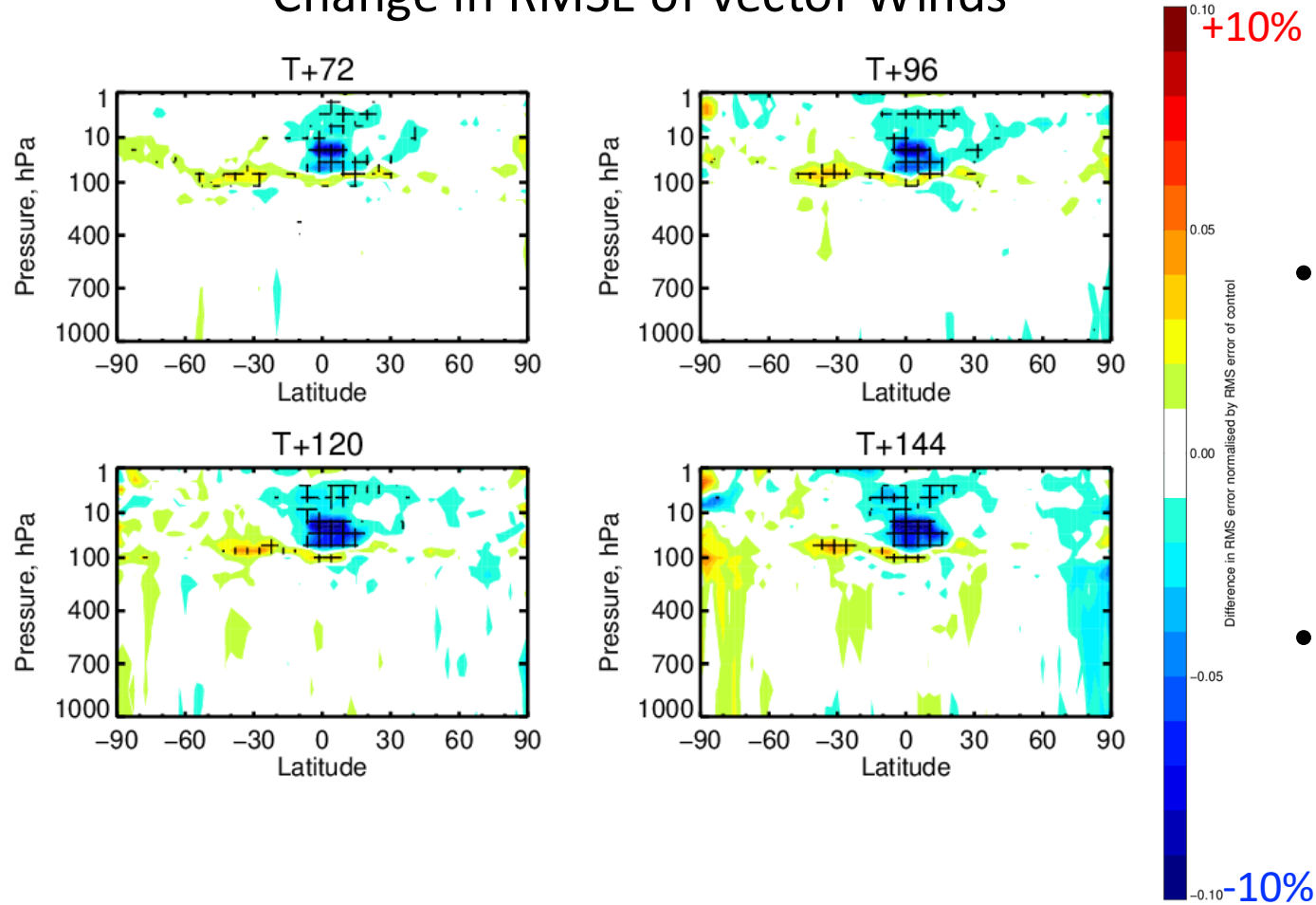
$C_\phi$  in surface layer:

$Ri > 0$ , Holtslag & Bruin (1988)

$Ri < 0$ , Dyer & Hicks (1974)

# Impact of changing functions (stratosphere)

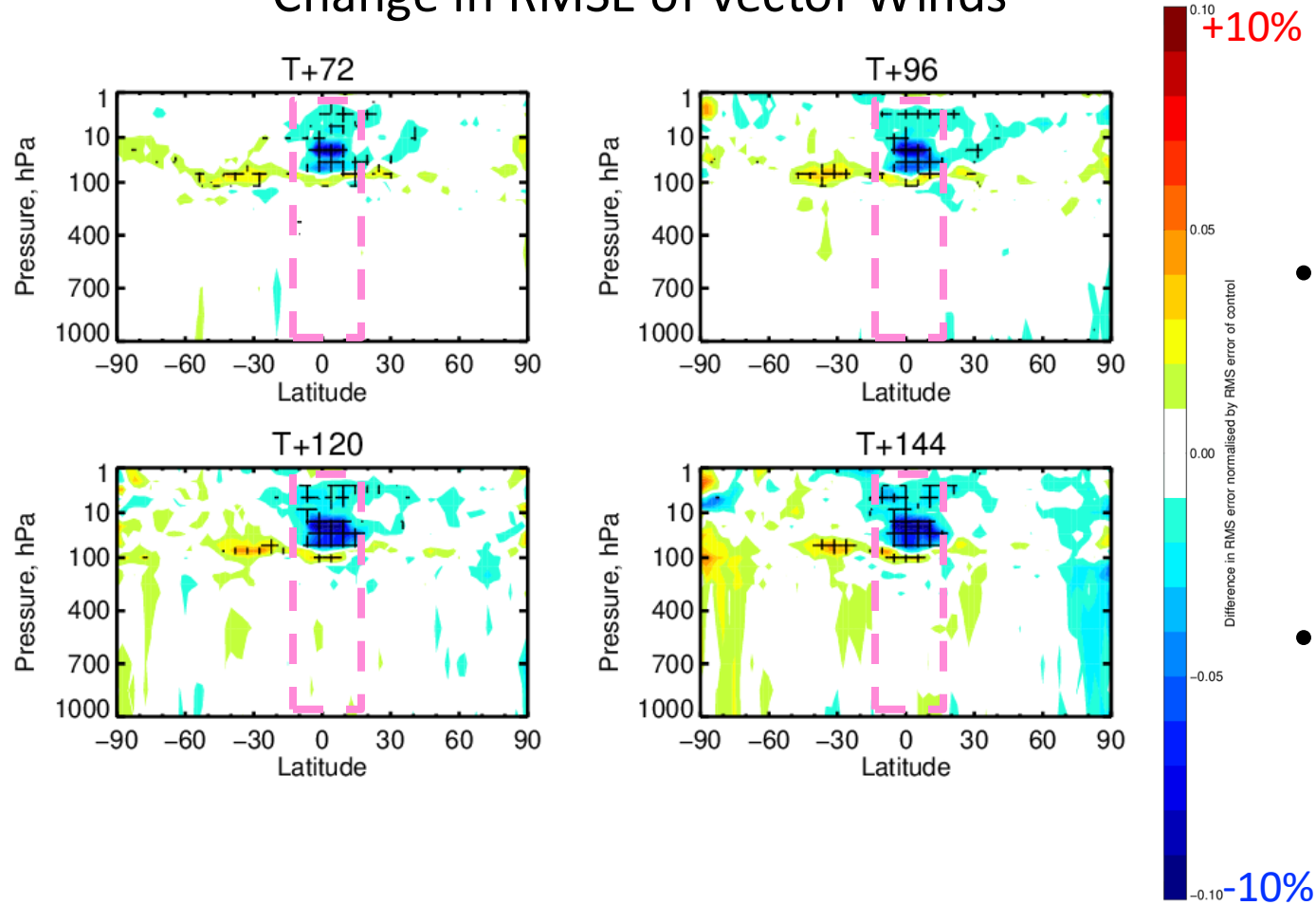
Change in RMSE of vector Winds



- Typically, the same exchange coefficients are used in the stratosphere as in the outer / mixed layer
- There is little constraint on the exchange coefficients in the stratosphere, where the flow is very stable
- Reducing diffusion in the stratosphere (above the tropopause) leads to improved winds and a better Quasi-biennial Oscillation of the winds in the tropics

# Impact of changing functions (stratosphere)

Change in RMSE of vector Winds



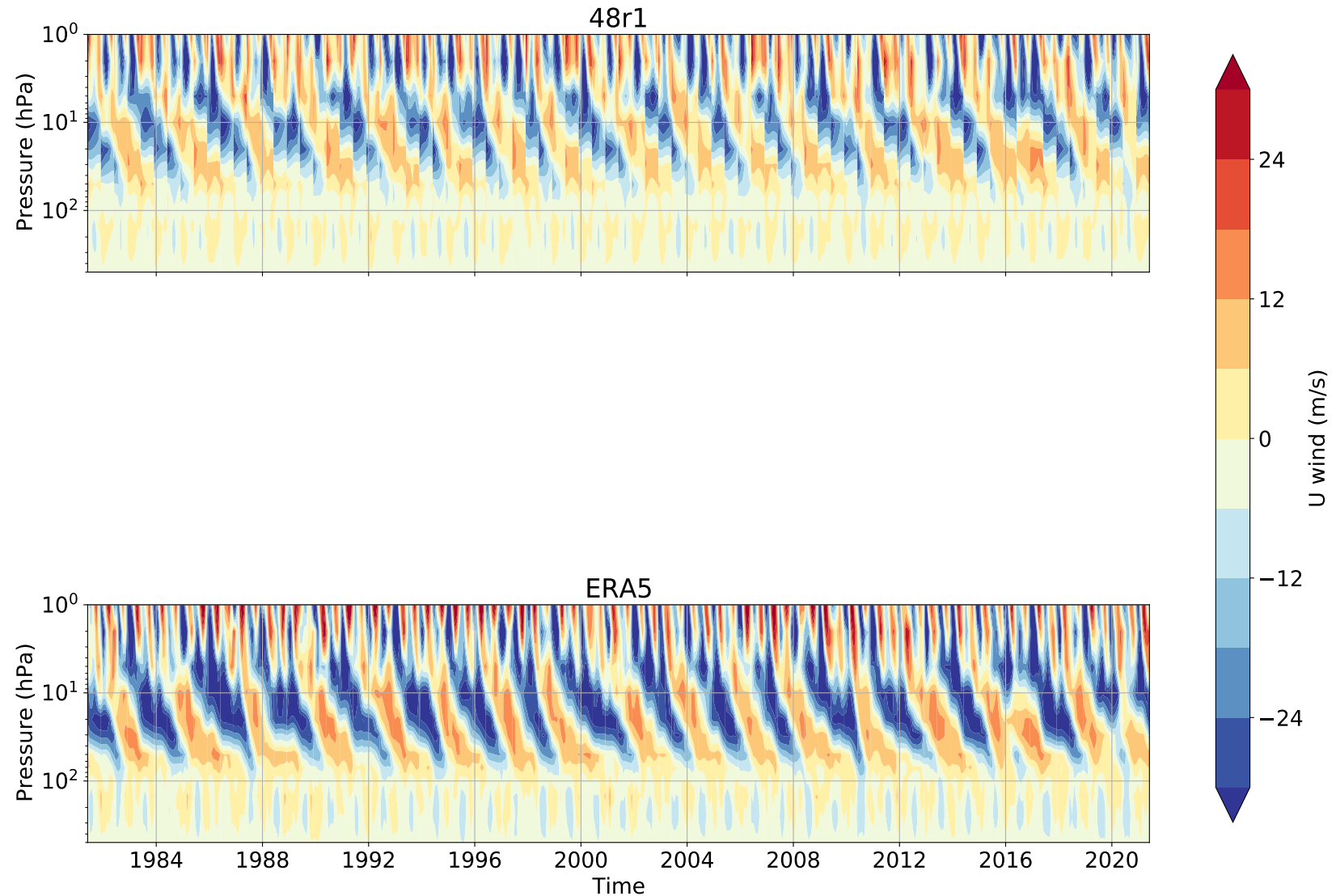
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# Turbulent diffusion in the stratosphere

## Zonal winds averaged between 5S – 5N

Seasonal hindcasts run with the  
ECMWF IFS, 7 months long

The Quasi-biennial Oscillation (QBO)  
has too weak amplitude and does not  
descend far enough



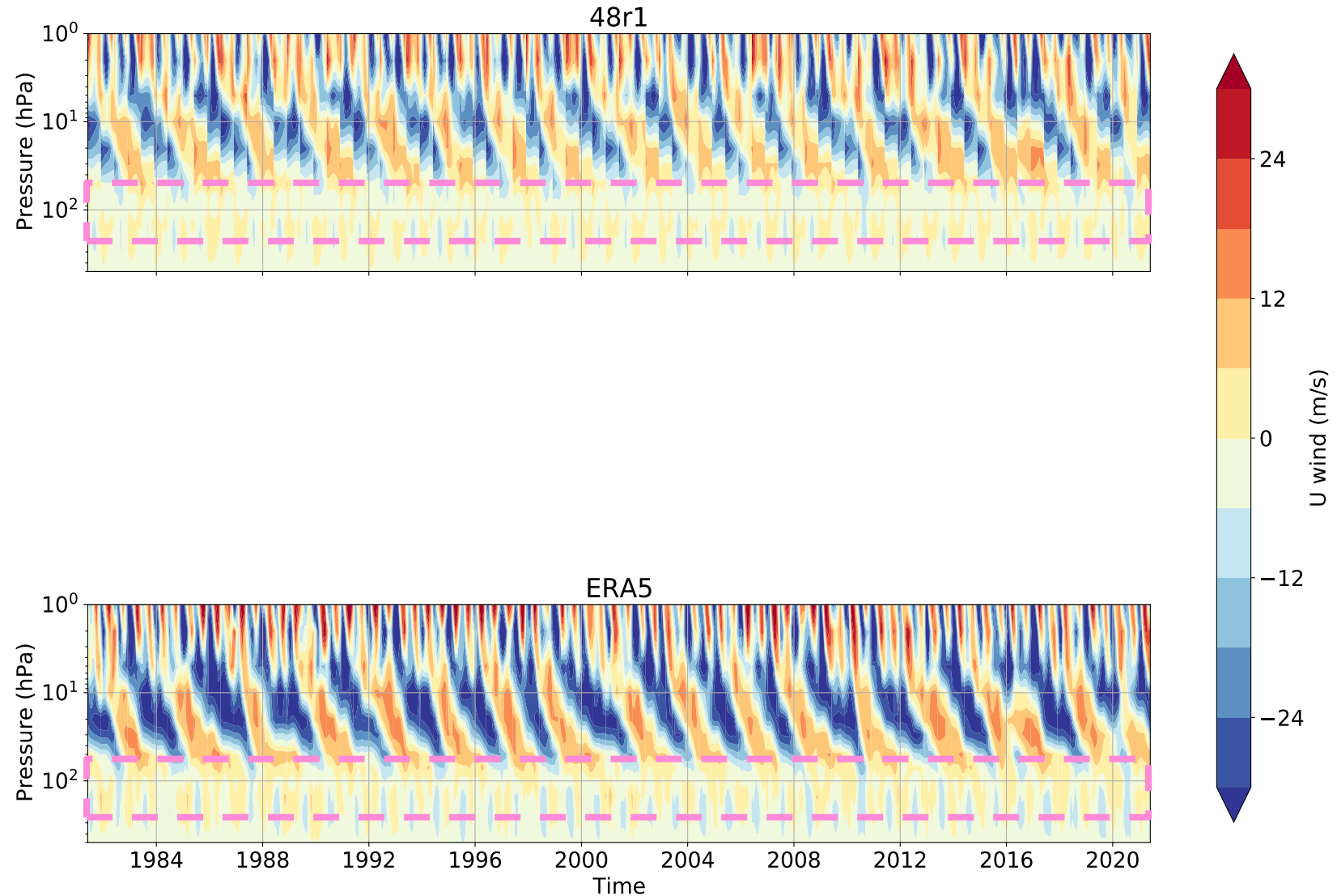


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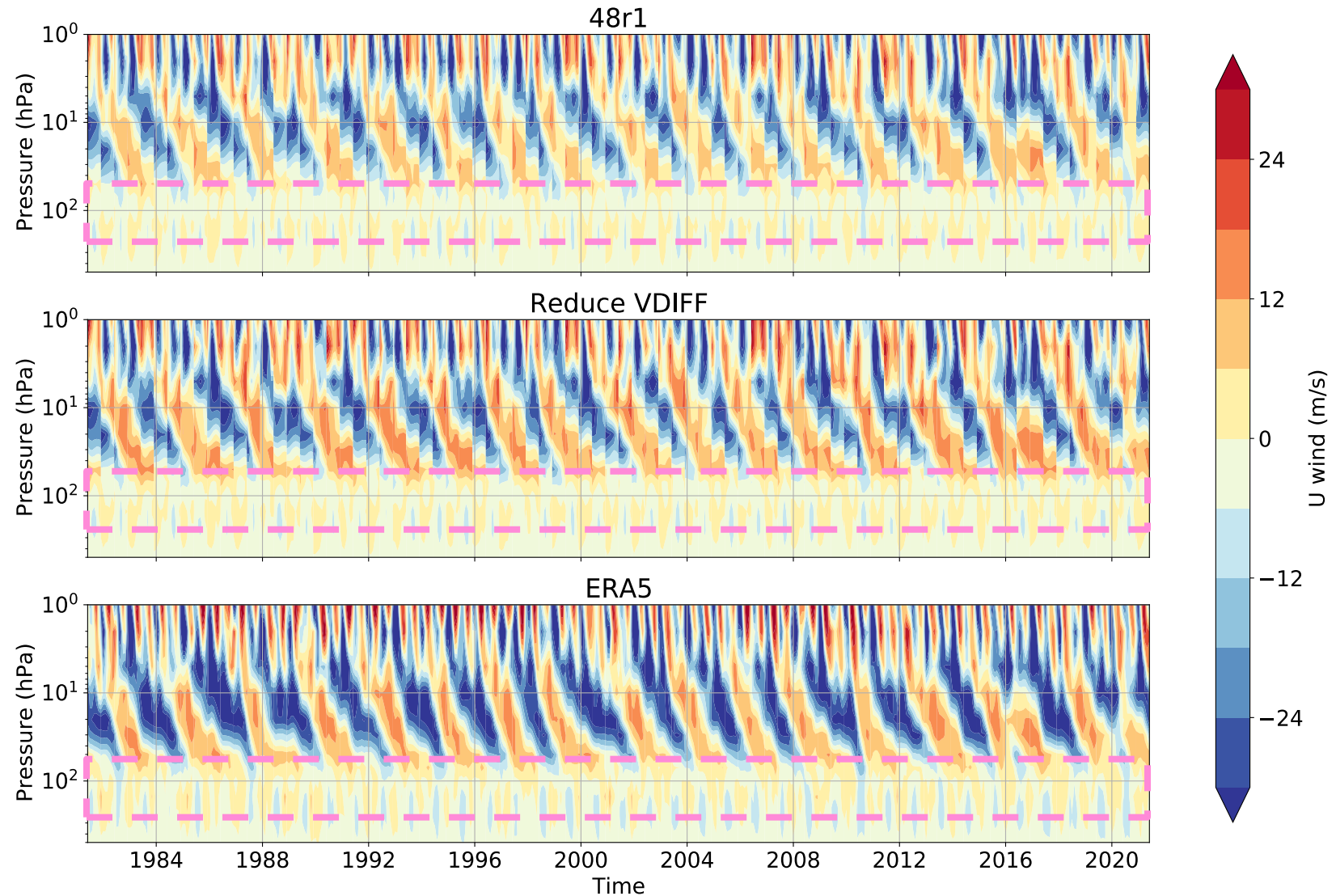
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Reducing vertical diffusion in the  
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amplitude and slightly improves its  
descent



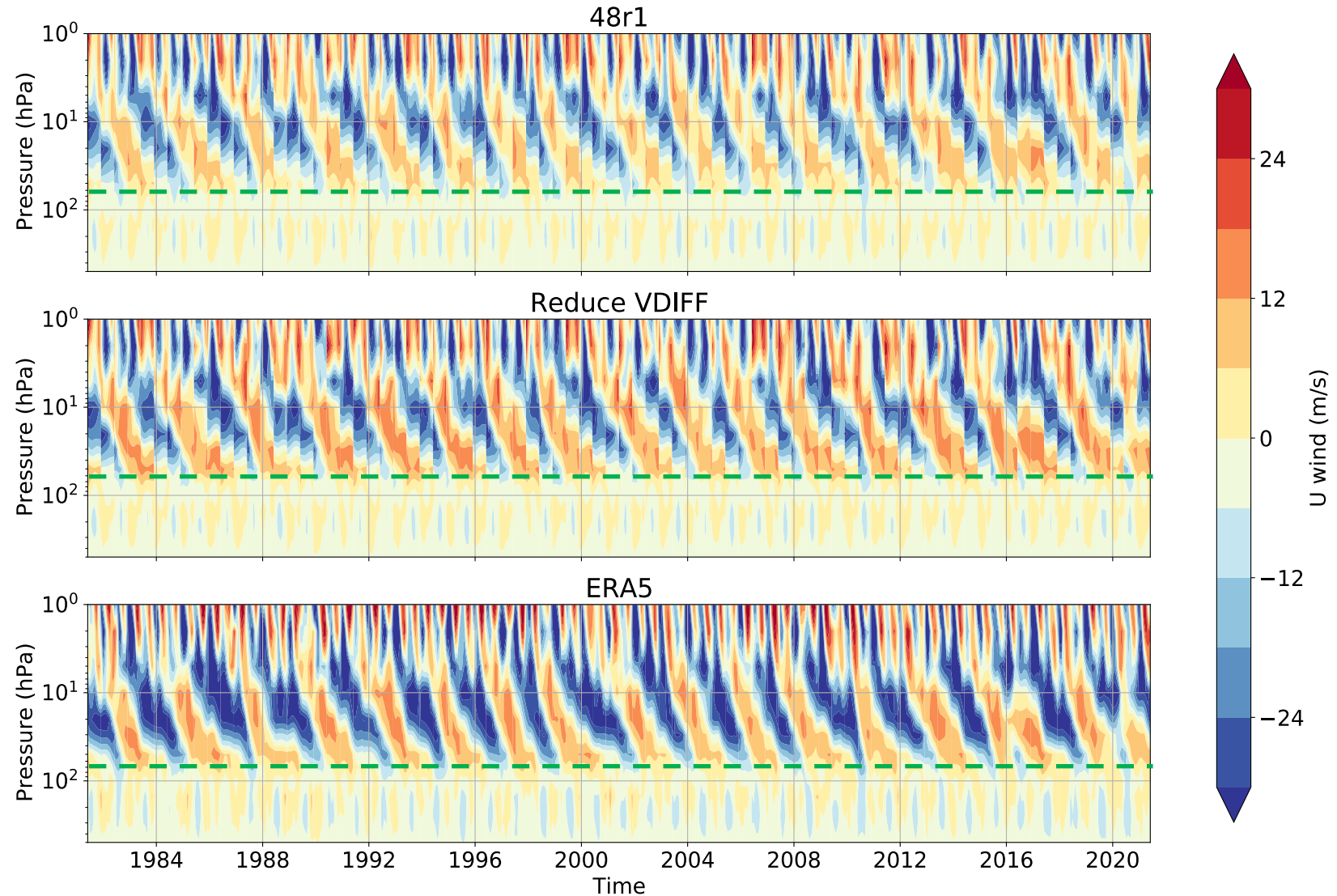
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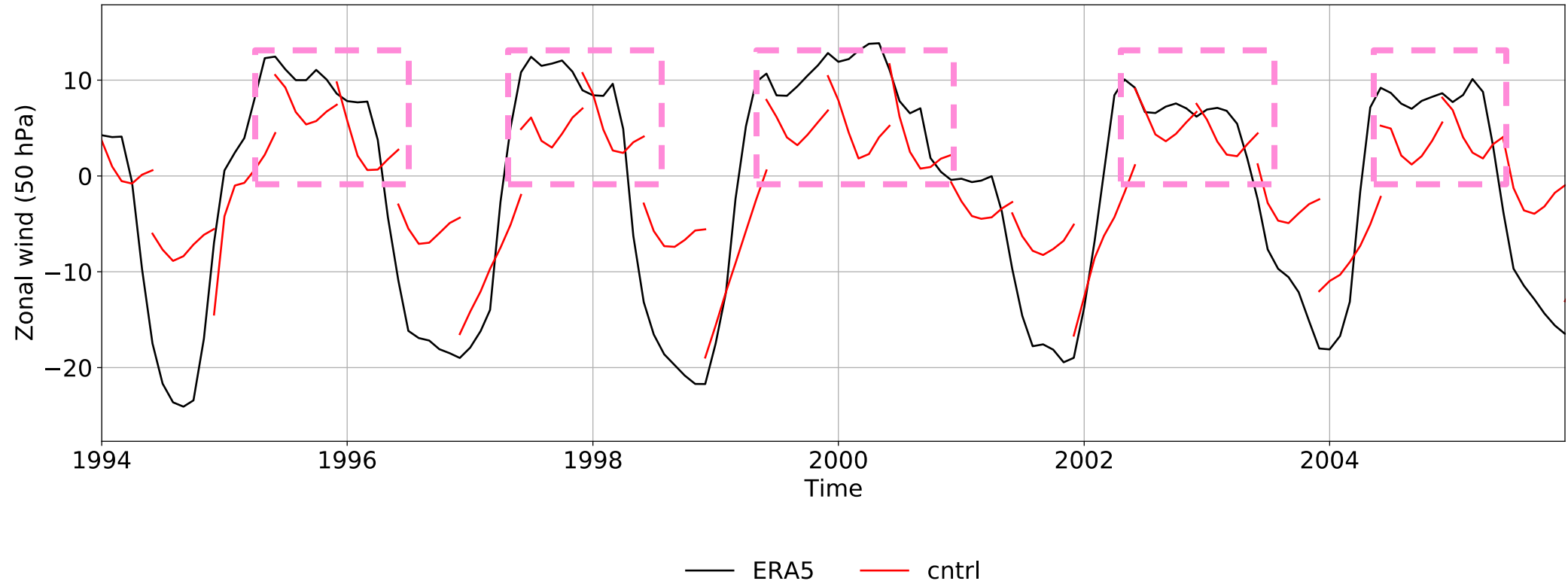
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# Turbulent diffusion in the stratosphere

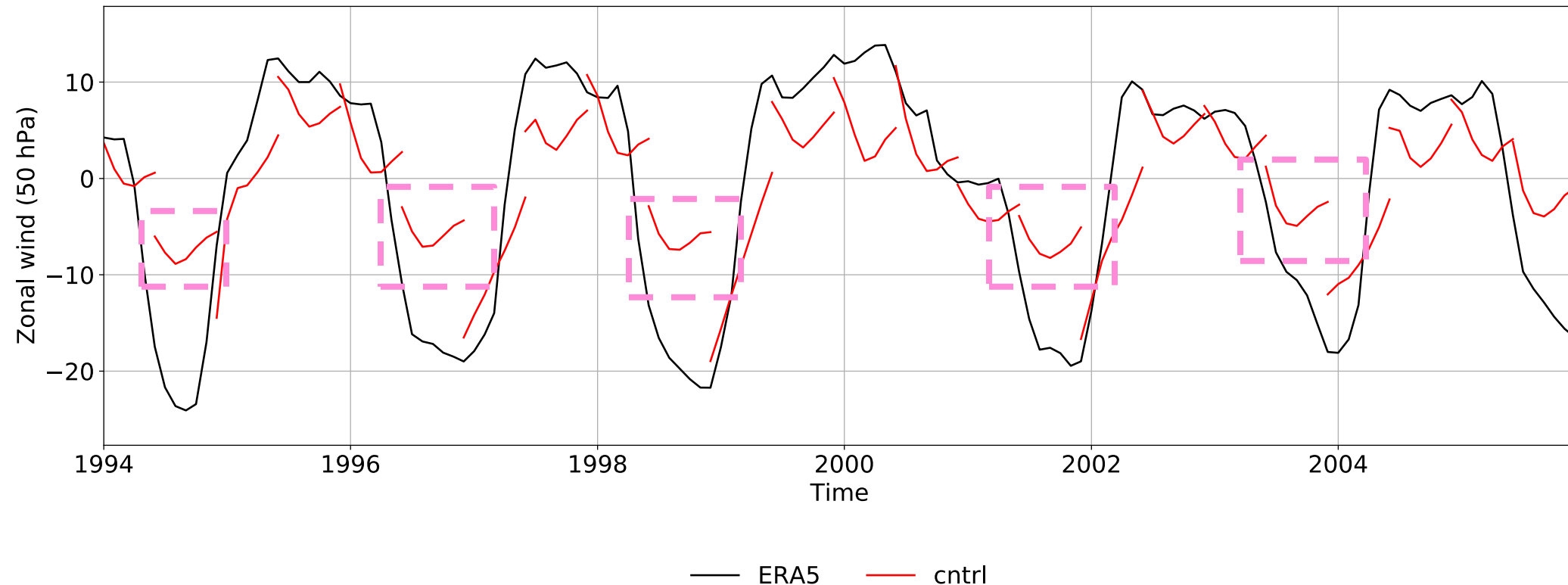
Current model has too weak winds in the QBO positive phase



Plot shows 50 hPa zonal winds averaged between 5S – 5N  
Seasonal hindcasts run with the ECMWF IFS, 7 months long

# Turbulent diffusion in the stratosphere

Current model has too weak winds in the QBO positive phase and negative phase

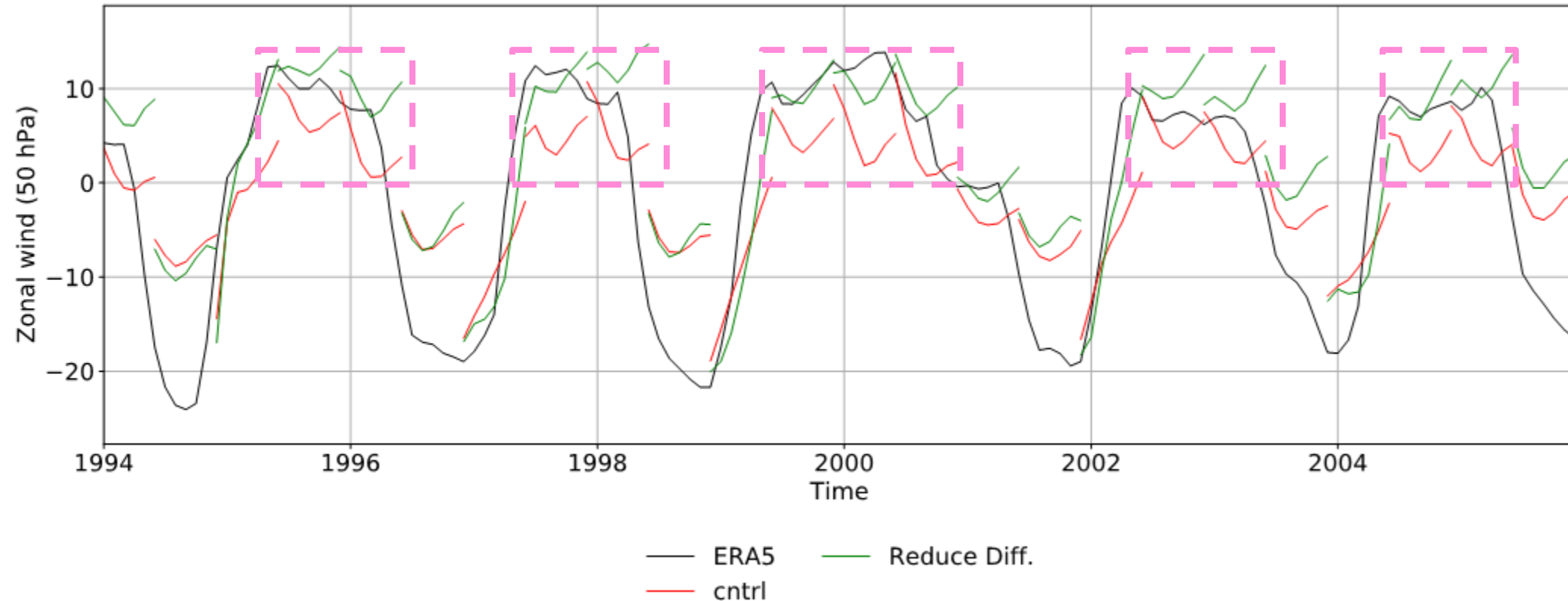


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Seasonal hindcasts run with the ECMWF IFS, 7 months long



# Turbulent diffusion in the stratosphere

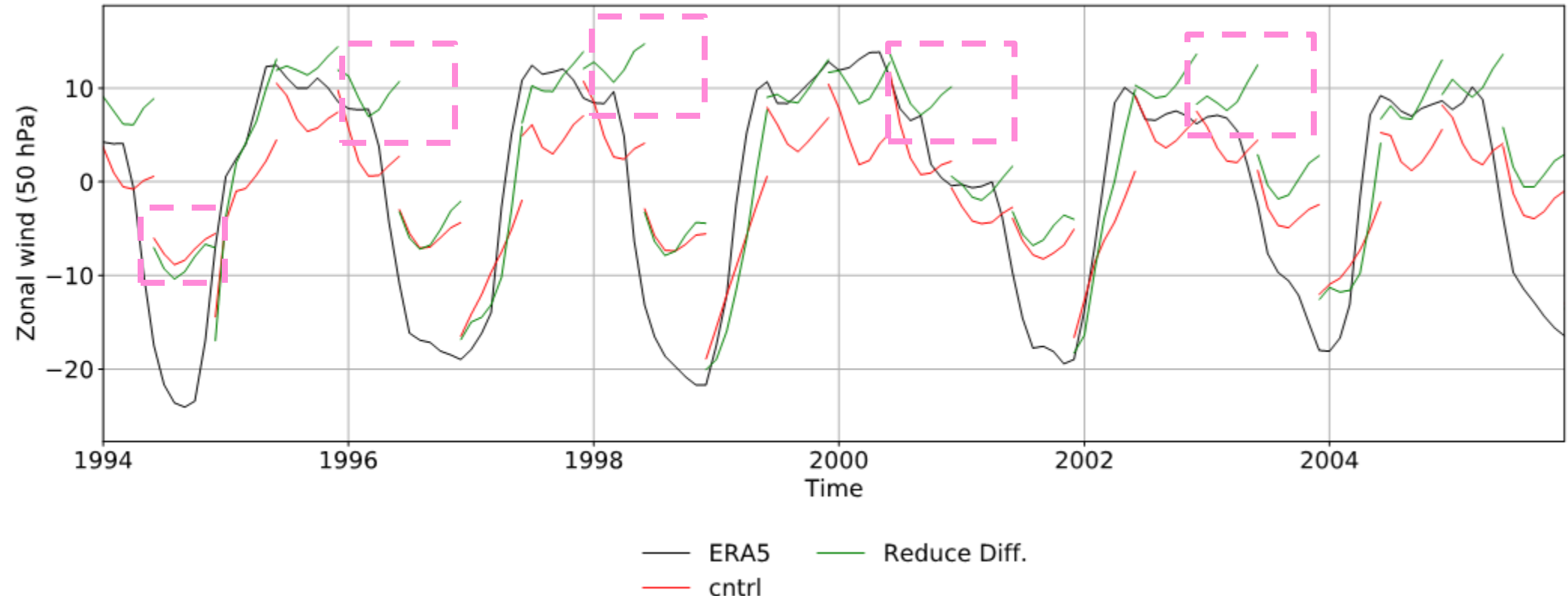
Reduced diffusion improves model winds in the QBO positive phase



Plot shows 50 hPa zonal winds averaged between 5S – 5N  
Seasonal hindcasts run with the ECMWF IFS, 7 months long

# Turbulent diffusion in the stratosphere

Reduced diffusion improves model winds in the QBO positive phase but does not make things better at the longer range



Plot shows 50 hPa zonal winds averaged between 5S – 5N  
Seasonal hindcasts run with the ECMWF IFS, 7 months long

# Turbulent kinetic energy (TKE) closure



# 'Local' turbulence closure: eddy diffusion above the surface

## Momentum

$$\overline{u'w'} \sim -K_M \frac{\partial \bar{u}}{\partial z}$$

## Thermodynamics

$$\overline{\theta'w'} \sim -K_H \frac{\partial \bar{\theta}}{\partial z}$$

$K_M, K_H$  and  $K_q$  are the exchange coefficients of momentum, heat and moisture

Their magnitude determines the transfer of these conserved quantities from turbulent eddies

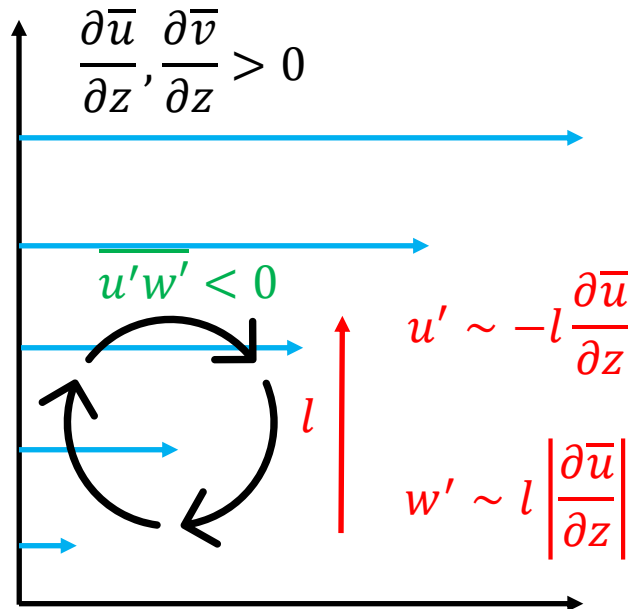
# 'Local' turbulence closure: eddy diffusion above the surface

## Momentum

$$\overline{u'w'} \sim -K_M \frac{\partial \bar{u}}{\partial z} = -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| f_M(Ri) \frac{\partial \bar{u}}{\partial z}$$

## Thermodynamics

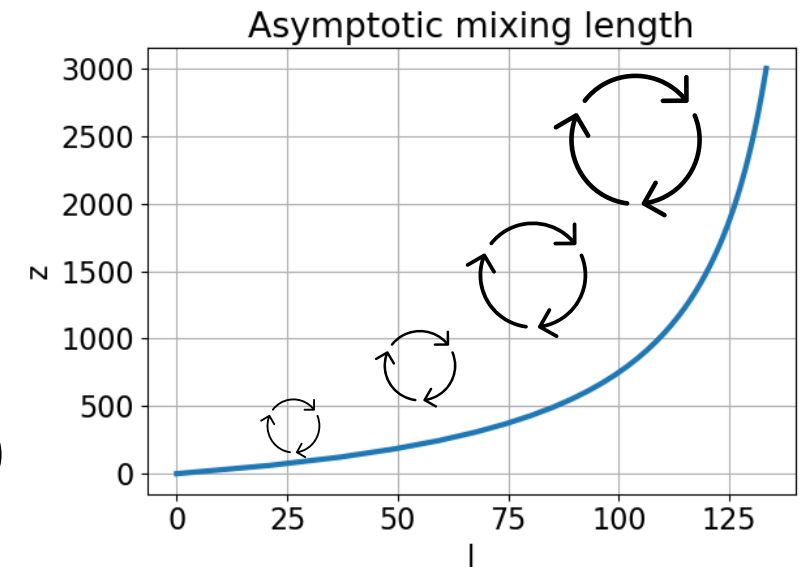
$$\overline{\theta'w'} \sim -K_H \frac{\partial \bar{\theta}}{\partial z} = -l^2 \left| \frac{\partial \bar{u}}{\partial z} \right| f_H(Ri) \frac{\partial \bar{\theta}}{\partial z}$$



Size of eddies get larger further away from the surface:

$$l \sim \frac{\kappa z \lambda}{\kappa z + \lambda}$$

$\kappa$  = von-Karman constant  
 $\lambda$  = asymptotic mixing length (150 m)



$f_M(Ri), f_H(Ri)$  determined empirically and depend on  $Ri(z)$ , since we are away from the surface

# Turbulent kinetic energy (TKE) closure

Momentum

$$\overline{u'w'} \sim -K_M \frac{\partial \bar{u}}{\partial z} = -C_k \chi_3(Ri_f^*) \sqrt{e_k} L_k \frac{\partial \bar{u}}{\partial z}$$

Thermodynamics

$$\overline{\theta'w'} \sim -K_H \frac{\partial \bar{\theta}}{\partial z} = -C_k C_3 \phi_3(Ri_f^*) \sqrt{e_k} L_k \frac{\partial \bar{\theta}}{\partial z}$$

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TKE  $\sqrt{e_k}$  - measure of the turbulence intensity

$C_k, C_3$  - closure constants

Stability functions  $\chi_3(Ri_f^*), \phi_3(Ri_f^*)$  - influence of stratification, uses flux Richardson number  $Ri_f^*$

Lengthscale  $L_k$  - defines the scale of the turbulence

# Turbulent kinetic energy (TKE) closure: it is prognostic

Advantage of the prognostic TKE is that it has 'memory', is advected and involves physical source terms :

$$\frac{\partial e_k}{\partial t} = \underbrace{-u \cdot \nabla e_k}_{\text{Advection}} - \underbrace{\frac{\partial}{\partial z} \left( K_{e_k} \frac{\partial e_k}{\partial z} \right)}_{\text{Turbulent diffusion}} - ST + BT - \epsilon_k$$

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TKE:

$$\sqrt{e_k} = \frac{\overline{u'u'} + \overline{v'v'} + \overline{w'w'}}{2}$$

Buoyancy production

$$BT = \frac{g}{\theta} \overline{\theta'w'} \approx -K_H \frac{g}{\theta} \frac{\partial \theta}{\partial z}$$

Shear production:

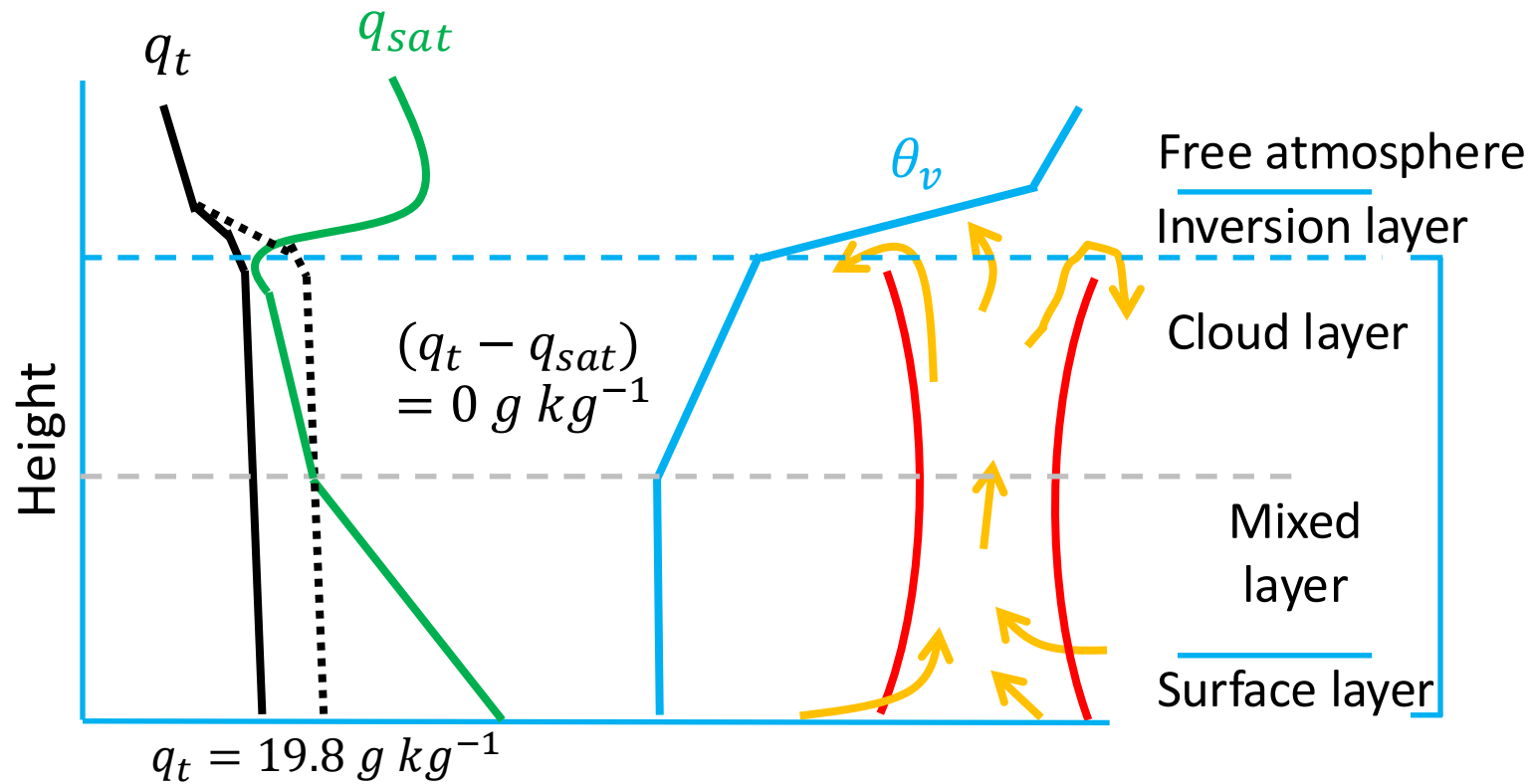
$$ST = -\overline{u'w'} \frac{\partial u}{\partial z} - \overline{v'w'} \frac{\partial v}{\partial z} \approx K_M \left| \frac{\partial u}{\partial z} \right|$$

Dissipation:

$$\epsilon_k = \frac{2e_k}{\tau_k} \approx C_\epsilon \frac{e_k^{\frac{2}{3}}}{L_\epsilon}$$

# Impact of TKE on low level cloud cover

# Stratoculums topped PBLs are very sensitive to mixing



The presence of stratocumulus is sensitive to:

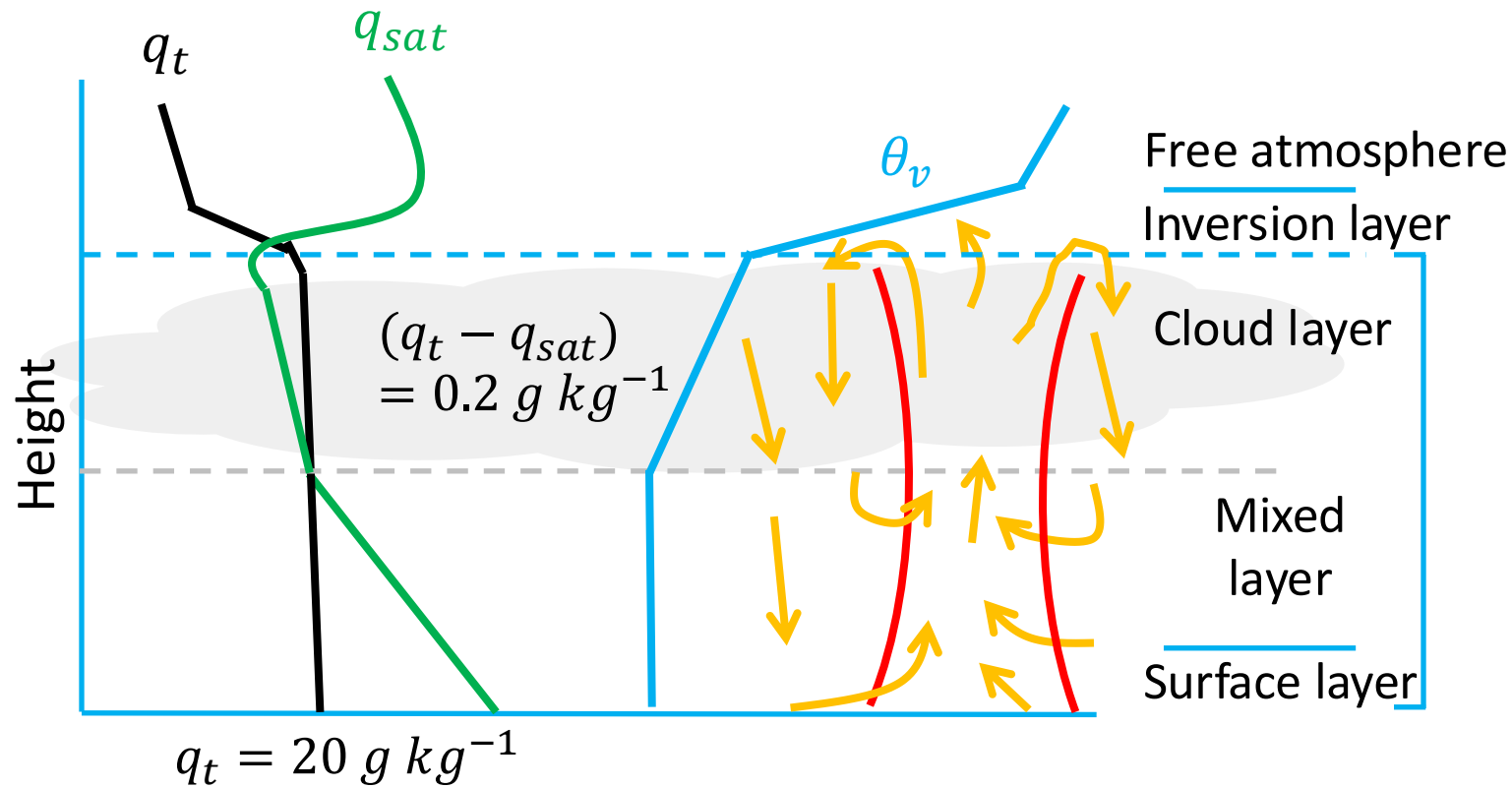
- Small variations in humidity

Mixing in stratocumulus clouds is more complex due to:

- Stronger entrainment from free atmosphere
- Condensation within cloud
- Radiative heating/cooling, which is essential for cloud evolution



# Stratoculums topped PBLs are very sensitive to mixing



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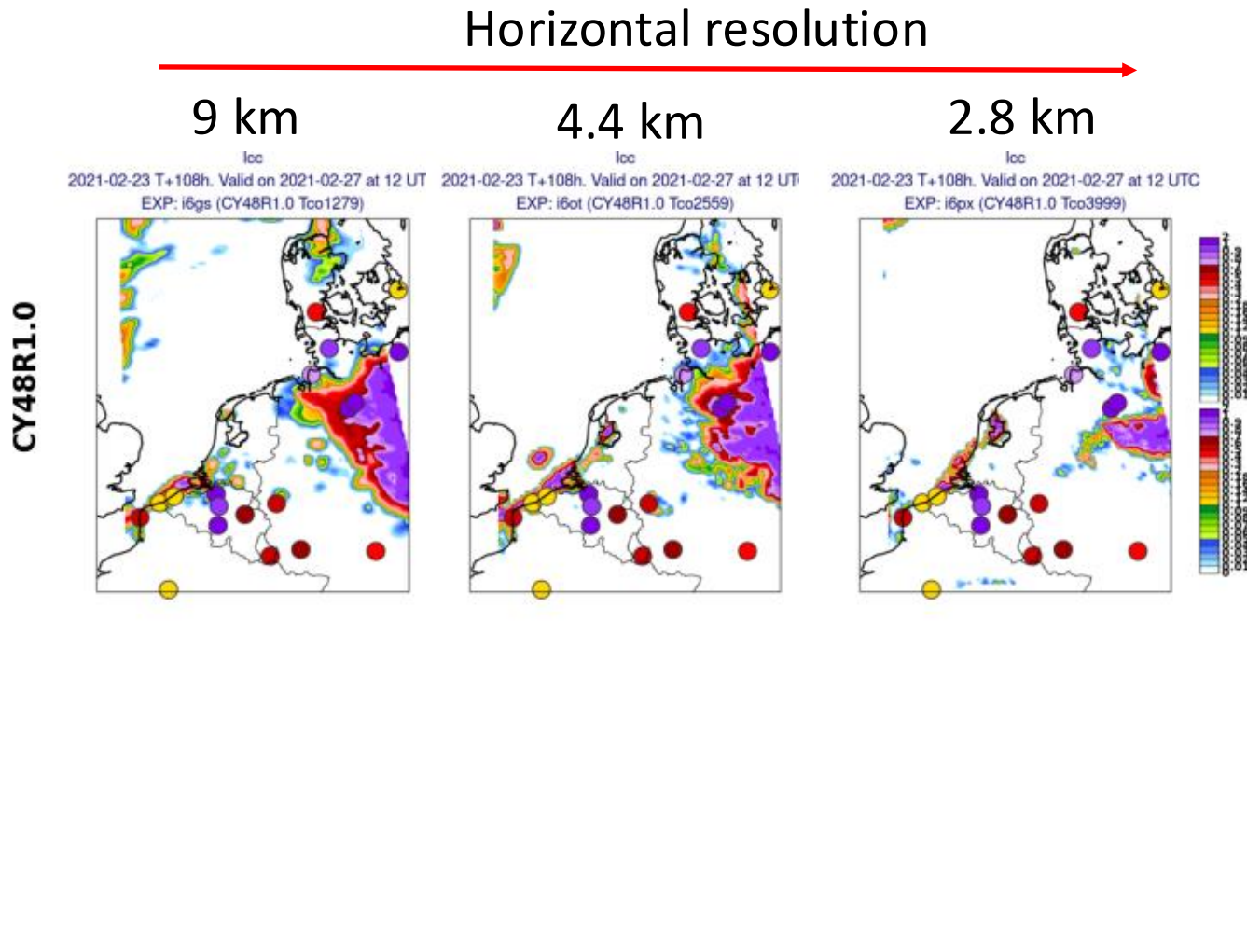
- Small variations in humidity
- Small variations in temperature

Mixing in stratocumulus clouds is more complex due to:

- Stronger entrainment from free atmosphere
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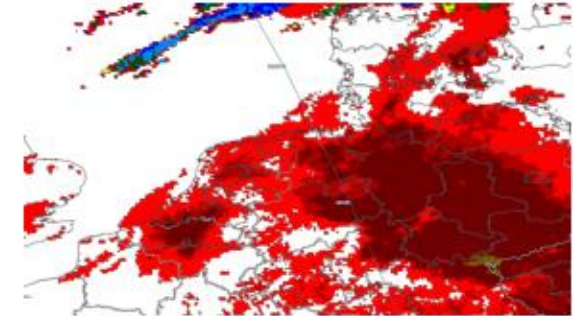
# Impact of TKE on low level clouds

Current  
turbulence  
scheme  
underestimates  
low cloud cover

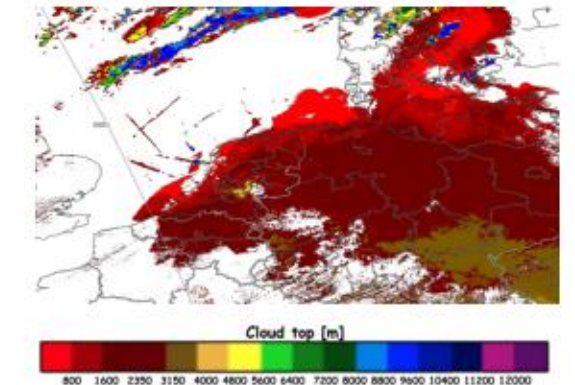


Observation of cloud top  
height

**CI Top Aqua 2021-02-27T12**



**CI Top NOAA 2021-02-27T12**



Figures c/o Ivan Bastak-  
Duran

# Impact of TKE on low level clouds

Horizontal resolution

9 km

4.4 km

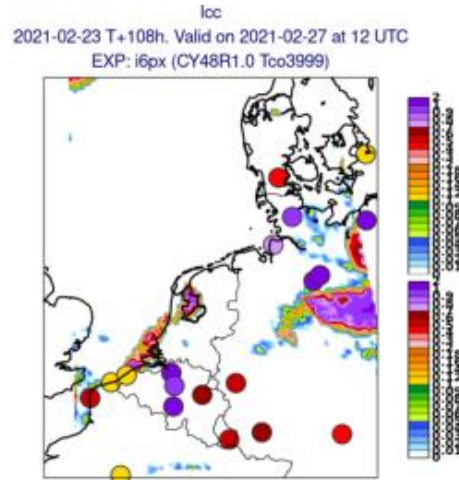
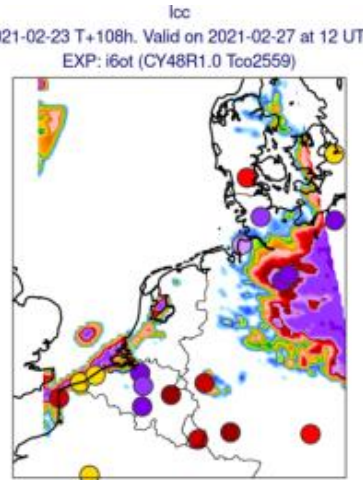
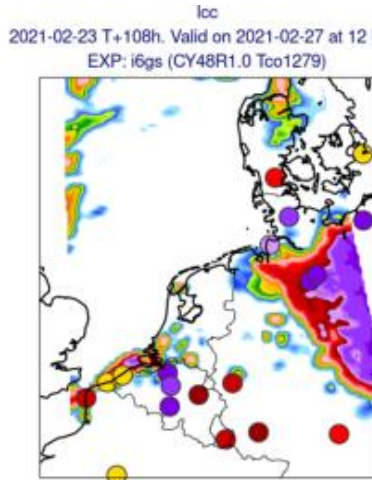
2.8 km

Observation of cloud top height

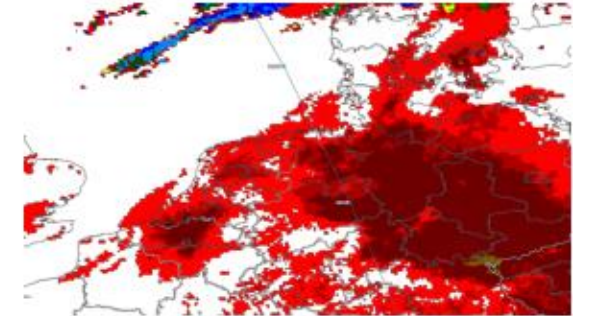
Current turbulence scheme underestimates low cloud cover

TKE scheme tends to have less mixing in these cases, and so can maintain low cloud

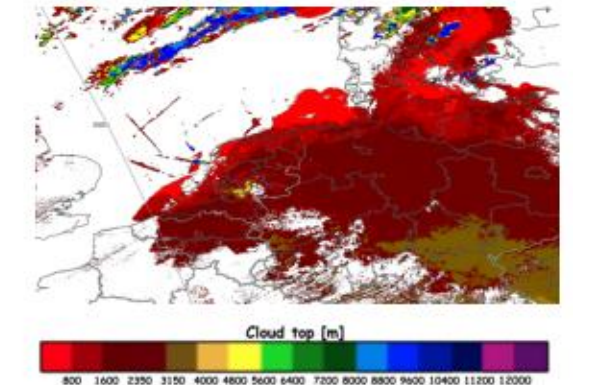
CY48R1.0



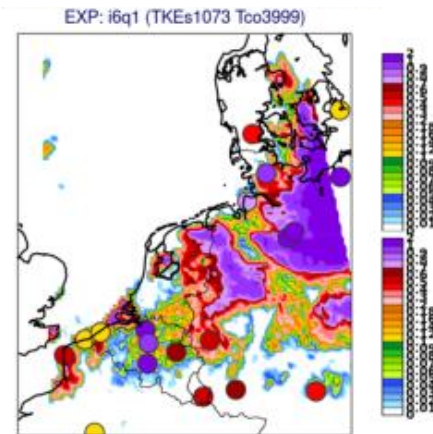
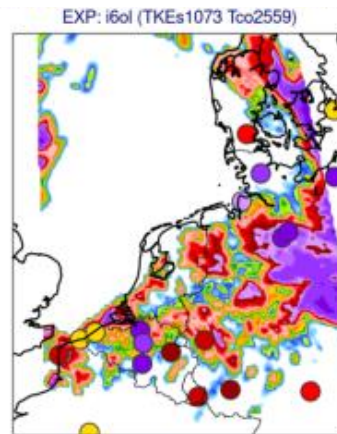
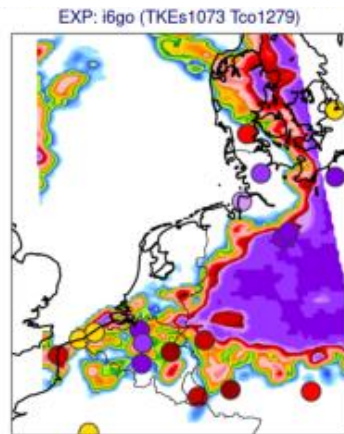
CI Top Aqua 2021-02-27T12



CI Top NOAA 2021-02-27T12



TKEs1073



Figures c/o Ivan Bastak-Duran



# Impact of TKE on low level clouds

Forecast lead time

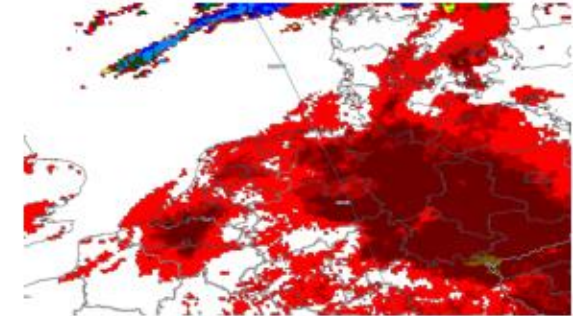
3 days

2 days

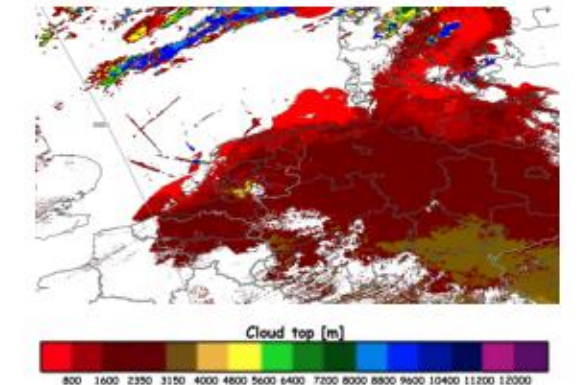
1 day

Observation of cloud top height

CI Top Aqua 2021-02-27T12



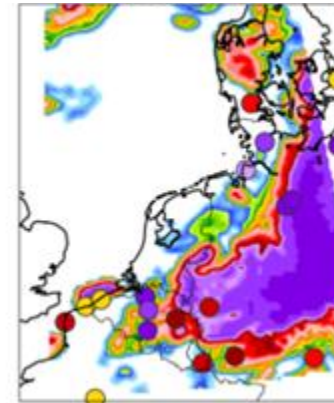
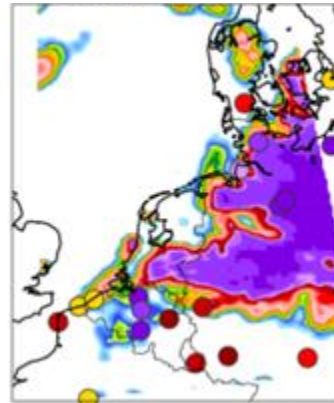
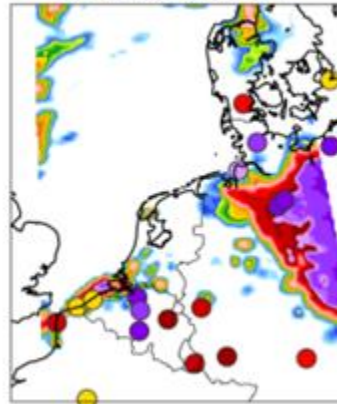
CI Top NOAA 2021-02-27T12



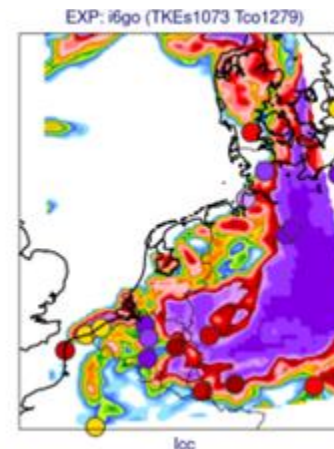
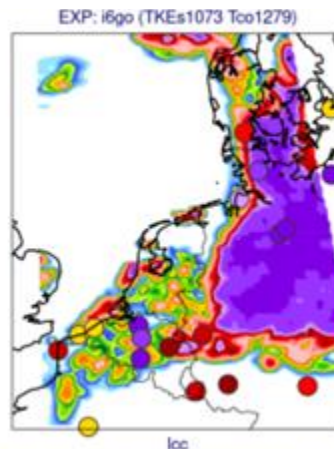
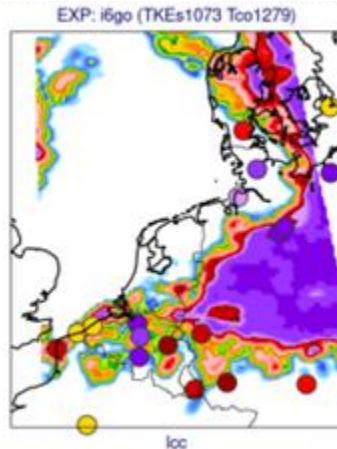
Current scheme cannot maintain low cloud – mixed too rapidly

TKE scheme has low cloud even at a lead time of 3 days

CY48R1.0



TKEs1073



2021-02-23 T+108h. Valid on 2021-02-27 at 12 UTC 2021-02-24 T+84h. Valid on 2021-02-27 at 12 UTC 2021-02-25 T+60h. Valid on 2021-02-27 at 12 UTC

Figures c/o Ivan Bastak-Duran

# Summary of BL parametrization

- **Outer layer diffusion and mass flux:**
  - Local similarity (derived from surface observations) is used in unstable regimes but does not work for stable
  - Local diffusion does not work for convective boundary layers, which is where mass flux is also added
- **IFS parametrization:**
  - Due to the uncertainty in the stability functions, different forms are used throughout the atmosphere
- **Sensitivity to changing stability functions:**
  - Reverting the stability function to their 'empirical' form degrades the forecast, due to reduced mixing
  - However, less mixing in the stratosphere improves the winds in the tropics
- **TKE:**
  - The TKE scheme benefits from having memory and being advected by the flow
  - TKE improves the representation of low cloud cover